

The University of Manchester  
Photon Science  
Institute

# The Photon Science Institute Building






MANCHESTER  
1824

The University of Manchester

# The AMPPS Project

Providing cutting-edge facilities for the School of Physics and Astronomy, the School of Mathematics and the Photon Science Institute

Estates Contact: Dave Smith 0161 275 2266



northwest development agency hbg FABER MAUNSELL AECOM SHEPPARD ROBSON CAPITA SYMONDS IRW JACOBS BABTIE mda CONSULTING



Photographs courtesy of Nicholas J Higham (c) 2006







# **ZEKE Rydberg states in a crowd and **C**ondensed **R**ydberg **C**lusters: a new state of matter?**

10th October 2008

Department of Physics

**University of Virginia, Charlottesville**

**Klaus Müller-Dethlefs**

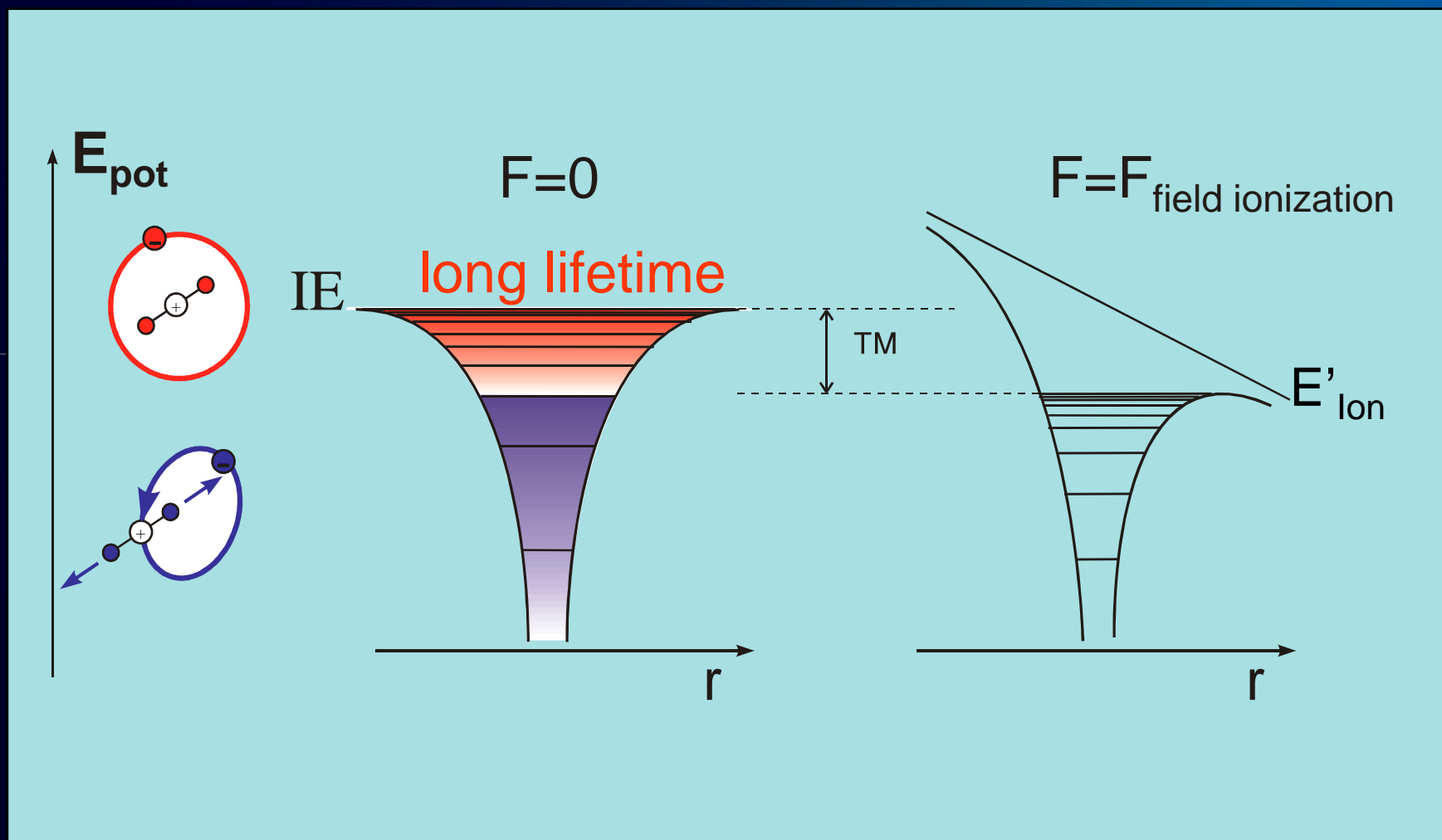
**The Photon Science Institute**

**School of Physics & Astronomy, School of Chemistry**

**The University of Manchester, United Kingdom**



# Pulsed Field Ionisation of ZEKE Rydberg States



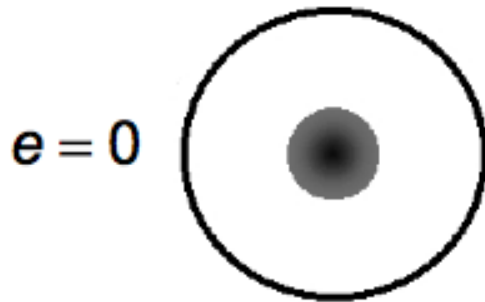
KMD et al CPL 1984, CPL 1988



## Stabilization via Angular Momentum Enhancement, Equal to Ultra-long Lifetimes

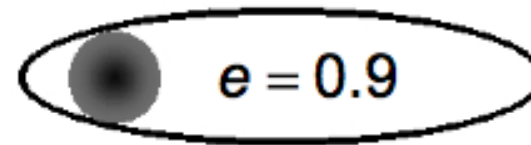
$$e = \sqrt{1 - \frac{\ell^2}{n^2}}$$

Eccentricity: It gives the shape of the orbit



Large  $\ell$  orbit: no collisional energy exchanges between the Rydberg electron and the core

**Long Rydberg Lifetime ( $\sim n^5$ )**



Small  $\ell$  orbit: many collisional energy exchanges between the Rydberg electron and the core

**Lifetime due to core-collisions  $\sim n^3$**

**Radiative lifetimes are also longer for large  $\ell$  orbits )**



## Energy States of H-atom

$$E = IE_H - R_H / n^2$$

## Energy States of “H-atom like” Rydberg State

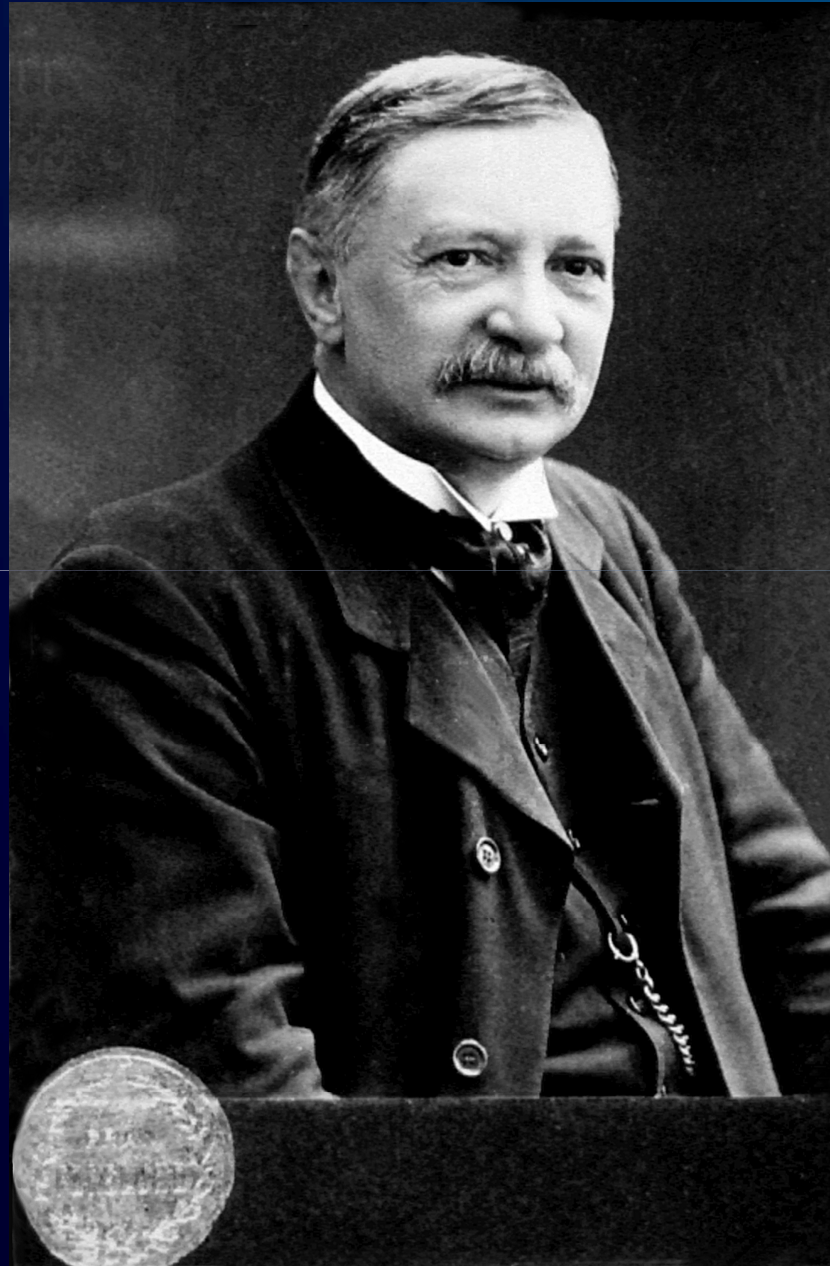
$$E = IE - R / (n - \delta)^2$$

*n*: principal quantum number

*δ* : quantum defect



# Rydberg





# *Setting the Scene* for a half day experiment

- ZEKE Rydberg states
  - one colour ZEKE Rydberg excitation
    - para-difluorobenzene
- Supersonic Jet Excitation
  - close to nozzle at high density
  - high collisional cooling
- Rydberg electron sees many ions!!
  - electrons should be bound collectively!



# *Setting the Scene* for a half day experiment

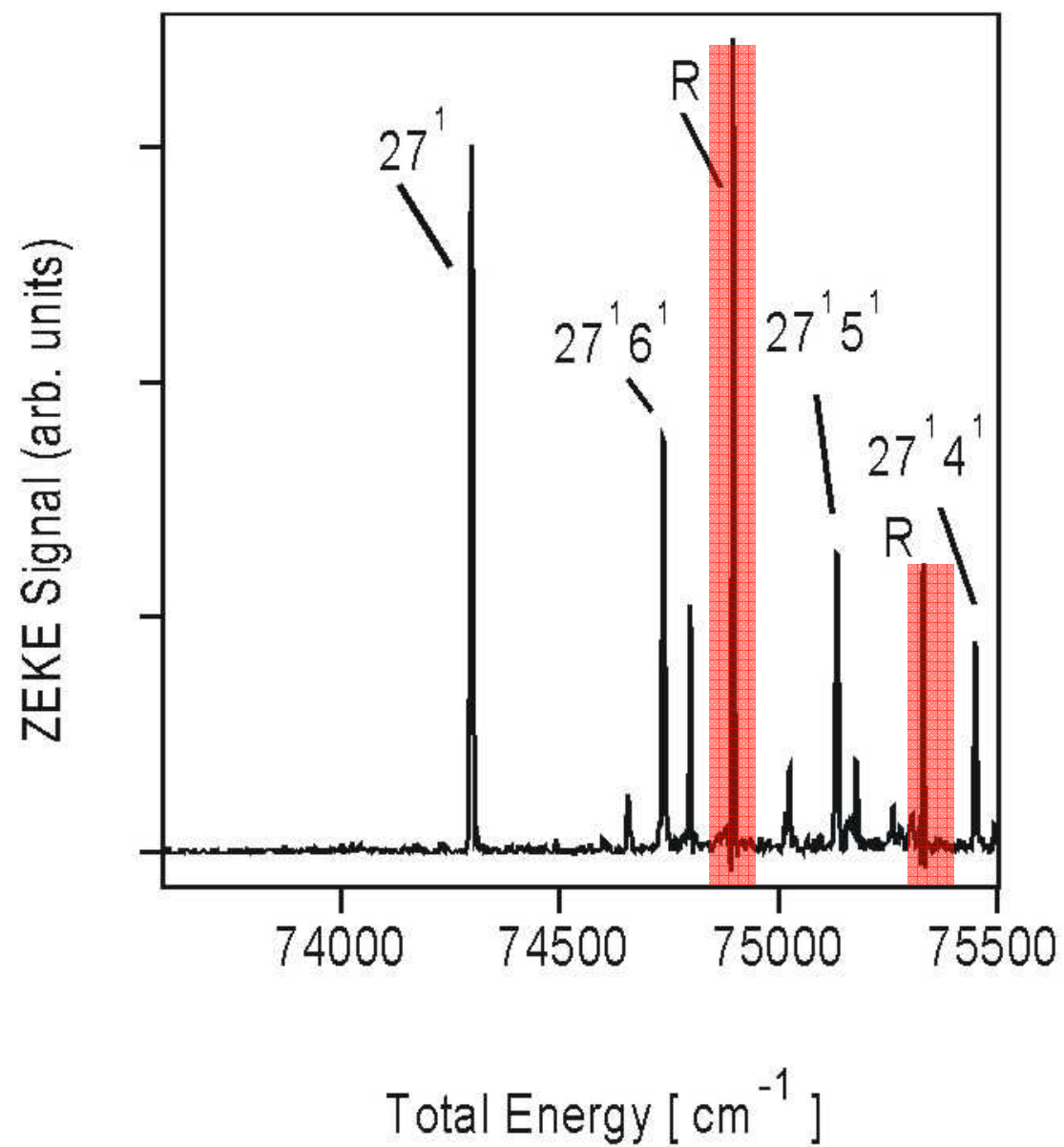
- ZEKE Rydberg states
  - one colour ZEKE Rydberg excitation
    - para-difluorobenzene (pDFB)

As shown by Reiser et al, *JPC* 1991,  
pDFB shows some very strong one-colour transitions  
where the  $S_1 \rightarrow S_0$  and the  $D_0 \rightarrow S_1$   
transition are simultaneously in resonance

So produce very highly excited ZEKE Rydberg states  
 $n > 200$  with only one laser!!



p-DFB:  $D_0 \leftarrow S_1 27^1$



# *Setting the Scene* for a half day experiment

- ZEKE Rydberg states
  - one colour ZEKE Rydberg excitation
    - para-difluorobenzene (pDFB)
- **Supersonic Jet Excitation**
  - close to nozzle at high density
  - high collisional cooling
- Rydberg electron sees many ions!!
  - electrons should be bound collectively!



## Jet characteristics:

at nozzle:  $2.5 \times 10^{19} \text{ cm}^{-3}$

pDFB density at nozzle:  $2.5 \times 10^{17} \text{ cm}^{-3}$

pDFB density at 1mm distance.:  $2.5 \times 10^{15} \text{ cm}^{-3}$

pDFB<sup>+</sup> maximum density:  $2.5 \times 10^{15} \text{ cm}^{-3}$

$10^5 \text{ ions/cm}$

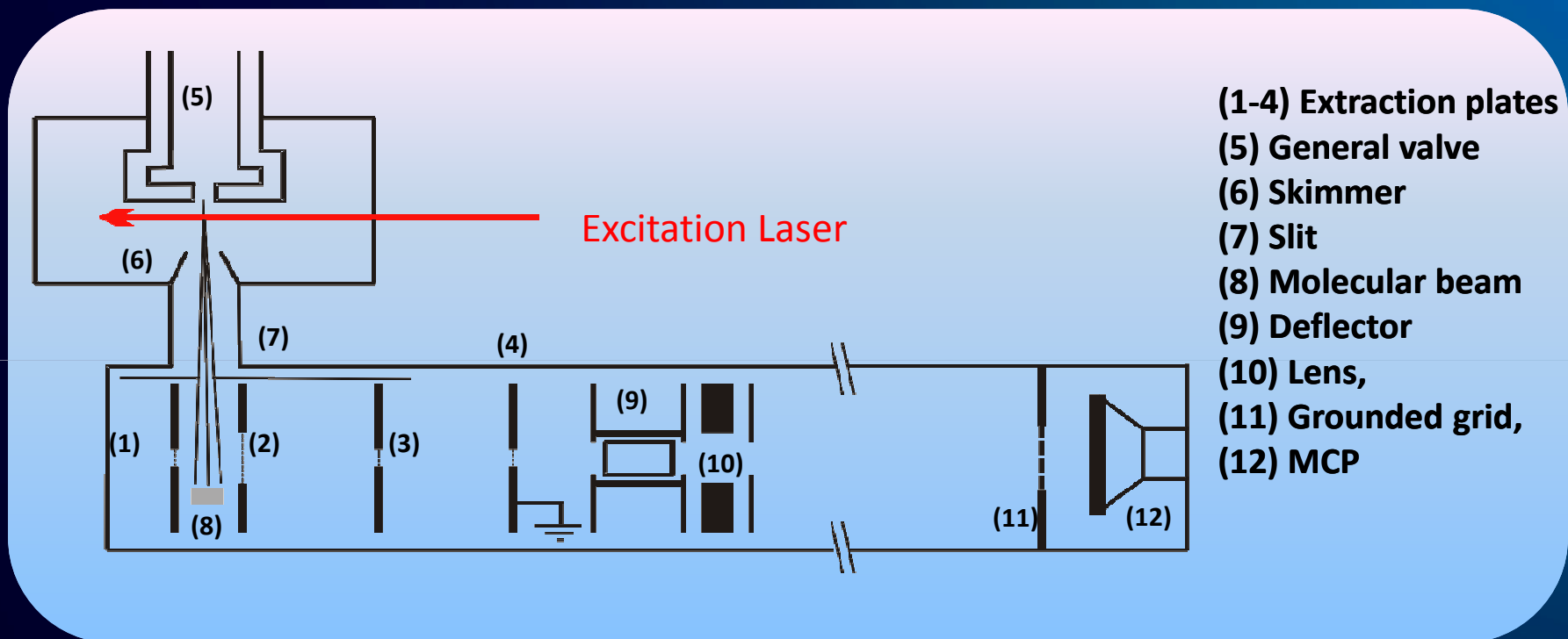
$10^4 \text{ ions/mm}$

$10 \text{ ions/}\mu\text{m}$

pDFB<sup>+</sup> average distance: *ca.* hundreds of nm's

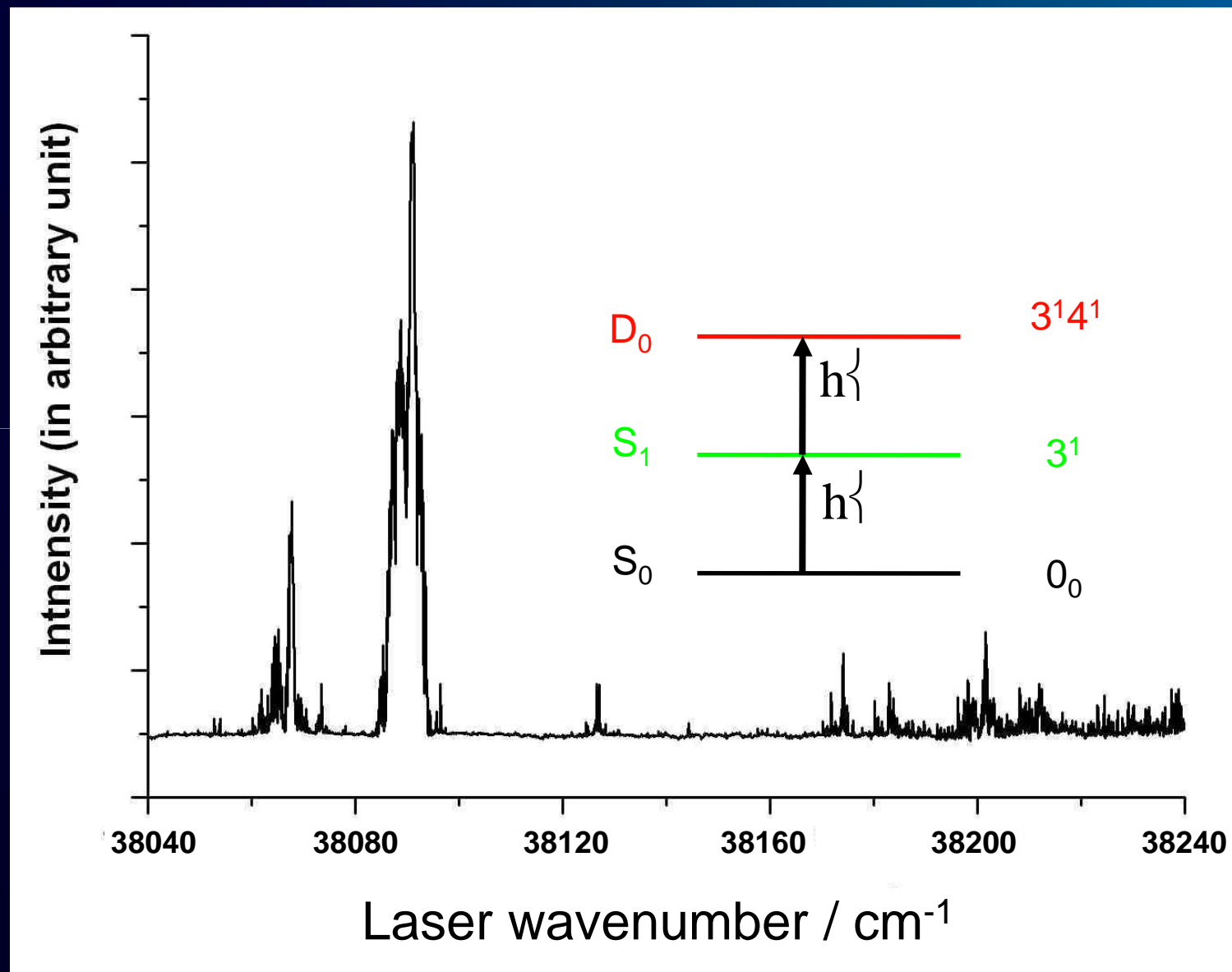
If something new such as *Condensed Rydberg Clusters* are formed, the expected distance between ion cores is expected to be around hundreds of nanometers, if  $\delta_{\text{pDFB}}$  is *ca.*  $10^{13}$  to  $10^{14} \text{ cm}^{-3}$

# Schematic diagram of the experimental set up for the TOF analysis of a molecular ultracold Rydberg plasma.

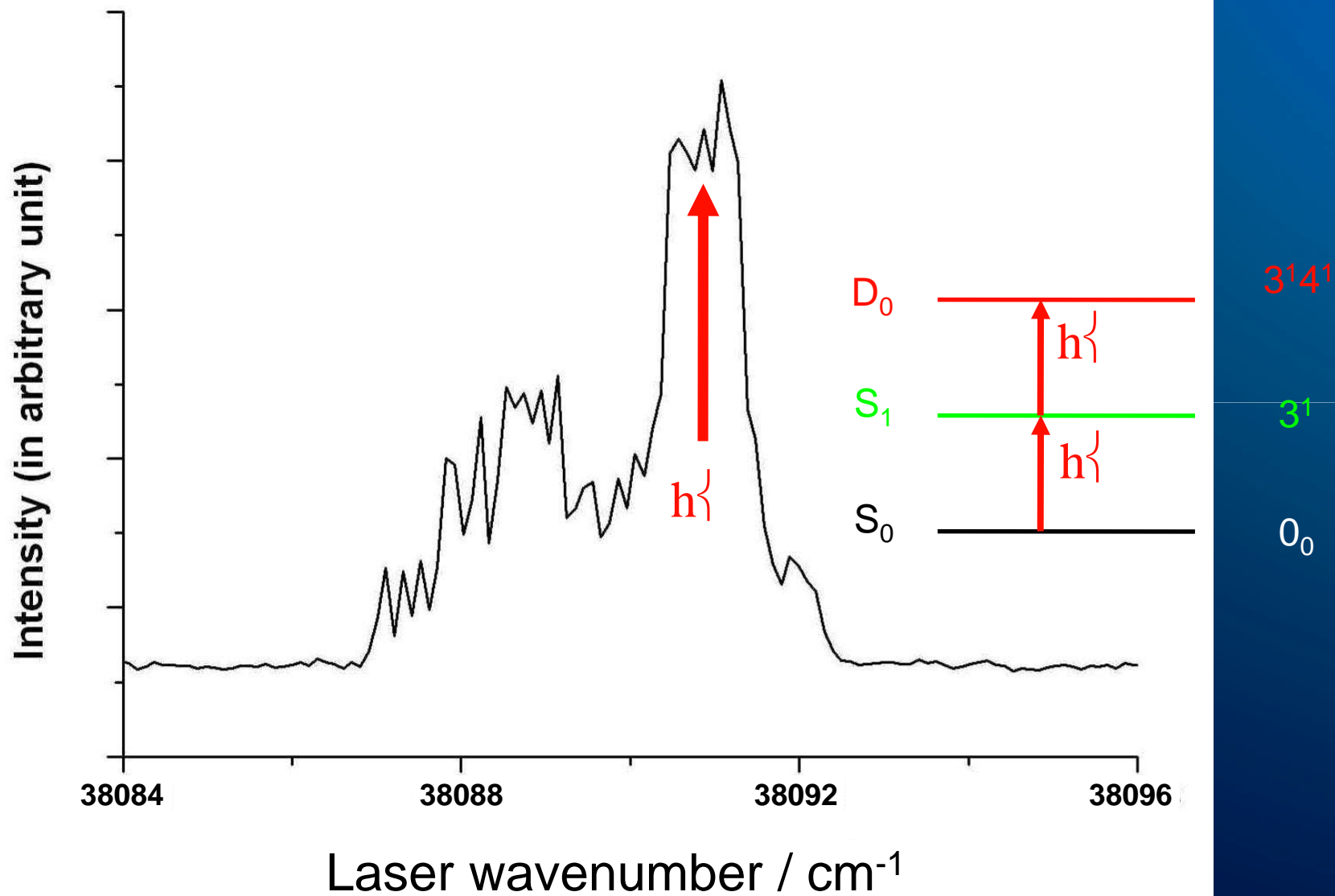




Para-difluorobenzene (pDFB) one-colour spectrum with simultaneous  $S_1$   $S_0$  and ZEKE Rydberg resonance



Para-difluorobenzene (pDFB) one-colour spectrum with simultaneous  $S_1$   $S_0$  and ZEKE Rydberg resonance



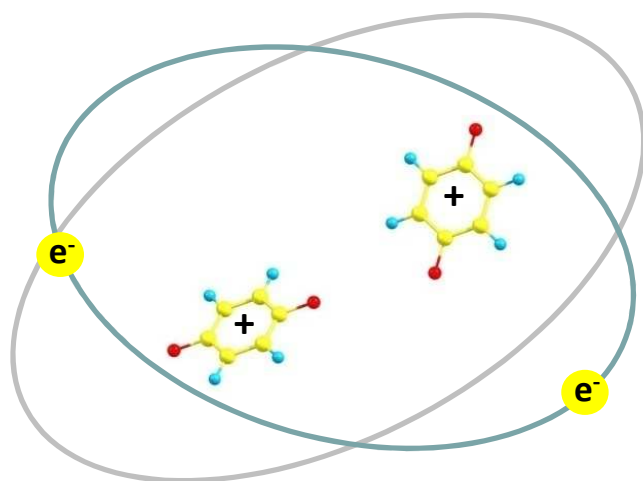


# *Setting the Scene* for a half day experiment

- ZEKE Rydberg states
  - one colour ZEKE Rydberg excitation
    - para-difluorobenzene (pDFB)
- Supersonic Jet Excitation
  - close to nozzle at high density
  - high collisional cooling
- Rydberg electron sees many ions!!
  - electrons should be bound collectively!

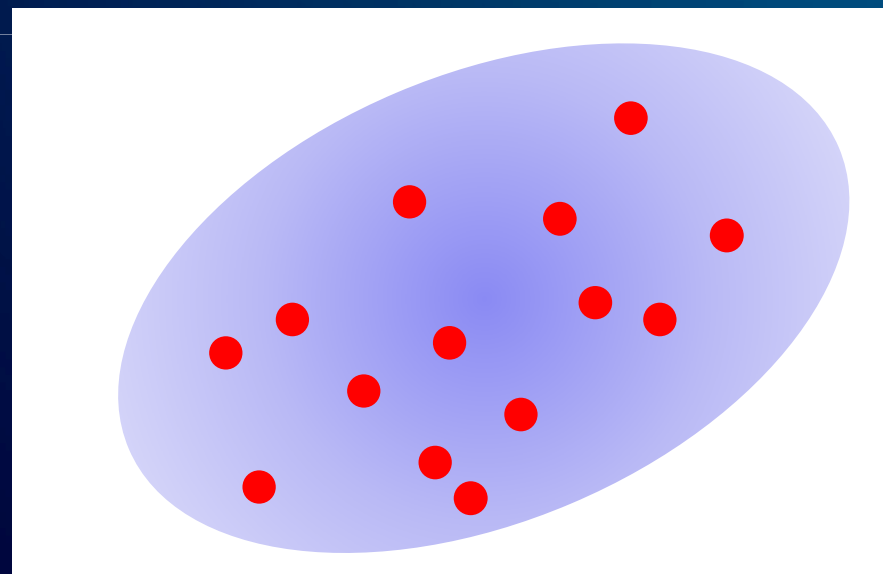
# Concept of CRC

We postulate **Condensed Rydberg Clusters (CRC)** as a moiety of  $n$  cations bound collectively by approximately the same number of electrons. The average distances between adjacent ionic molecular cores are assumed to be much smaller than the size of the Rydberg orbital, so that every Rydberg electron interacts with many cation cores.



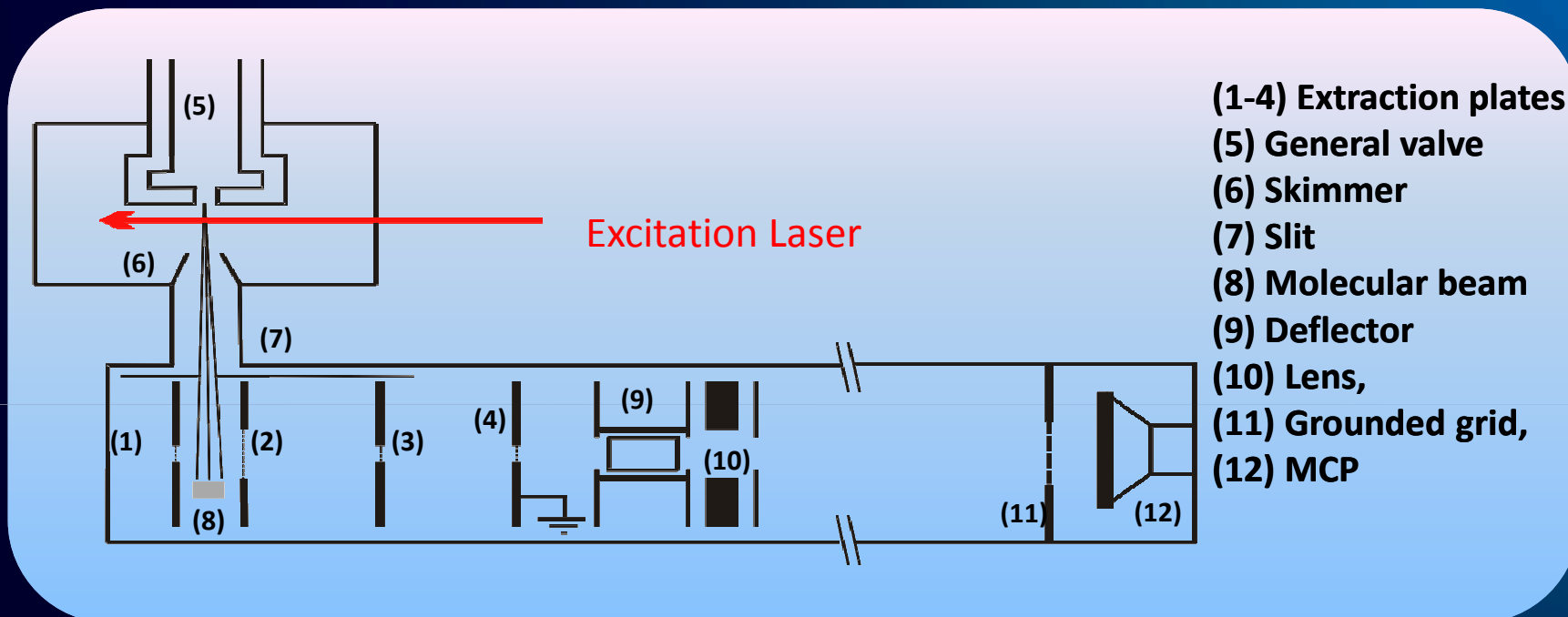
para-difluorobenzene

(pDFB)

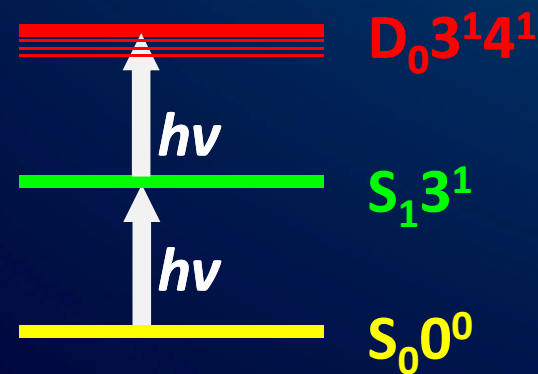


**Condensed Rydberg Cluster**

# Generation of Ultracold Molecular Plasma and CRC

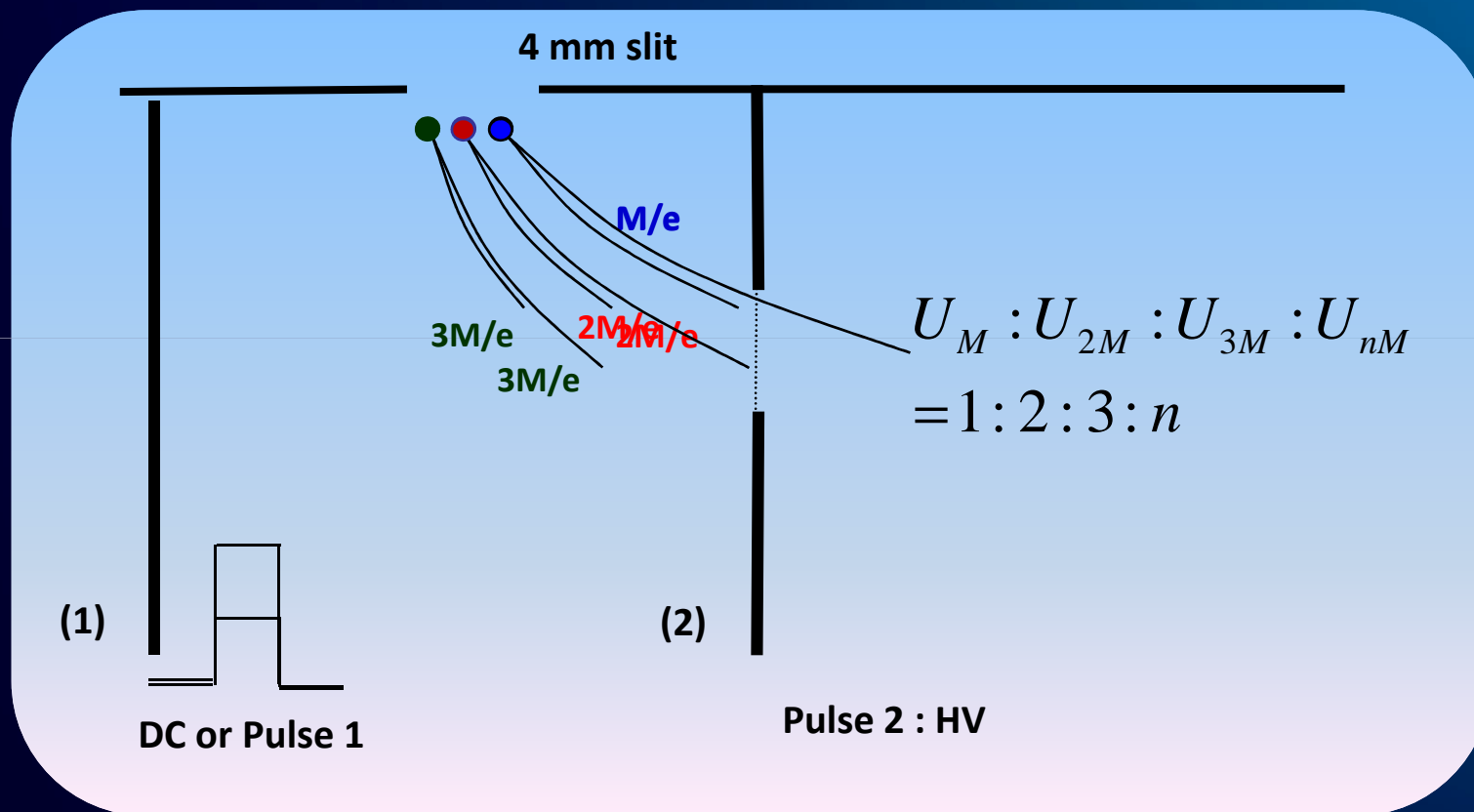


Density of  $\text{pDFB}^+$  up to  $10^{14}/\text{cm}^3$   
 ( $\Gamma_{\text{ion}} > 100$ ,  $\Gamma_{\text{electron}} > 40$ )





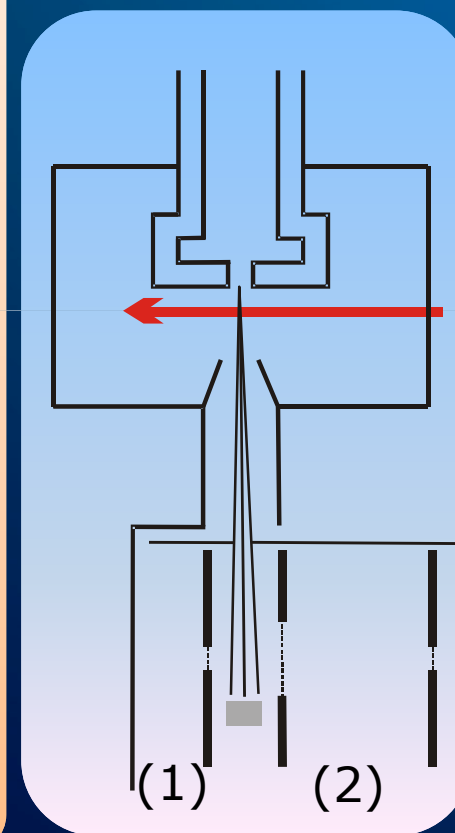
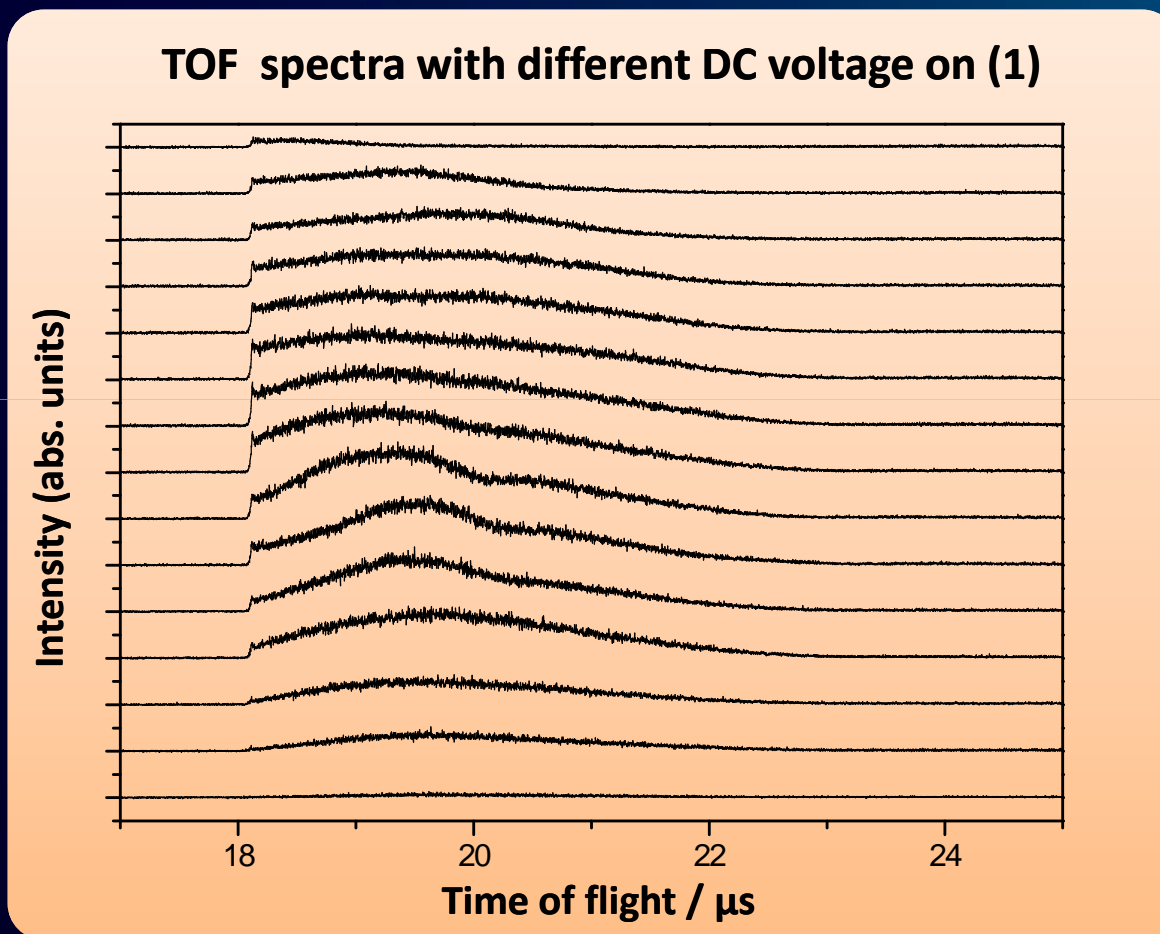
# Detection of CRC “Heavy Mass”



Method 1: Scan of **dc voltage** on aperture (1)

# Detection of “Heavy Mass”

## Results (1): Voltage scan on aperture 1



+4.8 kV pulse is applied on (2), pulse width is 1.22  $\mu\text{s}$ , laser crosses with molecular beam at 6 mm from the nozzle

# Possible signal source

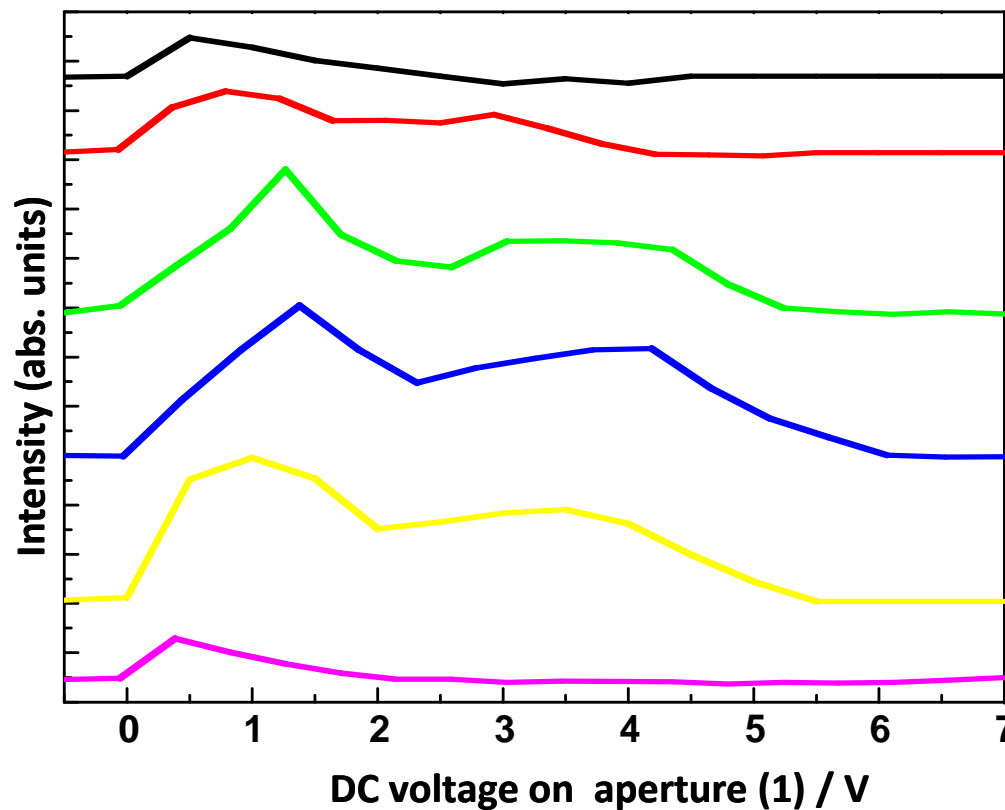
- Single pDFB cations ✗
- Single pDFB Rydberg's ✗
- Cations released from the plasma ✓
- something new released from or  
contained in the plasma: ✓  
**Condensed Rydberg Clusters CRC**



# Detection of “Heavy Mass”

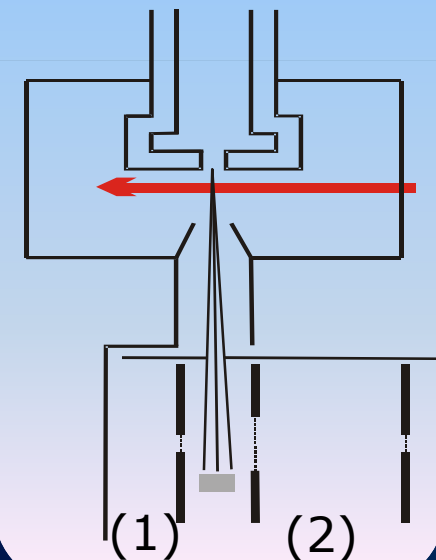
## Laser in Jet Position Sensitivity

DC voltage on (1) scan with different laser positions



Distance between Laser and Nozzle

$$\Gamma = \infty \frac{D}{T}$$



4.8 kV pulse is applied on (2), pulse width is 1.22  $\mu$ s, the delay between laser and pulse is 450  $\mu$ s.

# Questions arising:

- Plasma Formation?! ✓
- Can the “plasma” have a structure? ✓?
- Ultracold “condensate” of  **$m$  ions and  $n$  electrons** ( $m \approx n$ ) with periodic ion structure and electrons in a “conduction band”

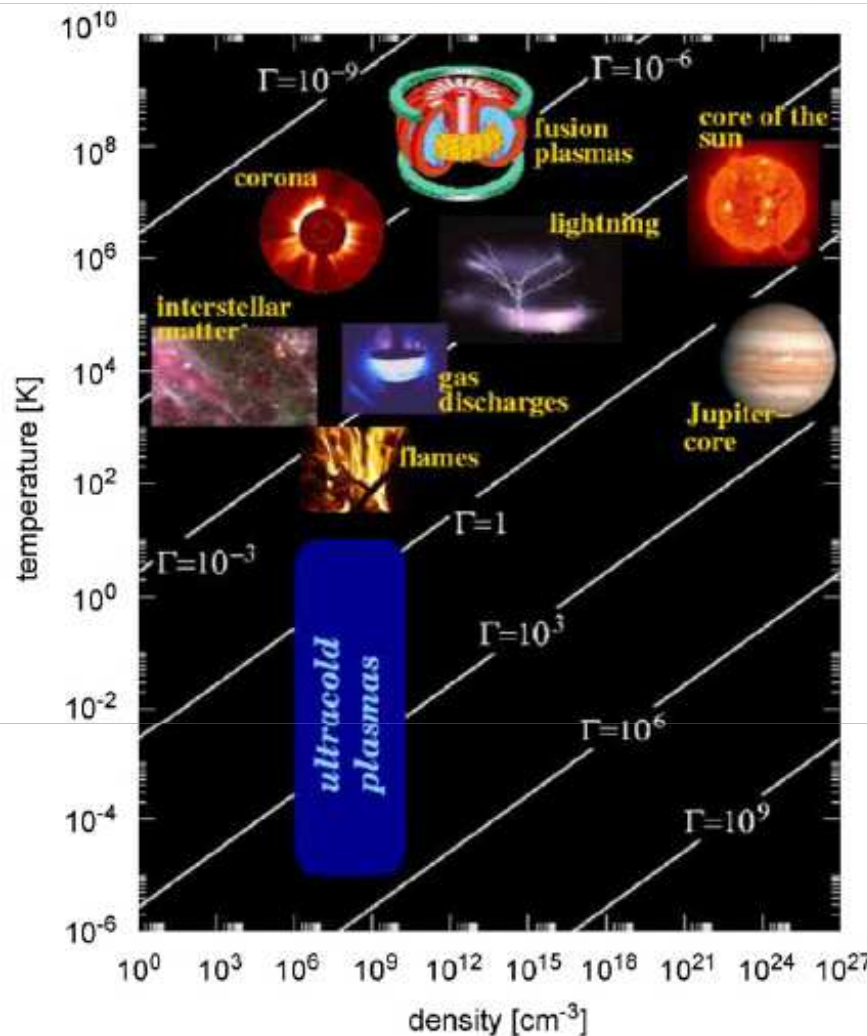
Contrast this to

- Cold Plasmas from MOT plus Rydberg excitation
- van der Waals clusters ✗

# Overview of neutral plasmas in the density-temperature plane

T.C. Killian etc., Phys.Report 449 (2007) 77

T. C. Killian, Science, 316 (2007) 705



$$\Gamma = \frac{E_{\text{interaction}}}{E_{\text{thermal}}} = 2.69 \times 10^{-5} \left[ \frac{n}{10^{12} \text{ cm}^{-3}} \right]^{1/3} \left[ \frac{T}{10^6 \text{ K}} \right]^{-1}$$

$E_{\text{interaction}}$ : Coulomb interaction energy of the plasma;  $E_{\text{thermal}}$ : thermal energy of ions or electrons.  $n$  is number density of plasma.



# How to obtain a *strongly coupled* Rydberg plasma

Plasma coupling parameter  $\Gamma$

$$\Gamma = \frac{e^2}{4\pi\epsilon_0 a} / k_B T$$

The experimental goal is to achieve strongly coupled plasma conditions for which the coulomb energy is much bigger than  $kT$ . Such a plasma could have a periodic shell structure and would have the characteristics of a metal

**Lower temperature and/or increase density!!**

Prediction is that  $\Gamma$  should be above 750 for crystalline phase

# How to obtain a *strongly coupled* Rydberg plasma

## Other efforts

(Killian, Gallagher, Pillet, Weidemüller...)

- Magneto-Optical trap to cool down atoms to microKelvin's.
- Density of ions is only up to  $10^9/\text{cm}^3$
- Applies to atomic systems only

## Our method

- Laser crosses with gas jet in the free expansion region ( $T \sim 0.1 \text{ K}$ )
- Density of ions can be higher than  $10^{14}/\text{cm}^3$ .
- Applies to molecular and atomic systems.

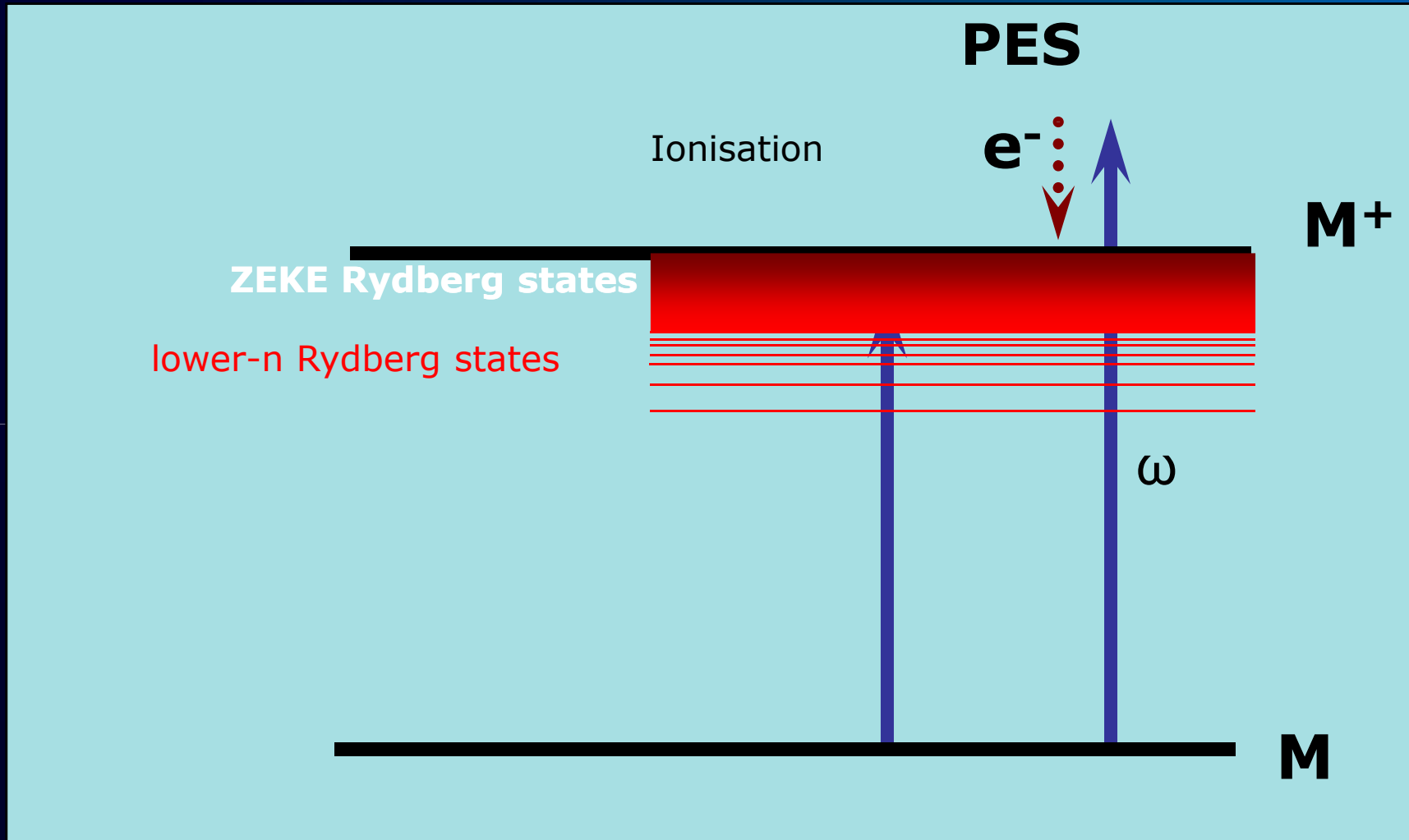
## References:

S. Kulin, T. C. Killian, S. D. Bergeson, and S. L. Rolston, Phys. Rev. Lett., 85 (2000) 318.

T. C. Killian, M. J. Lim, S. Kulin, R. Dumke, S. D. Bergeson, and S. L. Rolston, Phys. Rev. Lett., 86 (2001) 3759.

T. C. Killian, S. Kulin, S. D. Bergeson, L. A. Orozco, C. Orzel and S. L. Rolston, Phys. Rev. Lett., 83 (1999) 4776.

# Adjust the energy of the electron and hence the electron temperature





# Questions arising:

- Plasma Formation?!
- Can the “plasma” have a structure?
- Ultracold “condensate” of  **$m$  ions and  $n$  electrons ( $m \approx n$ )** with periodic ion structure and electrons in a “conduction band”

Contrast this to

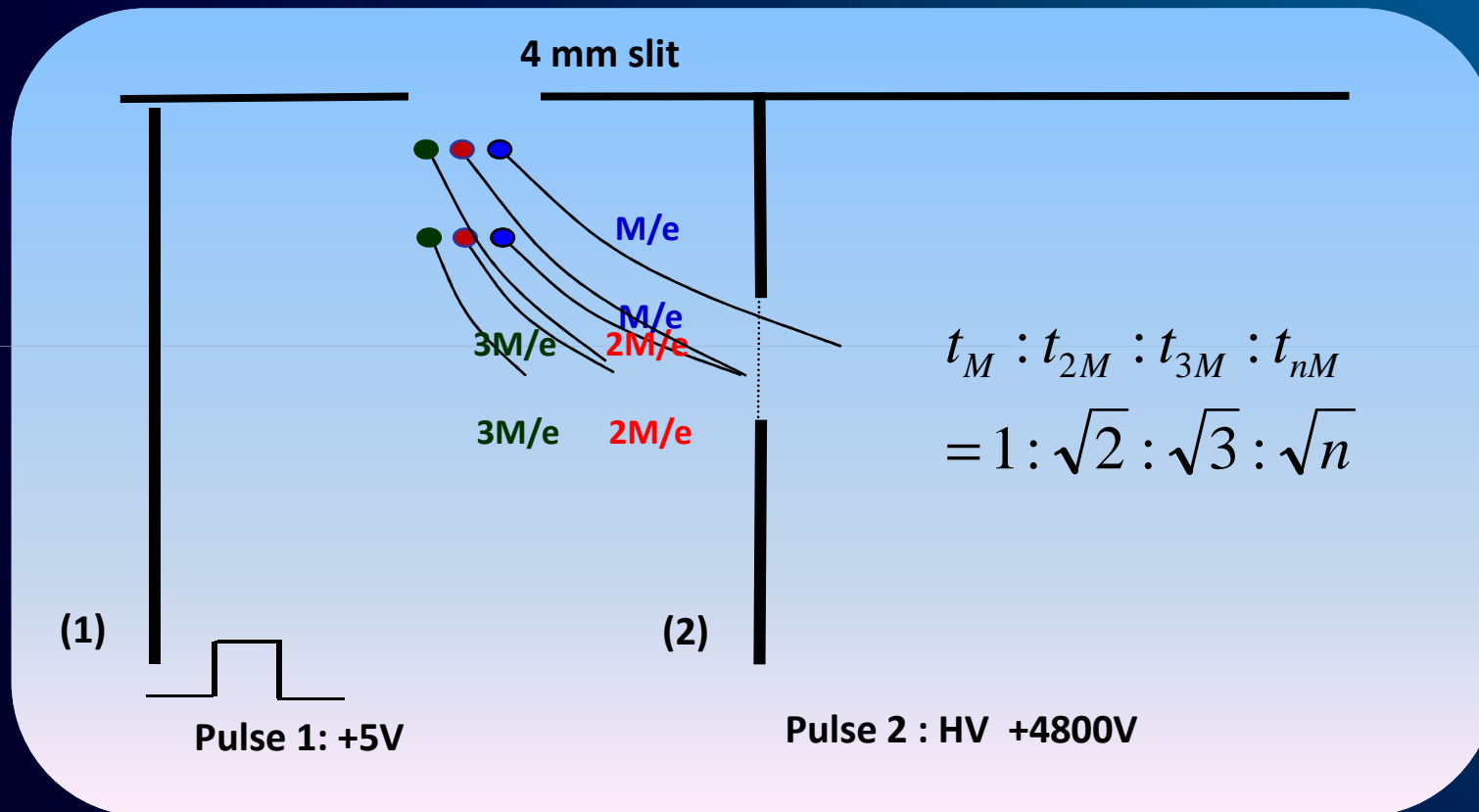
- Cold Plasmas from MOT plus Rydberg excitation
- van der Waals clusters

# Detection scheme

- First, a weak **pulsed** electric field ionises the (neutral?) CRC's into partially charged CRC's
- Second, a strong electric field completely ionises the partially charged CRC's into single cations.
- The time-of-flight of the (heavy) CRC to aperture 2 is used for its mass ( $M/e$ ) determination. Detection is by pulsed field ionisation in the strong field (between apertures 2 and 3) and detection at the time-of flight of the ***single cation*** to detector

# Detection of CRC

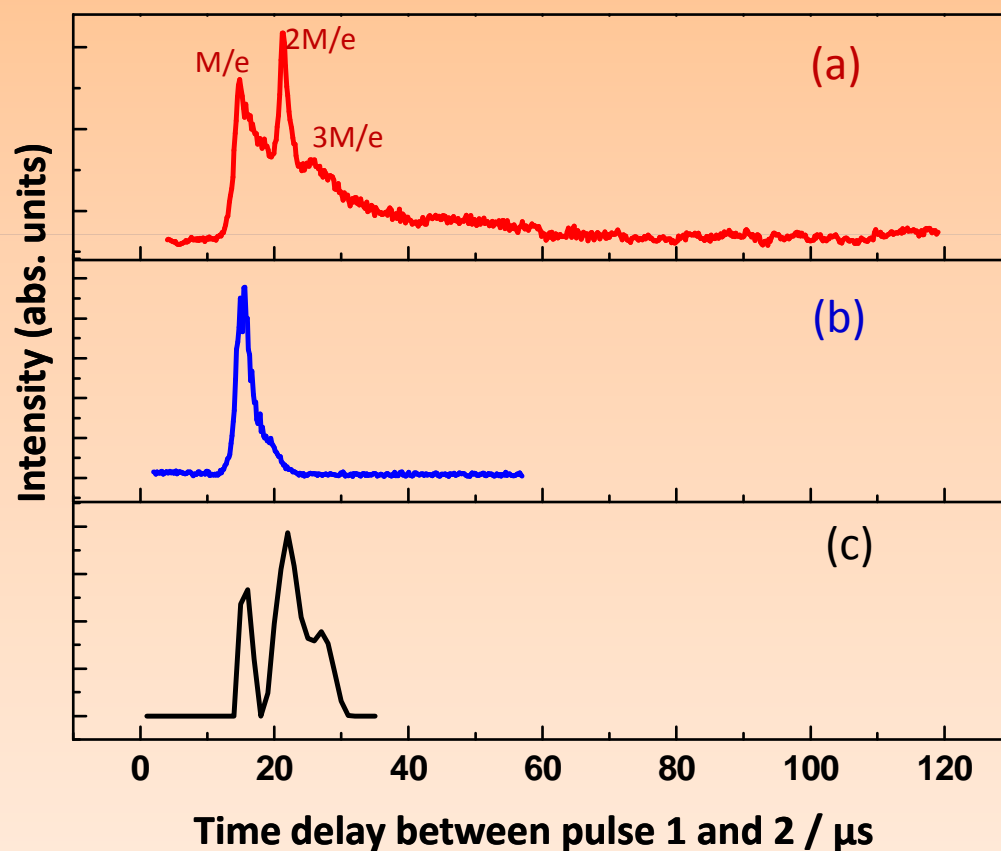
## Method 2: Time Scan of delay between pulse 1 and pulse 2



# Detection of CRC

## Results : Time Scan

Time scan for different positions of crossing the laser and molecular beam



Pulse 1 = 5 V, pulse width=210  $\mu\text{s}$ ; pulse 2 = 4.8 Kv, pulse width = 1.22  $\mu\text{s}$ .

(a) experimental result, laser crosses with molecular beam **2 mm from the nozzle**

(b) experimental result, laser crosses with molecular beam at **15 mm from the nozzle**

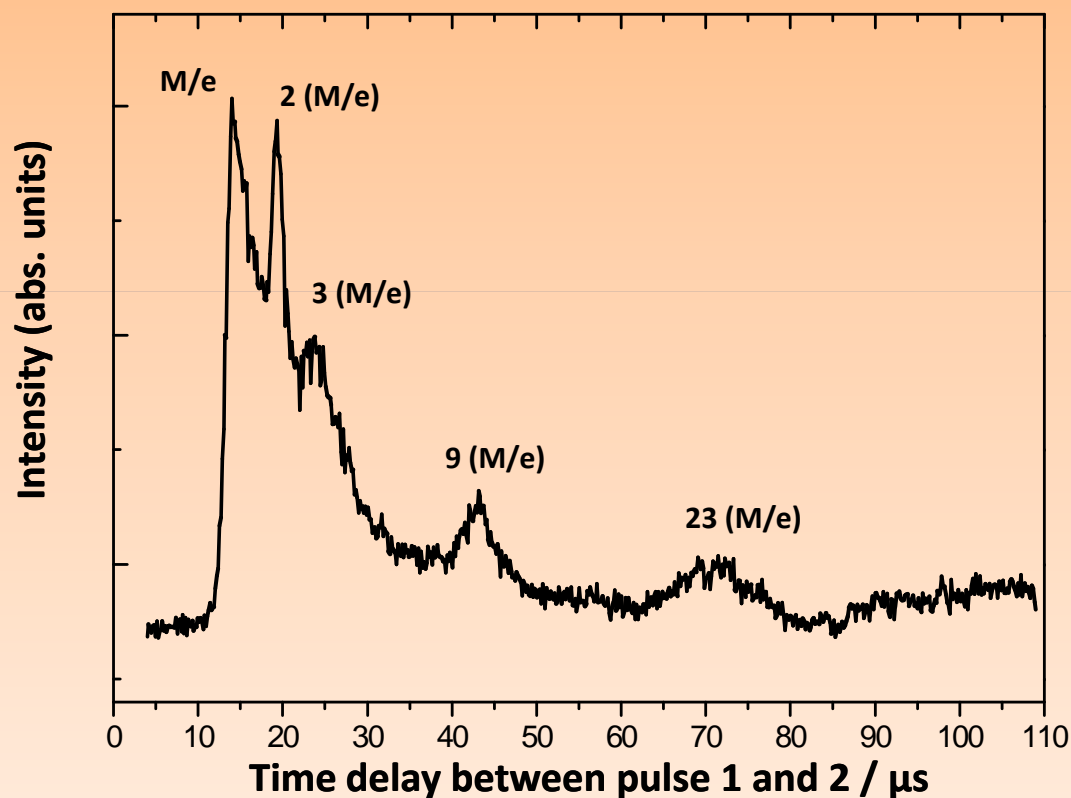
(c) simulations for (a)



# Detection of CRC

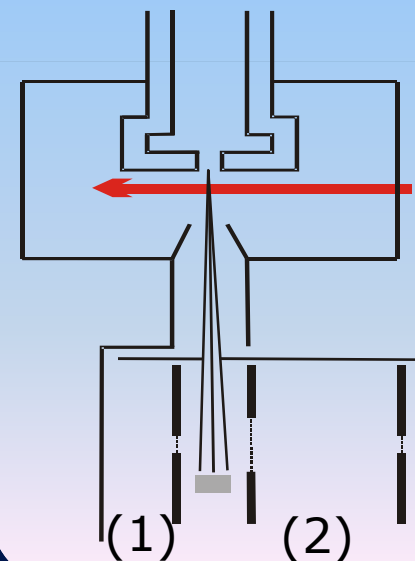
## Results: Looking for High CRC

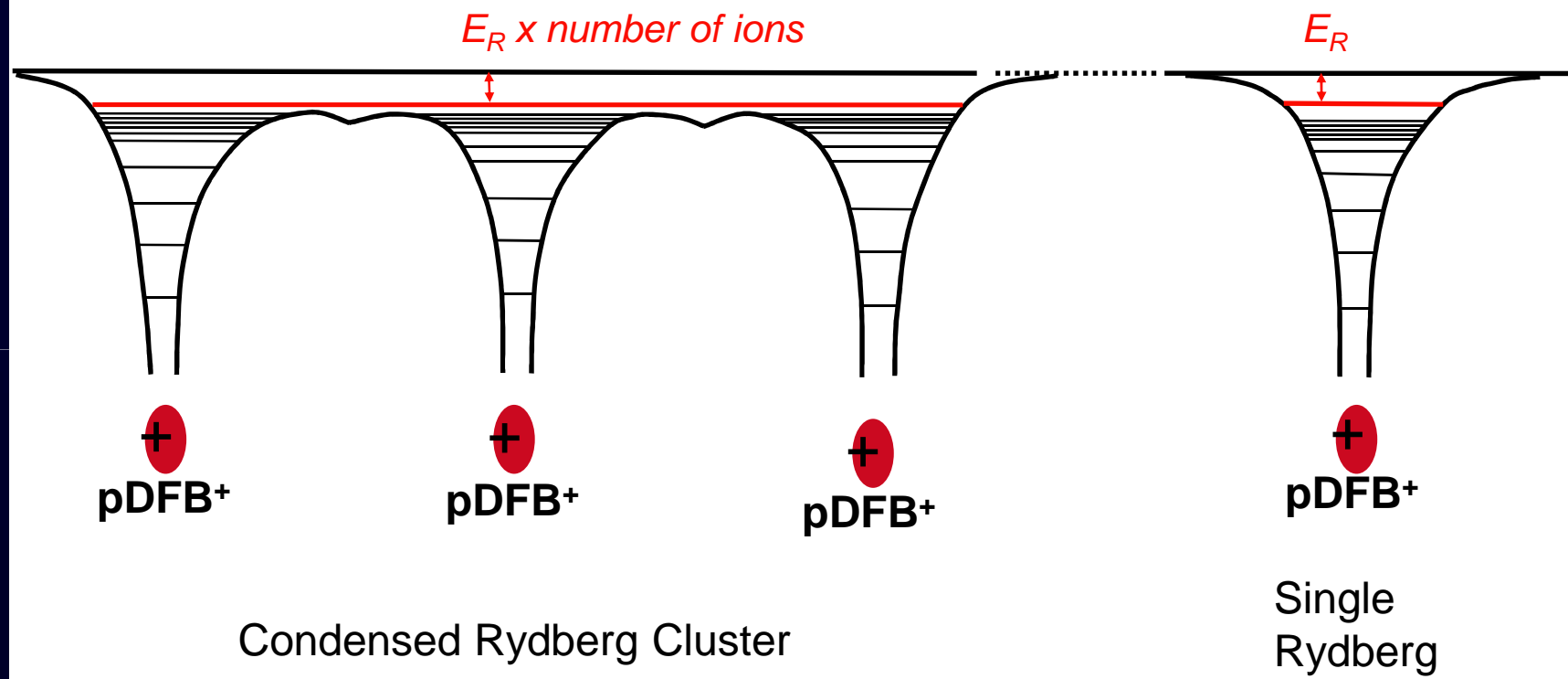
Two-pulse time delay scan

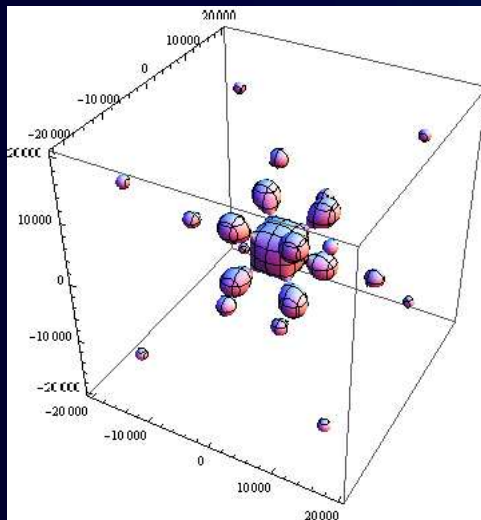


Laser crosses molecular beam at 6 mm from nozzle

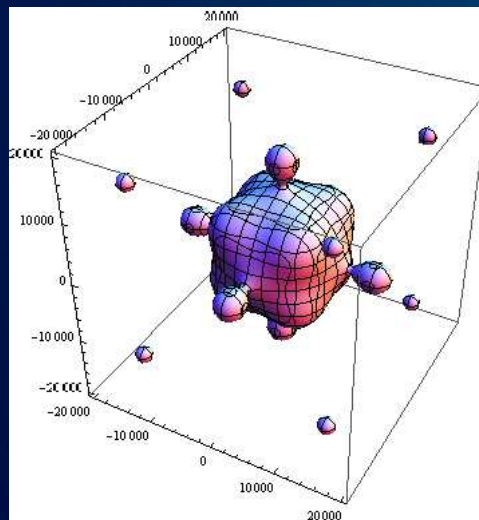
$$\Gamma = \infty \frac{D}{T}$$



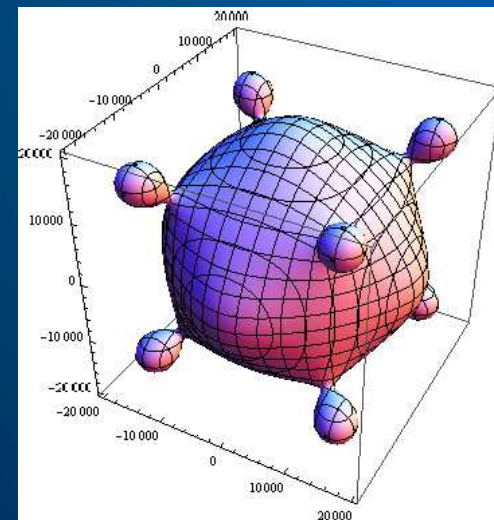




a

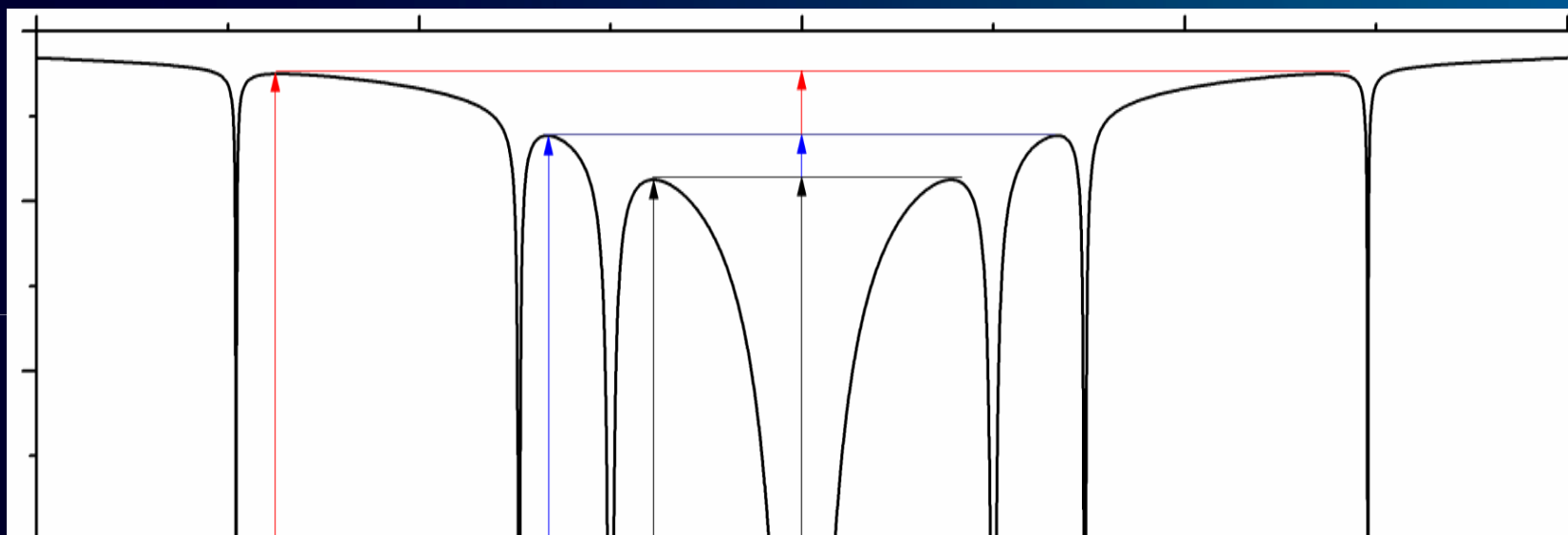


b



c

**Proposed equipotential shell structure of CRCs with  $M/e=23$ , where (a) is the first shell (cube structure); (b) is the second shell (octahedron structure) and (c) is the third shell (cube structure). All dimensions are in atomic units.**





# Conclusions

- Found a novel way to generate an ultracold (near) neutral **strongly coupled** plasma (applicable to molecules)
- Found that the life-time of the plasma is extremely long (milliseconds)
- Found a new state of matter  
**"Condensed Rydberg Clusters"**  
and the condition for its formation.
- Found the mass/charge ratio of **CRC** to be  **$2M/e$ ,  $3M/e$ ,  $9M/e$  and  $23M/e$**  ( $M$  is the mass of pDFB).

- ***What are Condensed Rydberg Clusters?***

Ultracold “condensate” of  
***m ions and n electrons ( $m \approx n$ )*** with  
periodic ion structure and electrons in a  
“conduction band”

- ***Proposed properties of  
Condensed Rydberg Clusters***

1. Under weak electric field conditions CRC will form from break-up of cold Rydberg plasma.

2. Under strong electric field conditions CRC will lose all electrons and fragment into single cations.

# Condition for the formation of CRC

- High density in Rydberg region
- Large ion Coulomb coupling parameter – choose the cross point of laser and molecular beam
- Large electron Coulomb coupling parameter– tuning laser to excite molecules into highly excited Rydberg region: low electron temperature

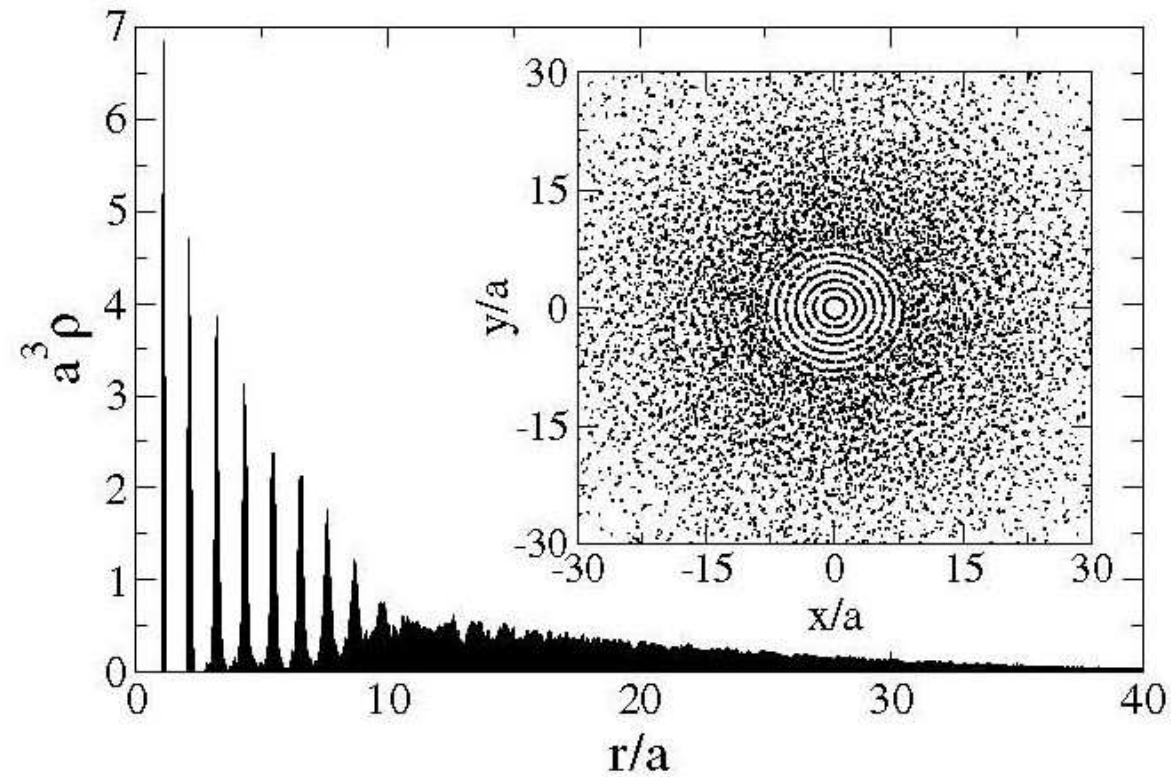


FIG. 3. Radial density after a time of  $t = 24 \mu\text{s}$  ( $\omega_{p0}t = 110$ ). The inset shows a two-dimensional cut through the plasma cloud, clearly revealing the formation of concentric shells. (For better contrast, cuts with  $x = 0$ ,  $y = 0$ , and  $z = 0$ , respectively, have been overlayed.)



# Outlook

## Evaporative cooling in plasma expansion?

ca  $10^{-4}$  of all cations can be described with a Maxwell Boltzmann velocity distribution at very much lower temperature (i.e.  $10\mu\text{K}$ )

*i.e.* the coldest molecular ions concentrate in the centre of the plasma

Starting with  $10^{14}$  ions/cm<sup>3</sup> there would be  $10^{10}$  cations at  $10\mu\text{K}$  in  $1\text{ cm}^3$ , *i.e.*  $10^7$  cations at  $10\mu\text{K}$  in  $1\text{ mm}^3$ ; with the plasma expanding by  $10^3$ , the coldest ions remain in the centre and there would still be  $10^7$  cations in  $1\text{ mm}^3$  in the centre, bound by the electron gas:

This could be useful for BEC of molecules

# Outlook

Choose NO as molecule:

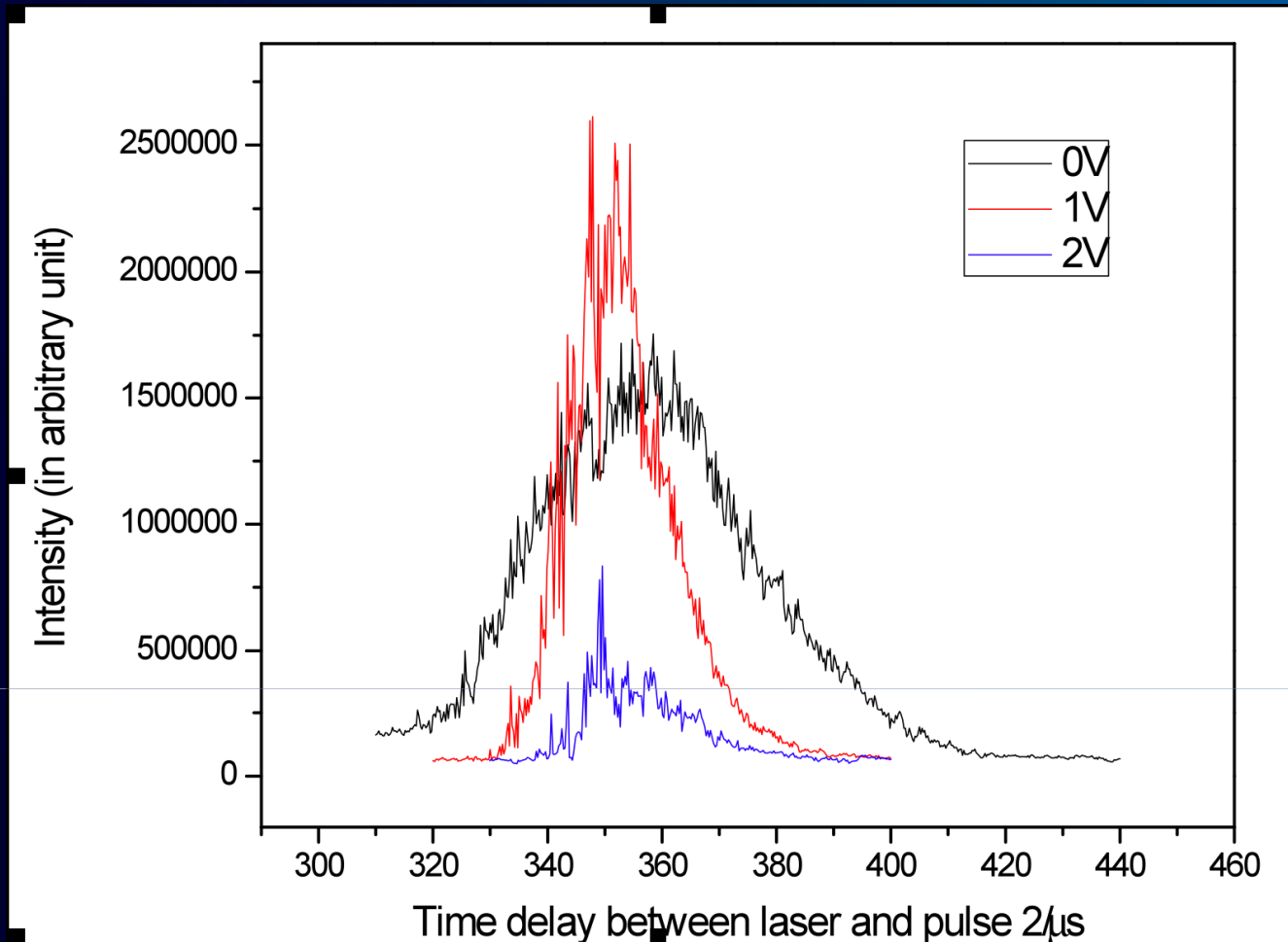
Density of NO<sup>+</sup> can be as high as 10<sup>17</sup> cm<sup>-3</sup>

All molecular cations in the plasma can be **in the same rotational state** by using rotationally selective two-photon, two-colour excitation *via* the A state:

N<sup>+</sup> = 0, 1, 2, etc

How are the rotational angular momenta coupled?  
What about parity?

**NO<sup>+</sup> is a Boson**



Signal from plasma, as a function of delay between laser and pulse 2, for a DC voltage of 0, 1 and 2 V applied to aperture 1 and a 3.6 KV pulse applied to aperture 2

## Co-workers

£££

**E**ngineering and  
**P**hysical  
**S**ciences  
**R**esearch  
**C**ouncil



The Royal Society

**PSI**

The University of Manchester

- **Dr Jingwei Guo**
- **Mervash Varnasseri**
- Dr Xin Tong
- Professor Chris Greene  
(VMP and U Colorado)  
Visiting Miller Professor,  
UC Berkeley



# Proposed Mechanism

- $\text{pDFB} + 2 h\nu \rightarrow \text{pDFB}^* (n > 200)$   
 $\rightarrow \text{pDFB}^+ + e^-$  (no kinetic energy)
- The size of Rydberg orbital is proportional to  $n^2$   
***so Rydberg orbit is several  $\mu\text{m}$***
- Density of pDFB up to  $10^{16}/\text{cm}^3$
- Even if only 0.1% pDFB are Rydberg excited, the average distance between adjacent Rydberg is much smaller than the average Rydberg orbital

# CRC need strongly-coupled plasma environment

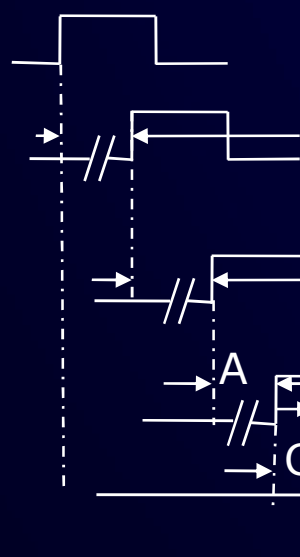
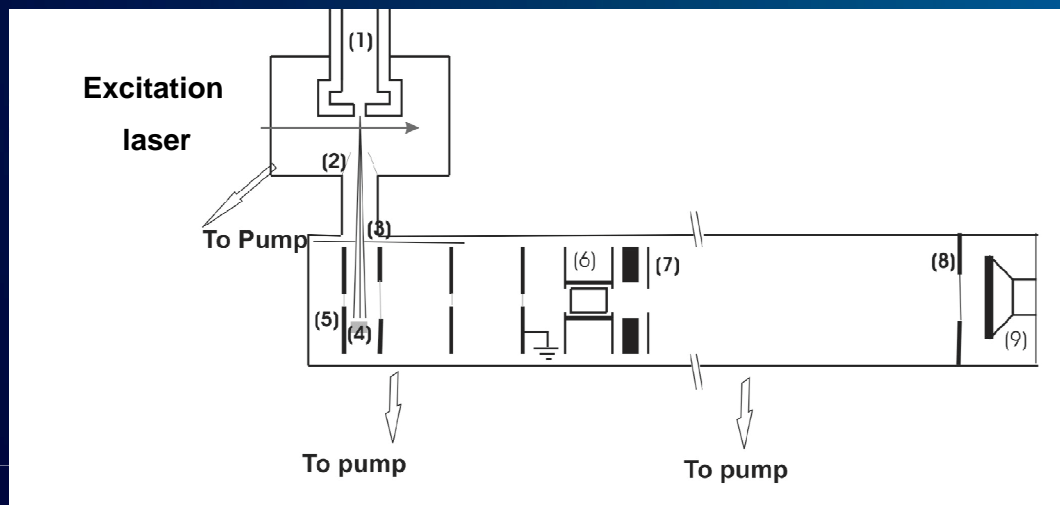
## For the ions

- Free expansion cools down the translational temperature of ions to 0.1 Kelvin
- Density of ions decreases rapidly with distance from the nozzle.
- combining these two factors, Coulomb coupling parameter  $\Gamma$  for ions can be bigger than 100

## For the electrons

- Resonant ionisation of pDFB generates ZEKE photoelectron (electron with very small negative or positive energy *around  $cm^{-1}$* )
- Coulomb coupling parameter for electrons can be bigger than 40.

# Schematic diagram demonstrating the synchronization of the TOF detection of Condensed Rydberg Cluster experiment



Trigger to the general valve controller (time 0 for the whole experiment)

Flashlamp discharge of the YAG laser 200 – 400  $\mu\text{s}$  delay

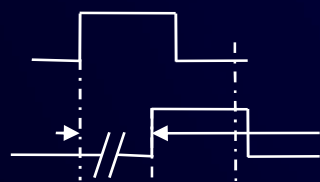
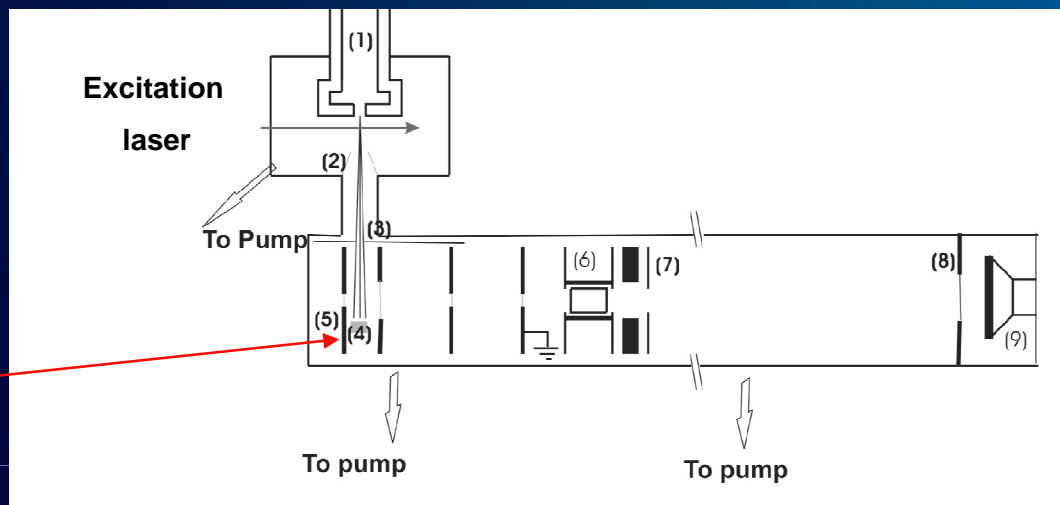
Q-switch of the YAG laser, 300  $\mu\text{s}$  delay

Pulse 1, 450 – 600  $\mu\text{s}$  delay

Pulse 2, 10 – 140  $\mu\text{s}$  delay

# Schematic diagram demonstrating the synchronization of the TOF detection of Condensed Rydberg Cluster experiment

Small DC voltage  
on aperture 1



Trigger to the general valve controller (time 0 for the whole experiment)



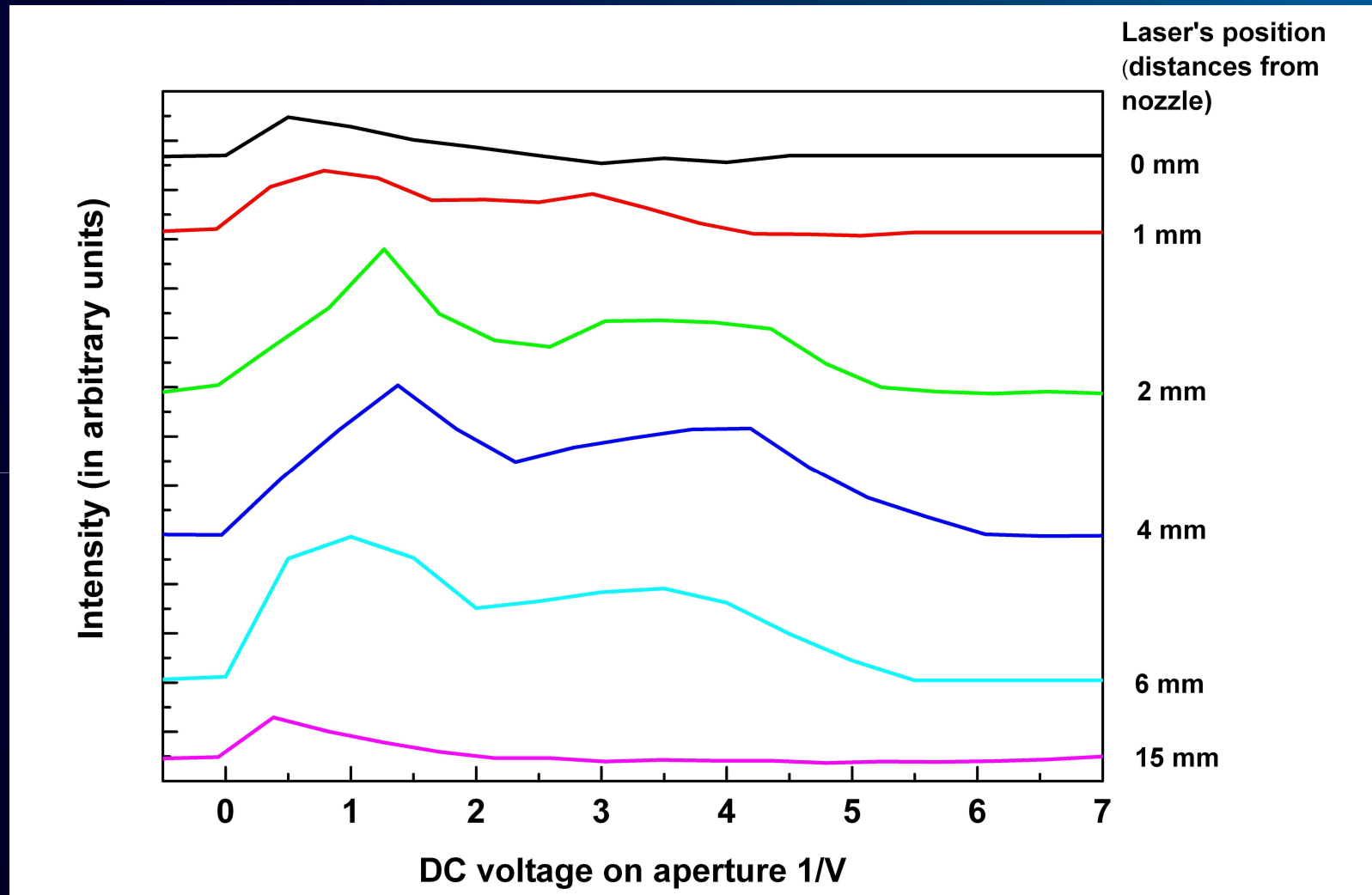
Flashlamp discharge of the YAG laser 200 – 400  $\mu\text{s}$  delay



Q-switch of the YAG laser, 300  $\mu\text{s}$  delay

4.8 kV pulse on aperture 2, 450  $\mu\text{s}$  delay after laser pulse.

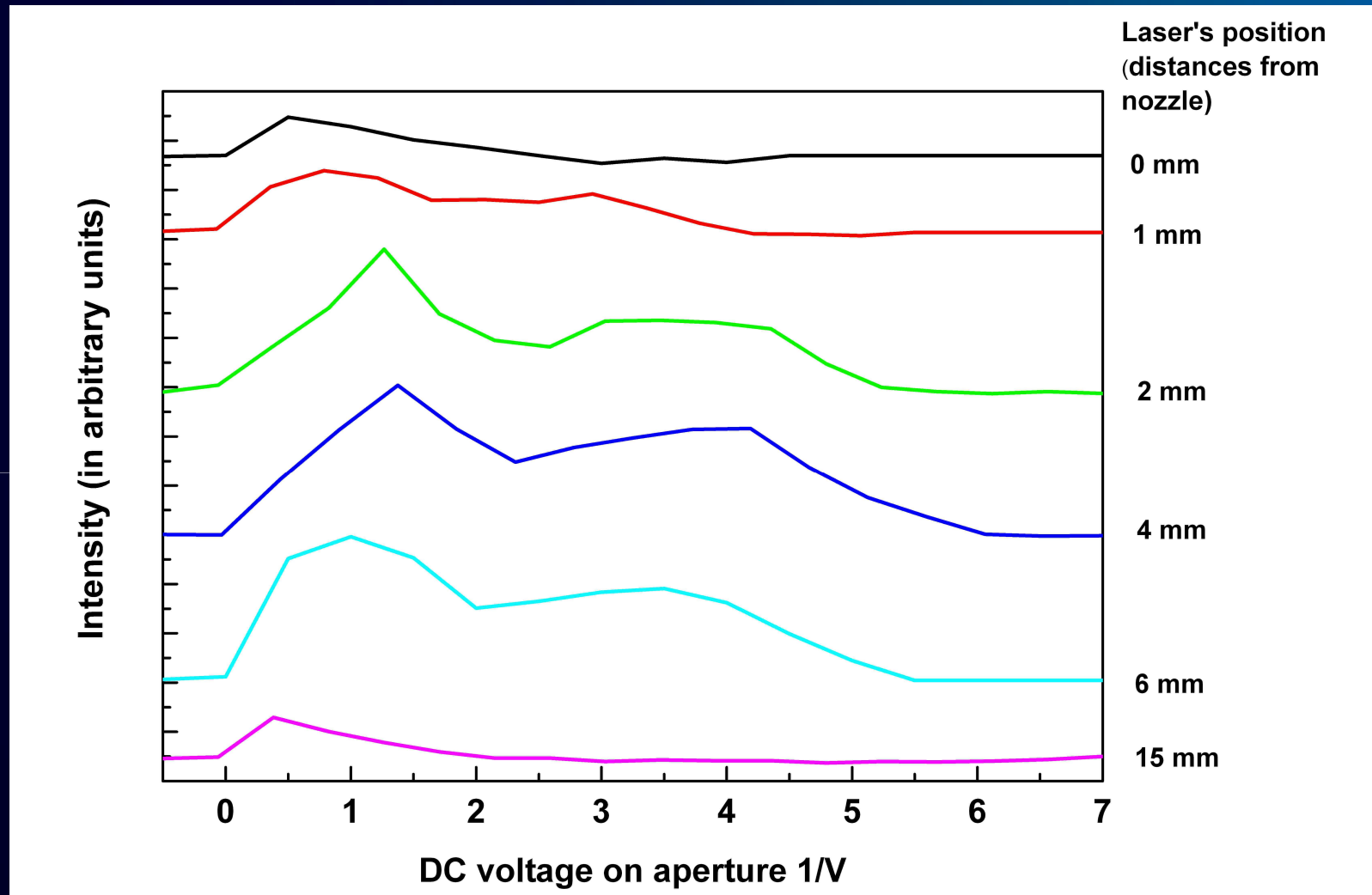
# DC (on aperture 1) voltage scan



+4.8 kV pulse is applied on aperture 2, pulsewidth is 1.22  $\mu$ s. The delay between laser pulse and ionisation pulse on aperture 2 is 450  $\mu$ s

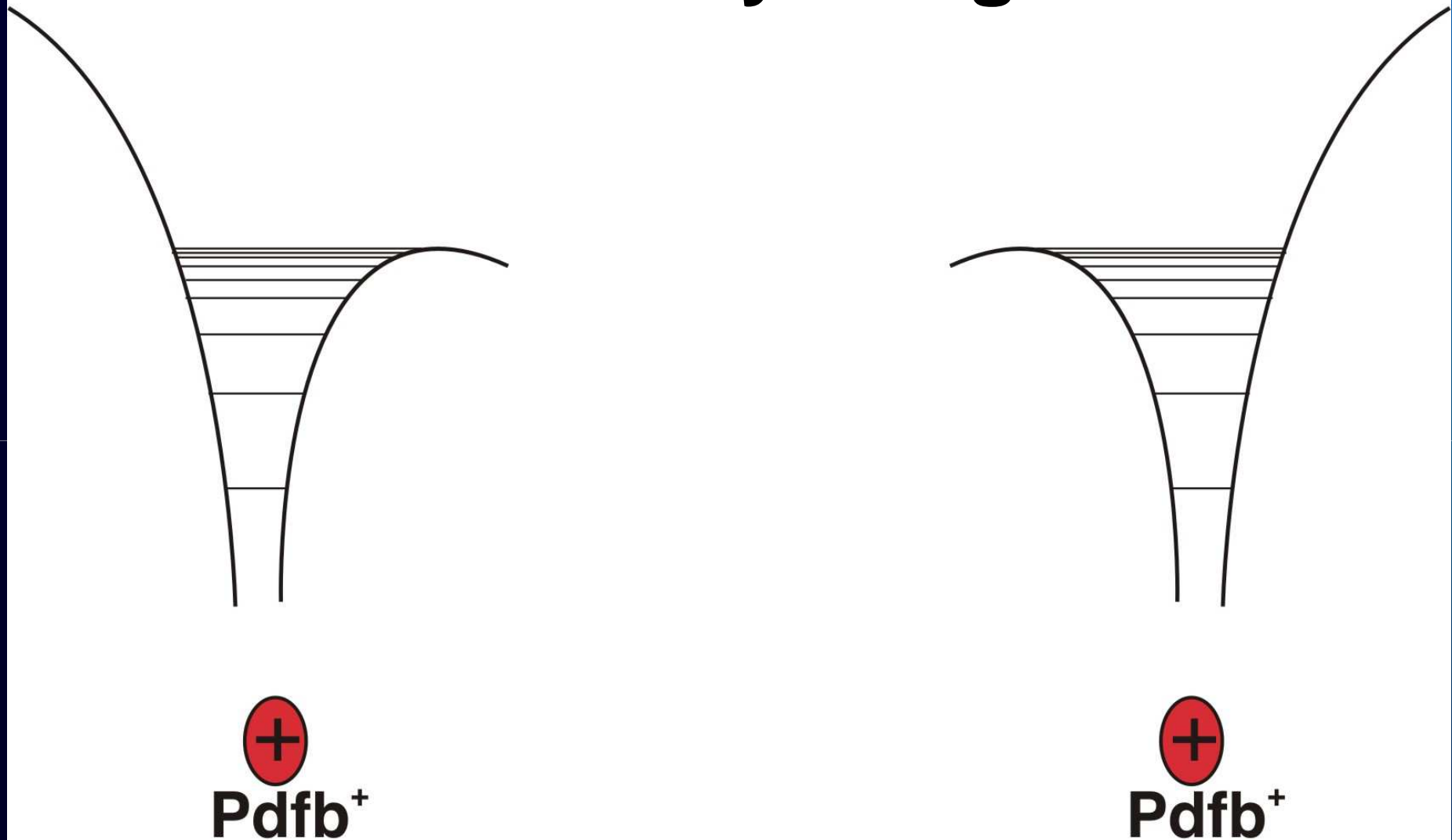


# DC (on aperture 1) voltage scan



4.8 kV pulse is applied on aperture 2, pulsewidth is  $1.22 \mu\text{s}$ . The delay between laser pulse and ionisation pulse on aperture 2 is  $450 \mu\text{s}$

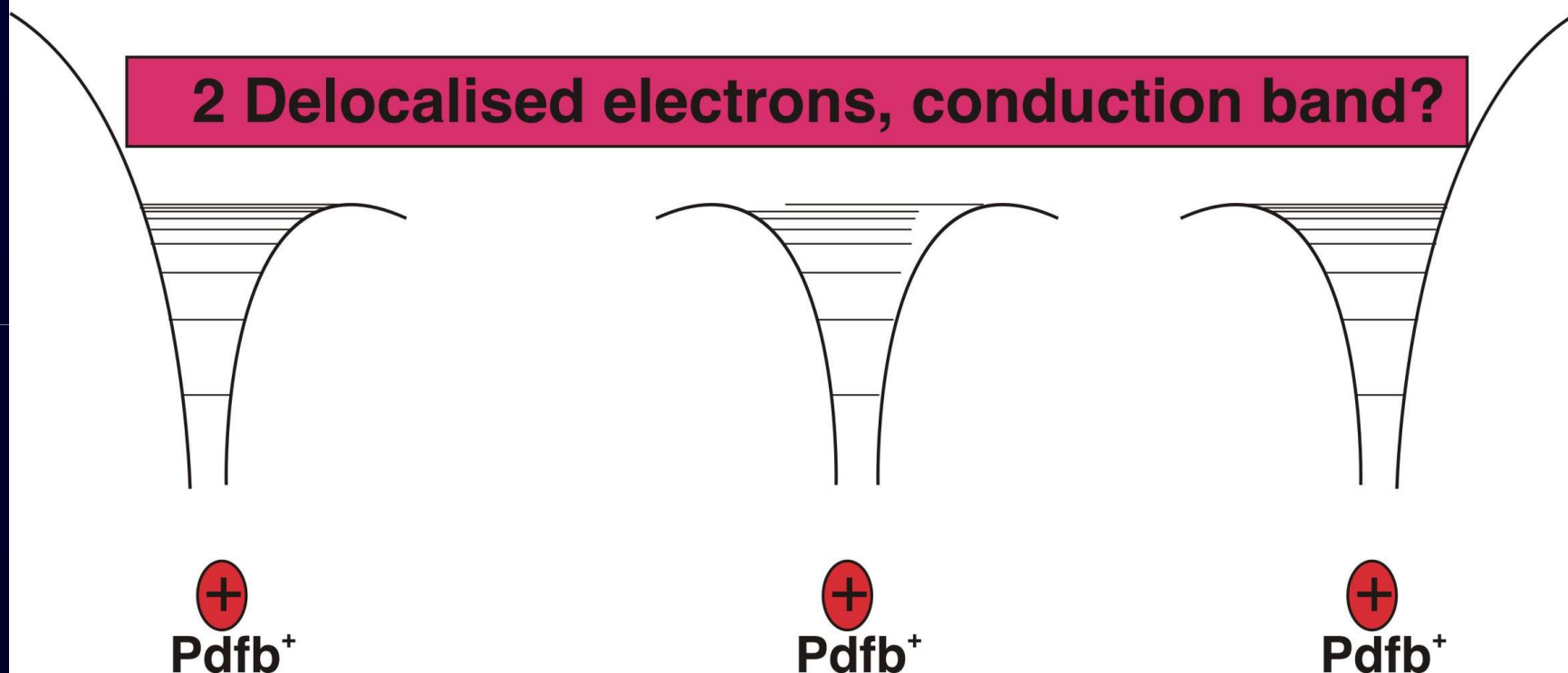
# Condensed Rydberg Dimer



**A few hundred nanometers**

# Condensed Rydberg Cluster Trimer<sup>+</sup>

2 Delocalised electrons, conduction band?



# Plasma expansion

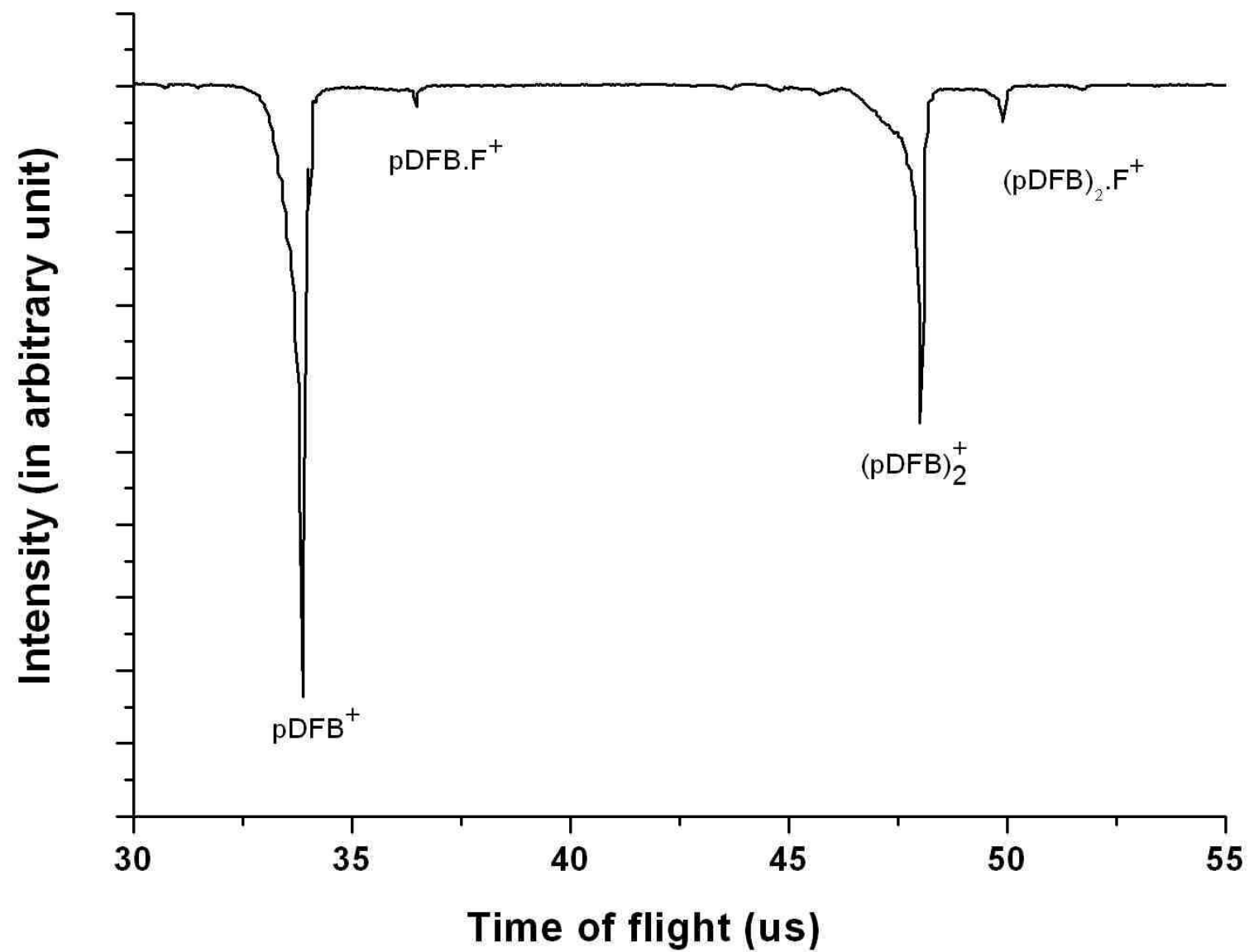
- Electron heat is the driving force of plasma expansion.<sup>1</sup>

$$v_{\text{expansion}} = \sqrt{k_B T_e / m_i}$$

- Electron inelastic collision and three-body recombination is the major process to control the electron temperature.<sup>2</sup>

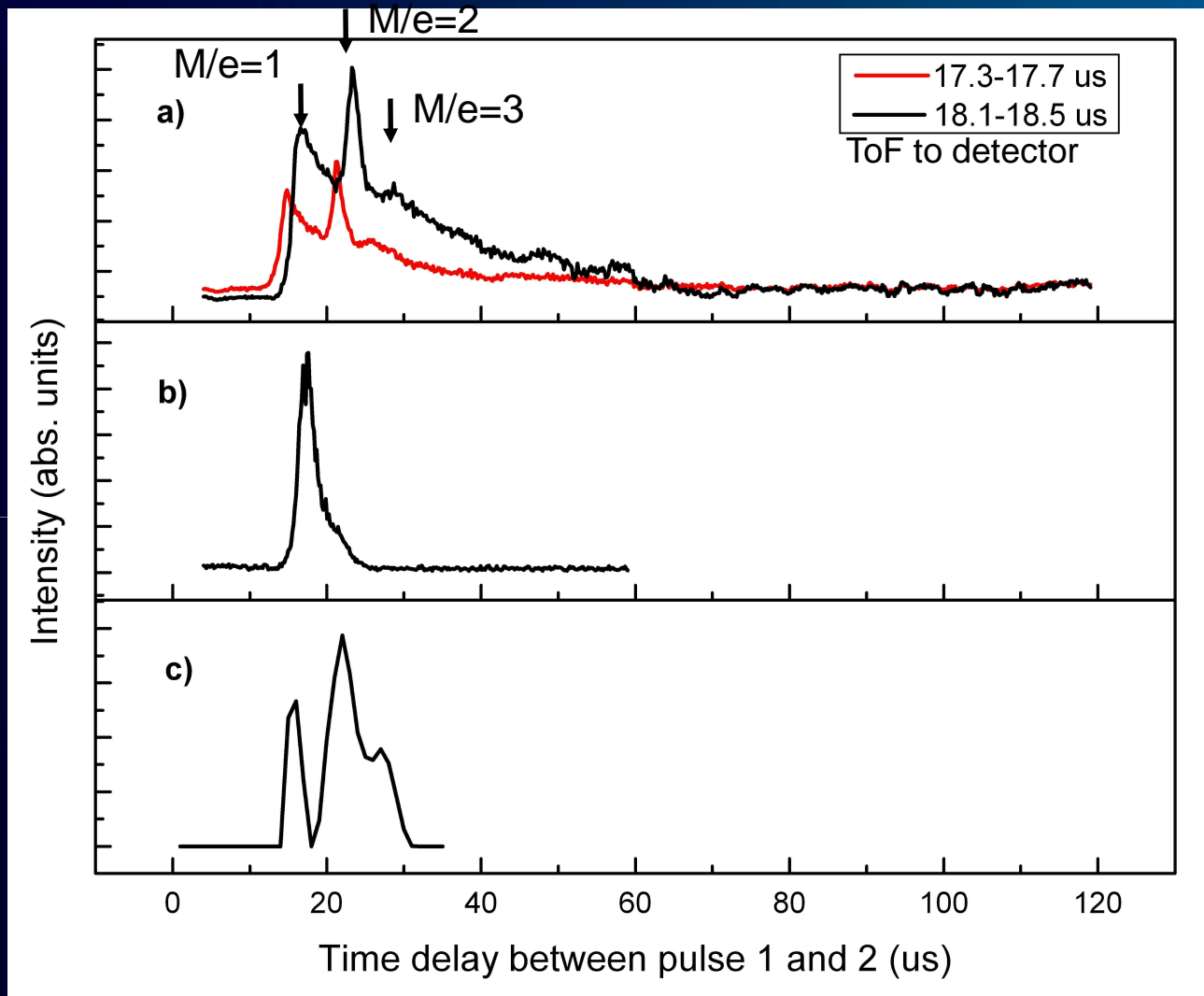
1. R. S. Fletcher, X. L. Zhang and S. L. Rolston, Phys. Rev. Lett., 96 (2006) 105003.

2. P. Mansbach and J. Kech, Phys. Rev. **181** 275 (1969).



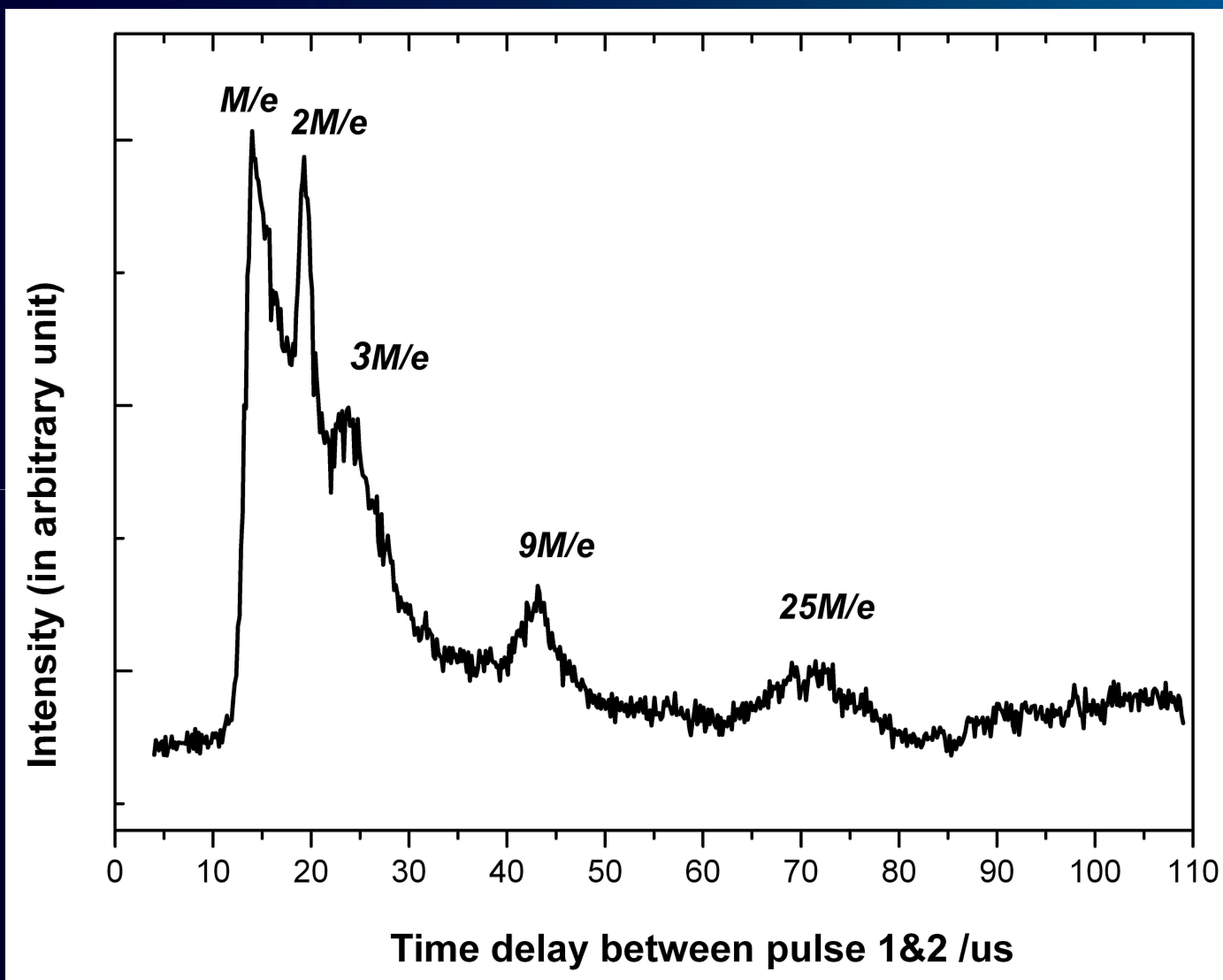


## Two-pulse time delay scan



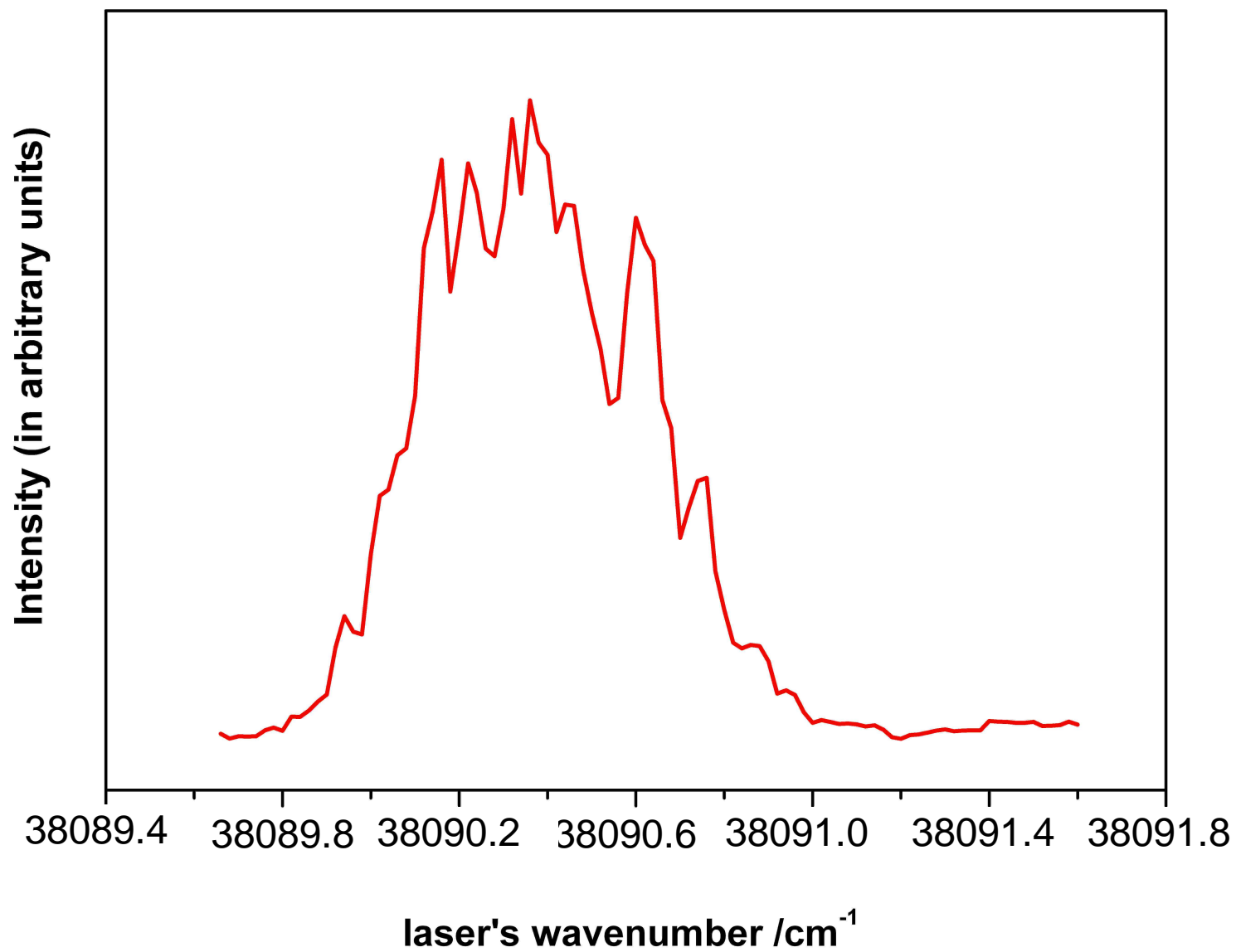
(a), the experimental result (with two different time gates), laser crosses with molecular beam 2 mm from the nozzle; (b) the experimental result, laser crosses with molecular beam at 15 mm from the nozzle; (c), simulations for (a).

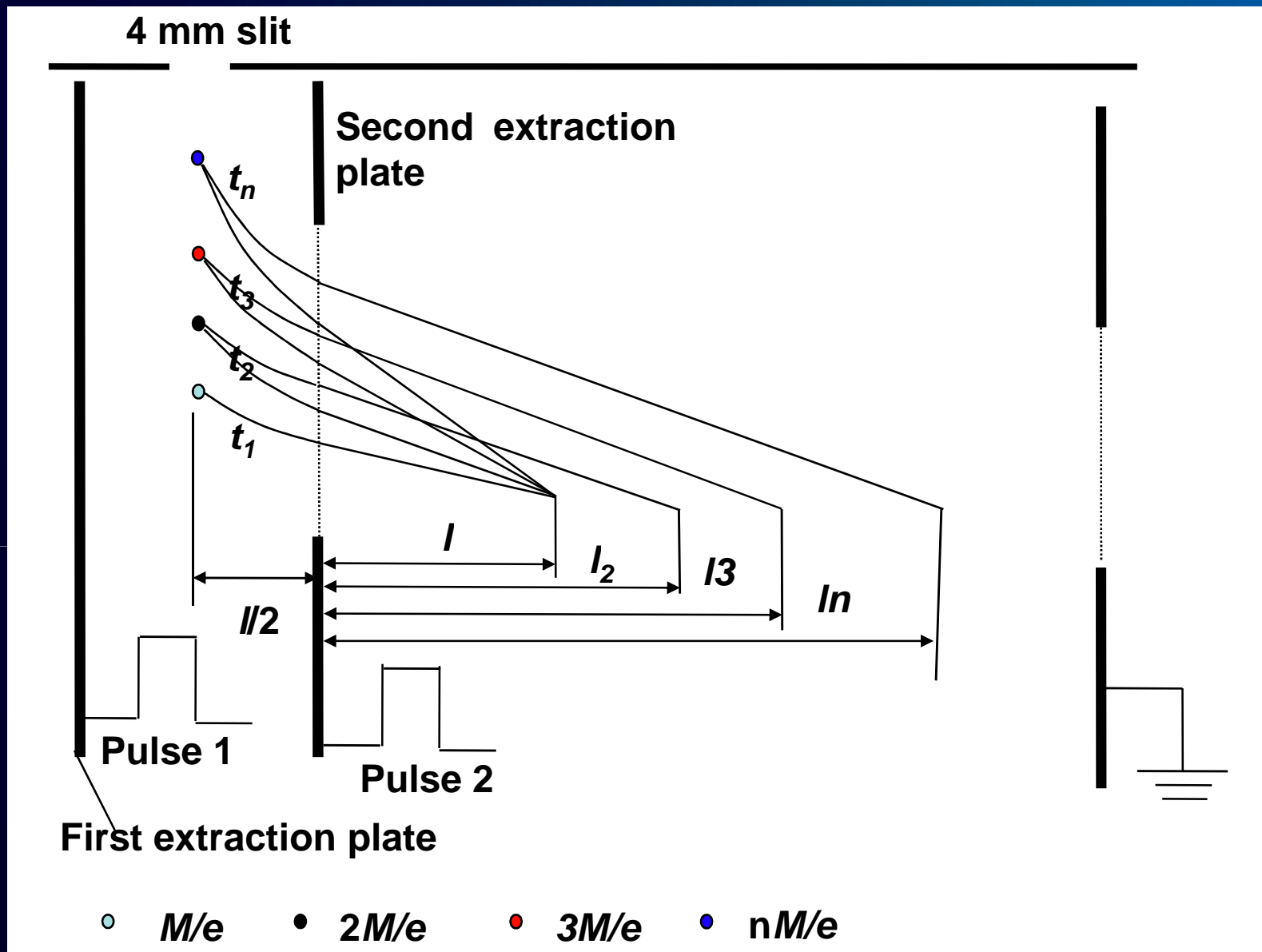
# Two-pulse time delay scan



Laser crosses molecular beam at 6 mm from nozzle

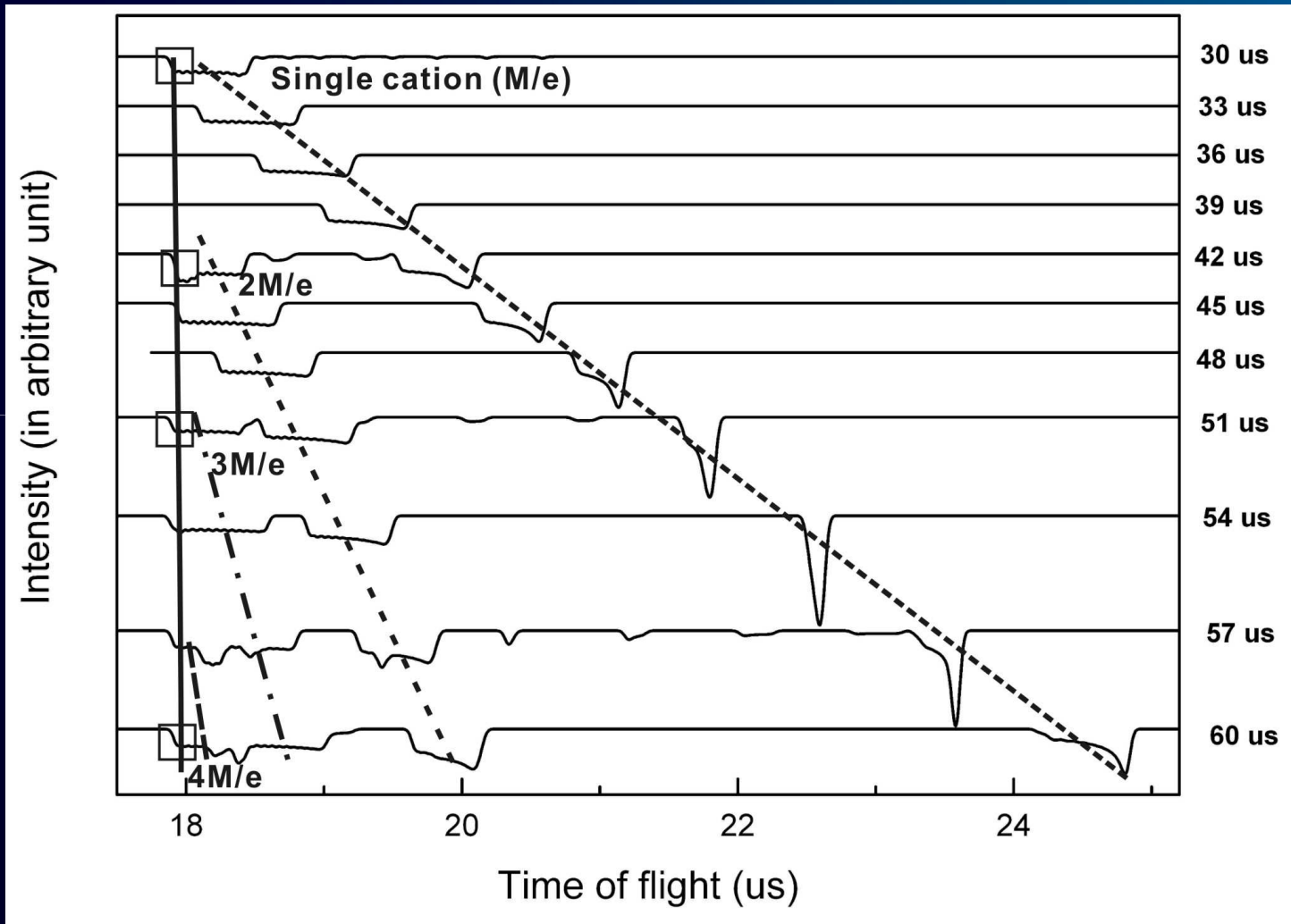
Spectral scan for signal at  $t_2 - t_1 = 70$  s





$$t_M : t_{2M} : t_{3M} : t_{nM} = 1 : \sqrt{2} : \sqrt{3} : \sqrt{n}$$

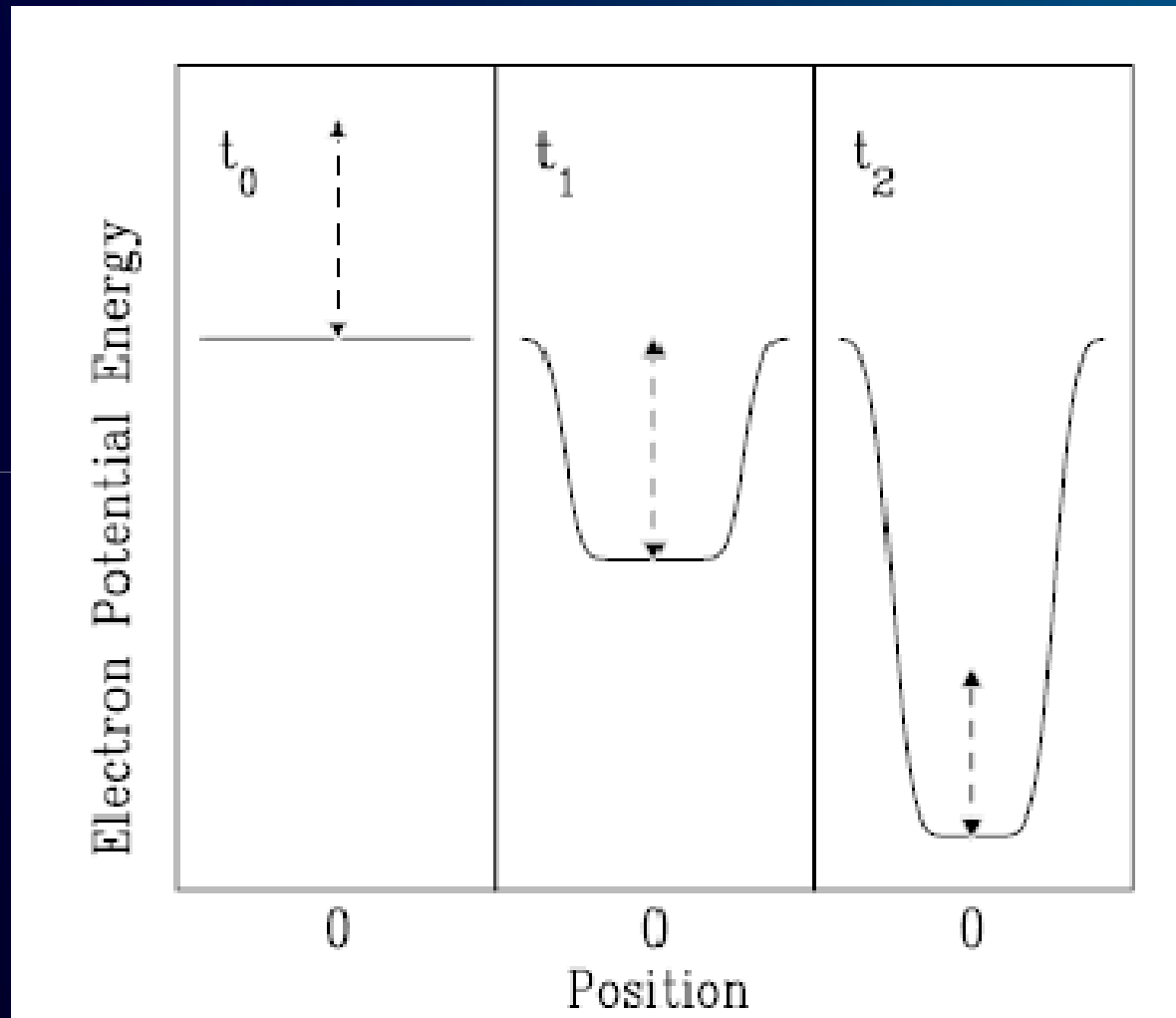
Simulation of the TOF spectra of pDFB<sup>+</sup>, with different delays between pulses 1 and 2 (labelled on the right hand side of each curve). Pulse 1 = 1.25 V, pulse 2 = 4.8 kV, both pulses are full energy extraction.



$$30: 42: 51: 60 = 1 : 1.4 : 1.7 : 2 = t_M : t_{2M} : t_{3M} : t_{4M}$$



# Plasma potential well evolution



T. C. Killian, S. Kulin, S. D. Bergeson, L. A. Orozco, C. Orzel and S. L. Rolston, Phys. Rev. Lett., 83 (1999) 4776.

# **Excitation to Rydberg states in an ultracold cloud of atoms**

**Laser excitation using pulsed dye laser**

**A substantial part of the atoms is laser excited into Rydberg states (typically  $n$  around 80, but newer experiments also excite to higher  $n$  or above threshold).**

# Electron temperature control

- Three-body-recombination heat up electron  $\sim T(-9/2)$ .
- Deexcitation of Rydberg molecules:

$$A_d = 7.2 \rho_e (27.2 \text{ eV} / K_B T_e)^{0.17} v^{2.66} a_0 \alpha c$$

- Excitation of Rydberg including

$$A_e = 55 \rho_e (K_B T_e / 27.2 \text{ eV})^{0.83} v^{4.66} a_0 \alpha c$$

$$v = \sqrt{-13.6 \text{ eV} / E_R}$$

Bohr length

the fine-structure constant

$$A_d / A_e = \frac{0.13 E_R}{-13.6 \text{ eV}} \left( \frac{27.2 \text{ eV}}{K_B T_e} \right)^{\frac{4.66}{2.66}} = - \frac{0.26 E_R}{K_B T_e}$$

# Definition of CRC

Condensed Rydberg Clusters (CRC) can be defined as dense clusters of highly-excited Rydberg molecules. The average distance between adjacent ionic cores is much smaller than the average Rydberg orbit, so that every Rydberg electron interacts with many cations. \*

Jingwei Guo, Klaus Muller-Dethlefs, in preparation to submit to *Science*