

# Bending Back Light: The Science of Negative Index Materials

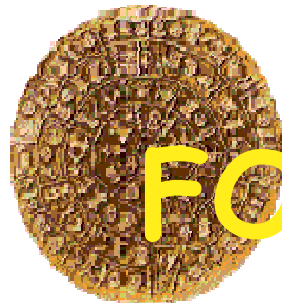
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Ames Lab. and Physics Dept. Iowa State  
University, Ames, Iowa, USA

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University of Crete, Heraklion, Crete, Greece



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IOWA STATE UNIVERSITY  
OF SCIENCE AND TECHNOLOGY

# Left-Handed Materials



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## History:

- Permittivity  $\epsilon$ , permeability  $\mu$  and index of refraction  $n$  negative
- Reversal of Snell's Law, perfect focusing, flat lenses, etc.
- Impedance match  $z = \sqrt{\mu/\epsilon}$  and  $n = -1$ , zero reflection
- $\lambda \gg a$  in LHM, while  $\lambda = a$  in PBG

Both PBG and LHM exhibit properties not found in naturally materials

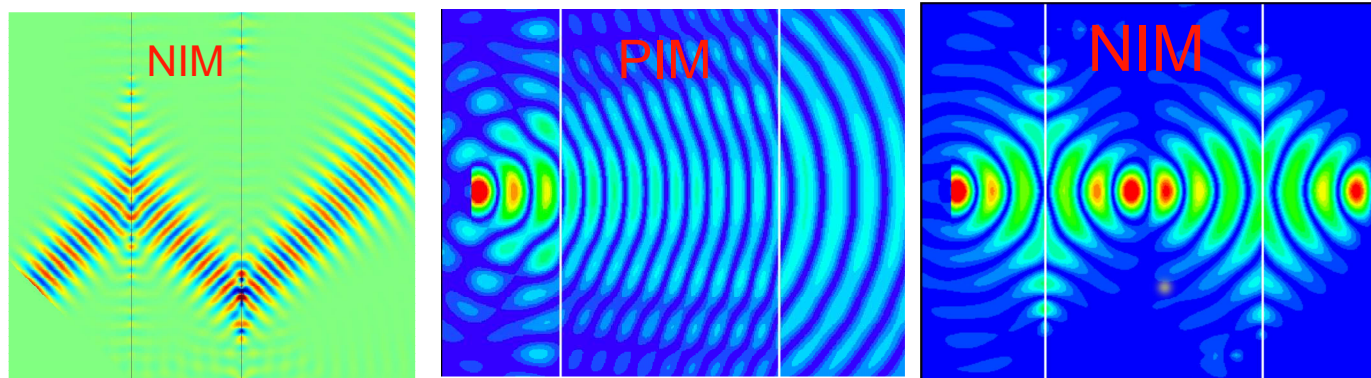
## Vision:

- Understanding the physics and the exotic properties of LHMs
- Perfect Lens. Near-field optical microscopy, nano-lithography
- Wireless and optical communications. RF sensing.
- Antenna and microwave device miniaturization

- 
- ✓ Breakthroughs and new concepts in materials processing at nanoscale
  - ✓ Search for new materials that exhibit  $n < 0$  at THz or optical regime

# Some reviews articles from our group

- 1) Bending Back Light: The Science of Negative Index Materials  
*Optics and Photonics News*, June 2006
- 2) Negative index materials: New frontiers in optics,  
*Adv. Mater.* **18**, 1941 (2006)
- 3) Photonic metamaterials: Magnetism at optical frequencies,  
*IEEE J. of Selected Topics in Quant. Electr.* **12**, 1097 (2006)
- 4) Negative Refractive Index at Optical Wavelengths  
*Science* **315**, 47 (2007)



# Collaborators

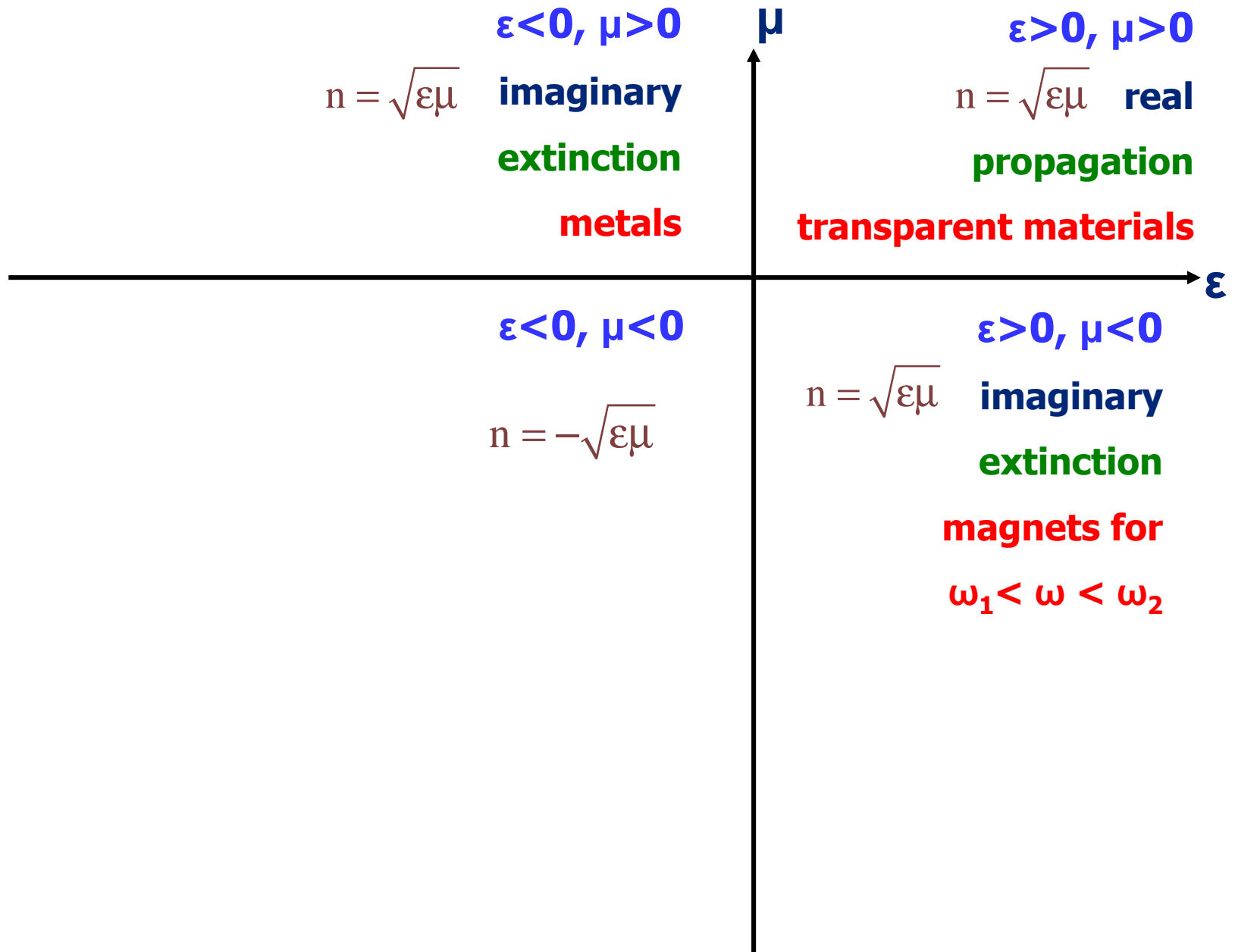
P. Markos (Ames & Slovakia), Th. Koschny (Crete & Ames)  
E. N. Economou, M. Kafesaki, N. Katsarakis (Crete, Greece)  
S. Foteinopoulou (Crete & Ames)  
E. Ozbay's group (Bilkent, Turkey)  
Lei Zhang, J. Zhou, R. Moussa & G. Tuttle (Ames, USA)  
D. R. Smith (Duke, USA); J. B. Pendry (Imperial, UK)  
V. Sandoghdar (ETH, Zurich)  
M. Wegener's group, S. Linden (Karlsruhe, Germany)  
Boeing's group (Seattle, USA)

<http://cmpweb.ameslab.gov/personnel/soukoulis>

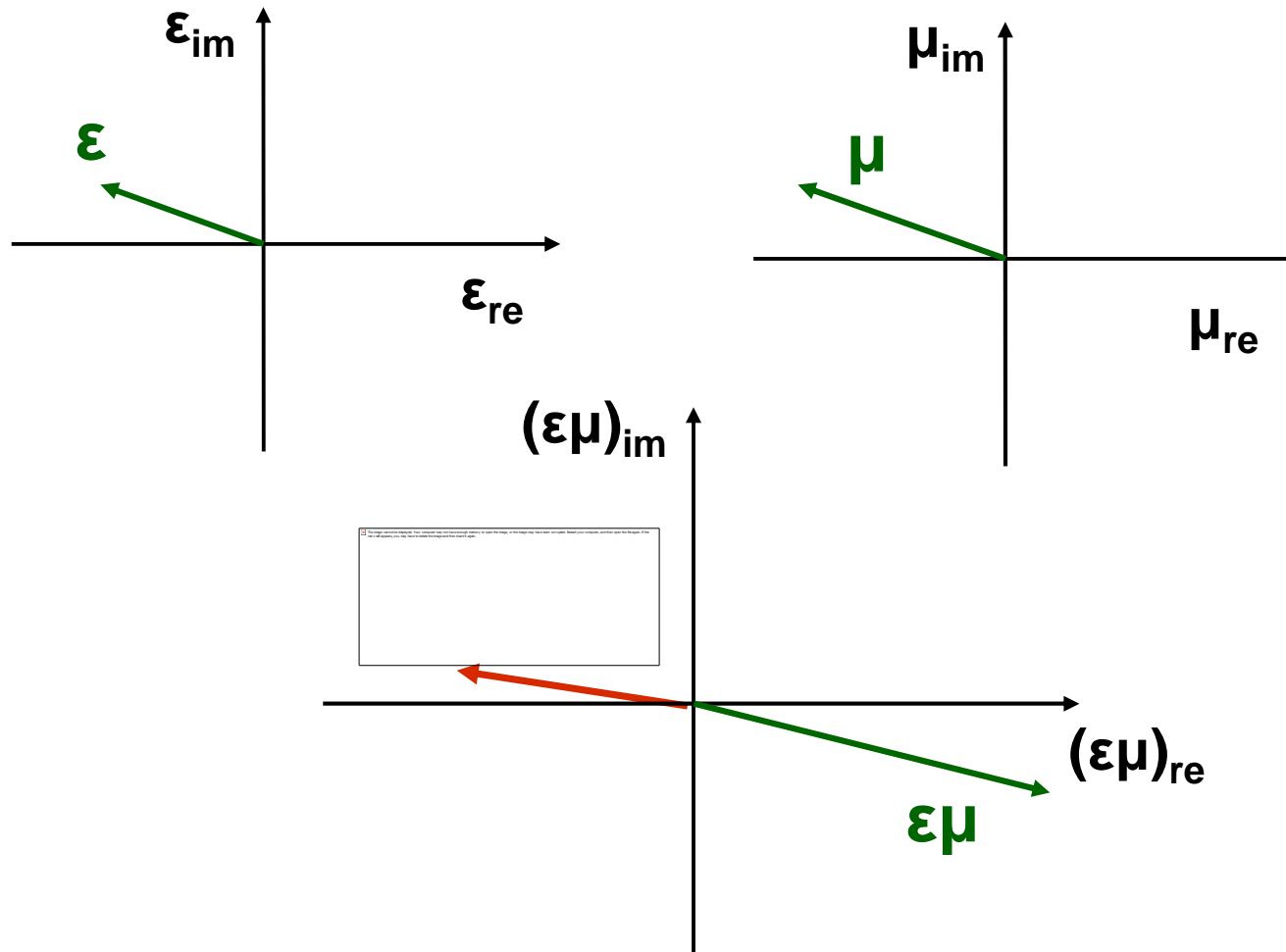


# Outline of Talk

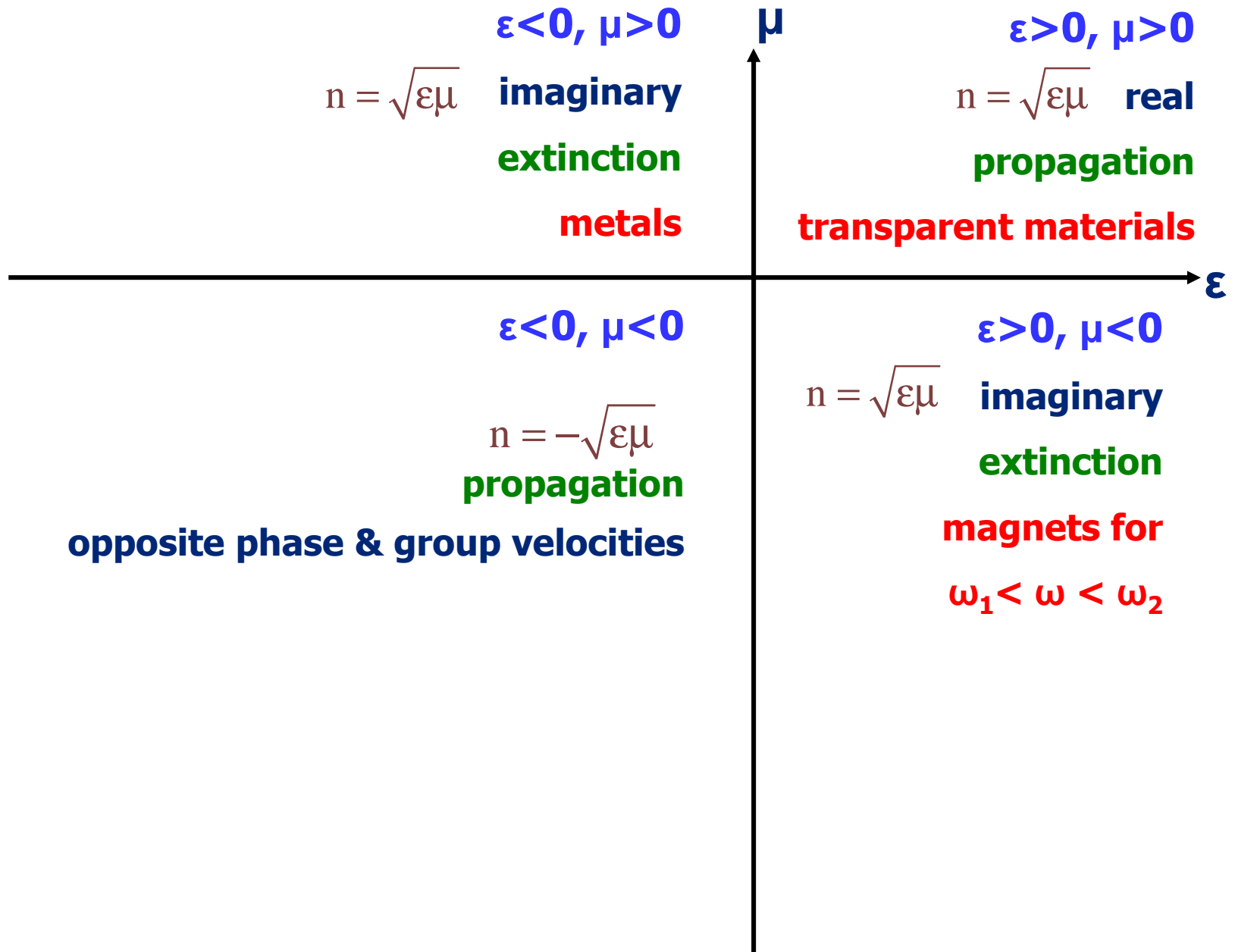
- **Brief history of left-handed materials**
- Electric and Magnetic Response of SRRs and LHMs
- 6, 100 and 200 THz response of SRRs (Karlsruhe, Crete, Ames)
- Upper limits of the SRRs? Simulation results and their interpretation by a LC model. Experiments.
- Breaking of scaling. Top-down approach does not work.
- Diamagnetic response and current density.
- LHM by Double Layer Cut - Wires. Negative  $n$  at optical frequencies.
- **Negative group** and phase **velocities** in NIMs!
- No negative  $n$  with only cut wires.
- Losses can give a negative  $n$ , without LH propagation.
- Concluding Remarks (EIT, Chiral, Losses, 3d structures with DLW)



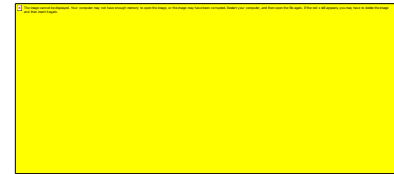
## Negative $\epsilon$ and $\mu$ lead to negative $n$



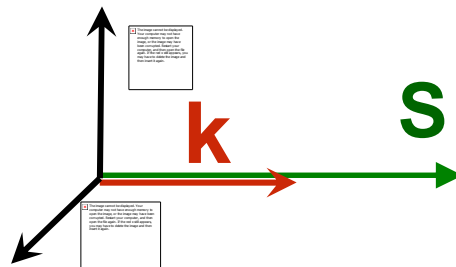
Causality requires  $\text{Im}(n) > 0 \Rightarrow \text{Re}(n) < 0$



# Group vs phase velocity

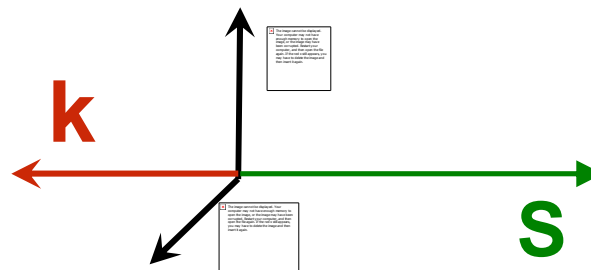


➤ For  $\epsilon > 0$ ,  $\mu > 0 \Rightarrow (\mathbf{k}, \mathbf{E}, \mathbf{H})$  right-handed set



$$\mathbf{S} = \frac{c}{4\pi} \mathbf{E} \times \mathbf{H} = \langle u \rangle \mathbf{v}_g$$

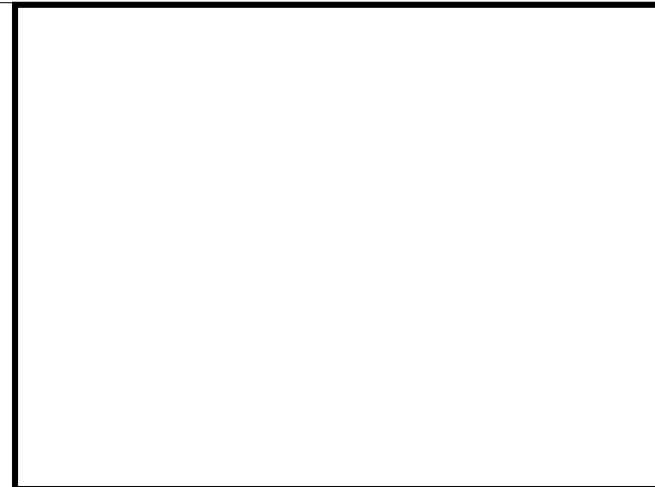
➤ For  $\epsilon < 0$ ,  $\mu < 0 \Rightarrow (\mathbf{k}, \mathbf{E}, \mathbf{H})$  left-handed set



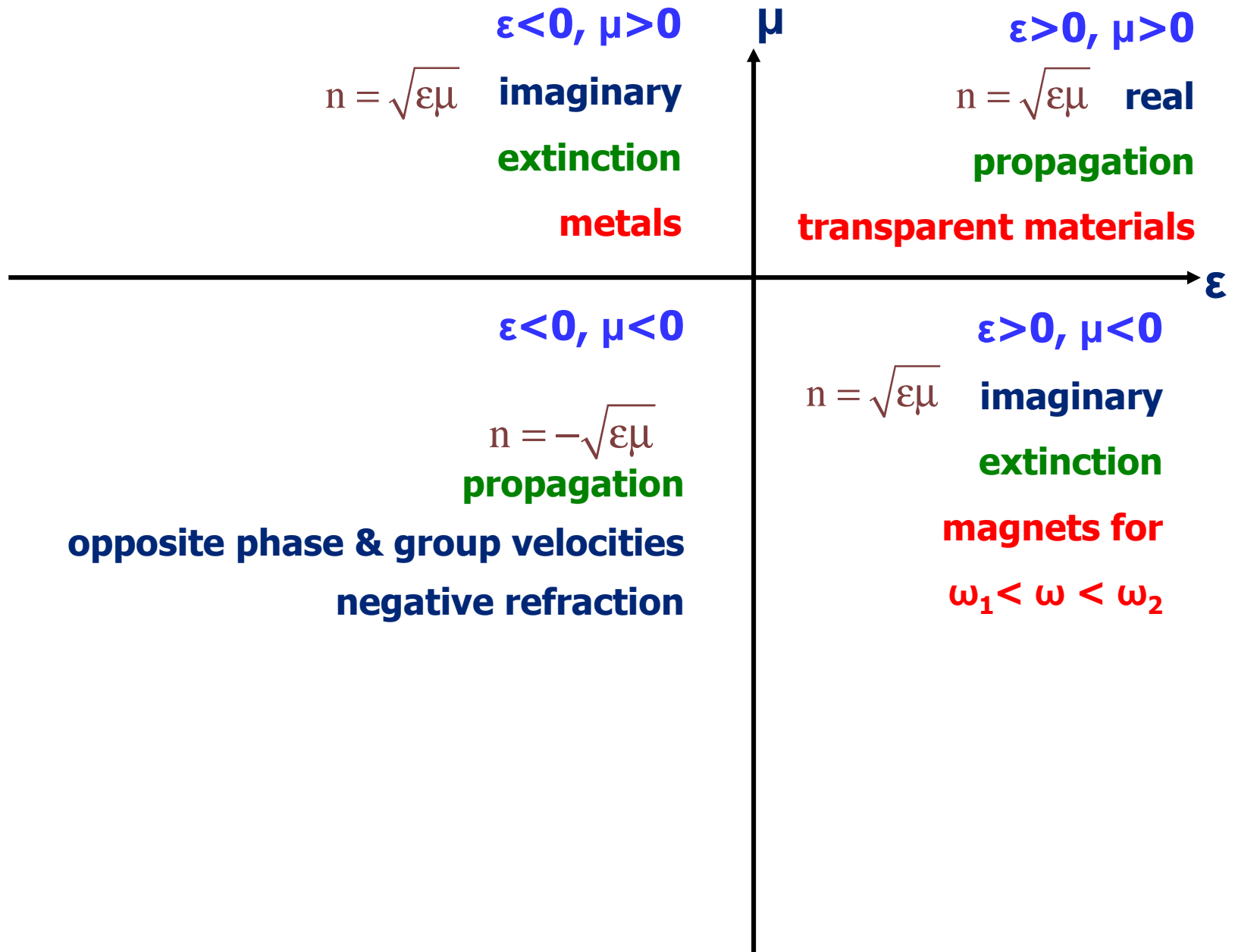
$\mathbf{v}_g \updownarrow \mathbf{v}_{ph}$

# Opposite phase and group velocities !!

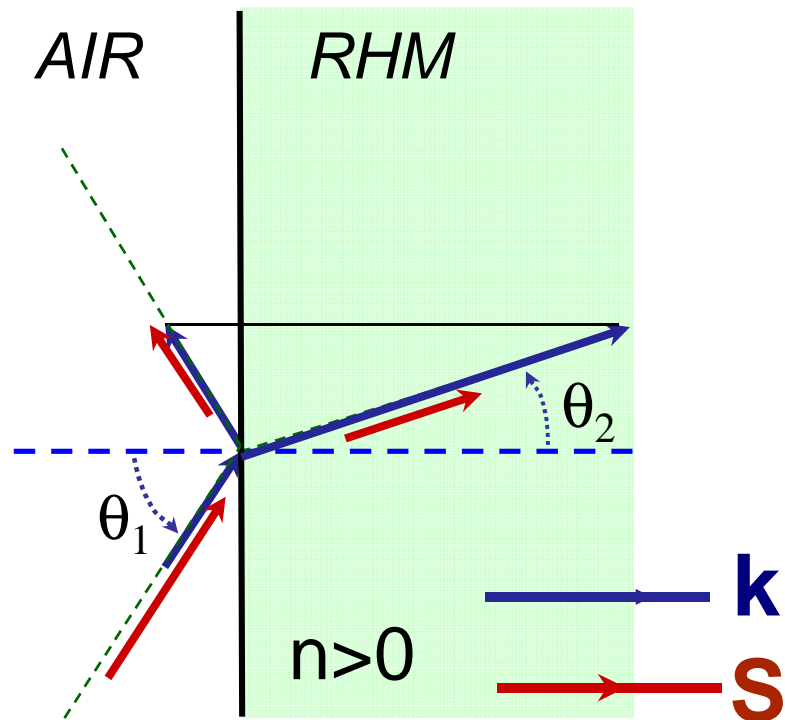
↑  
Source



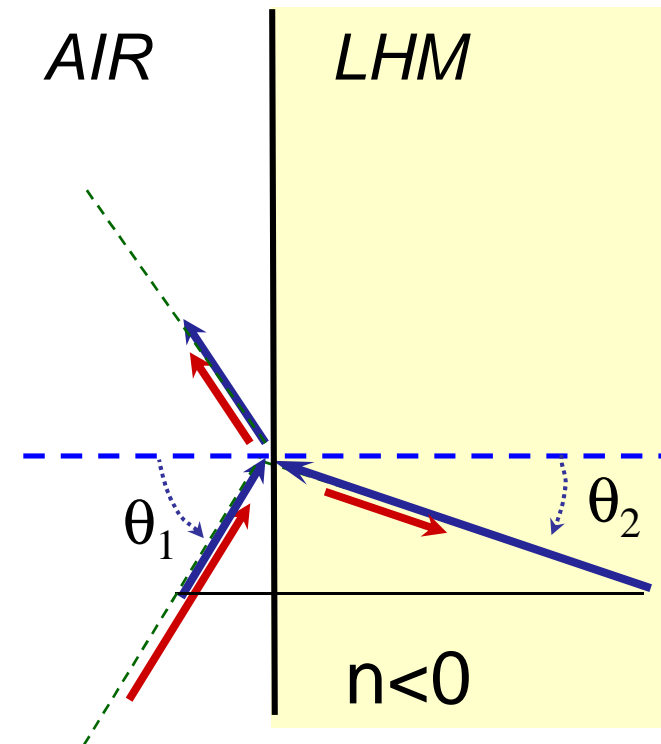
Left-handed slab



## Negative refraction - Snell's law reversion



$$\sin \theta_2 = \frac{1}{n} \sin \theta_1$$

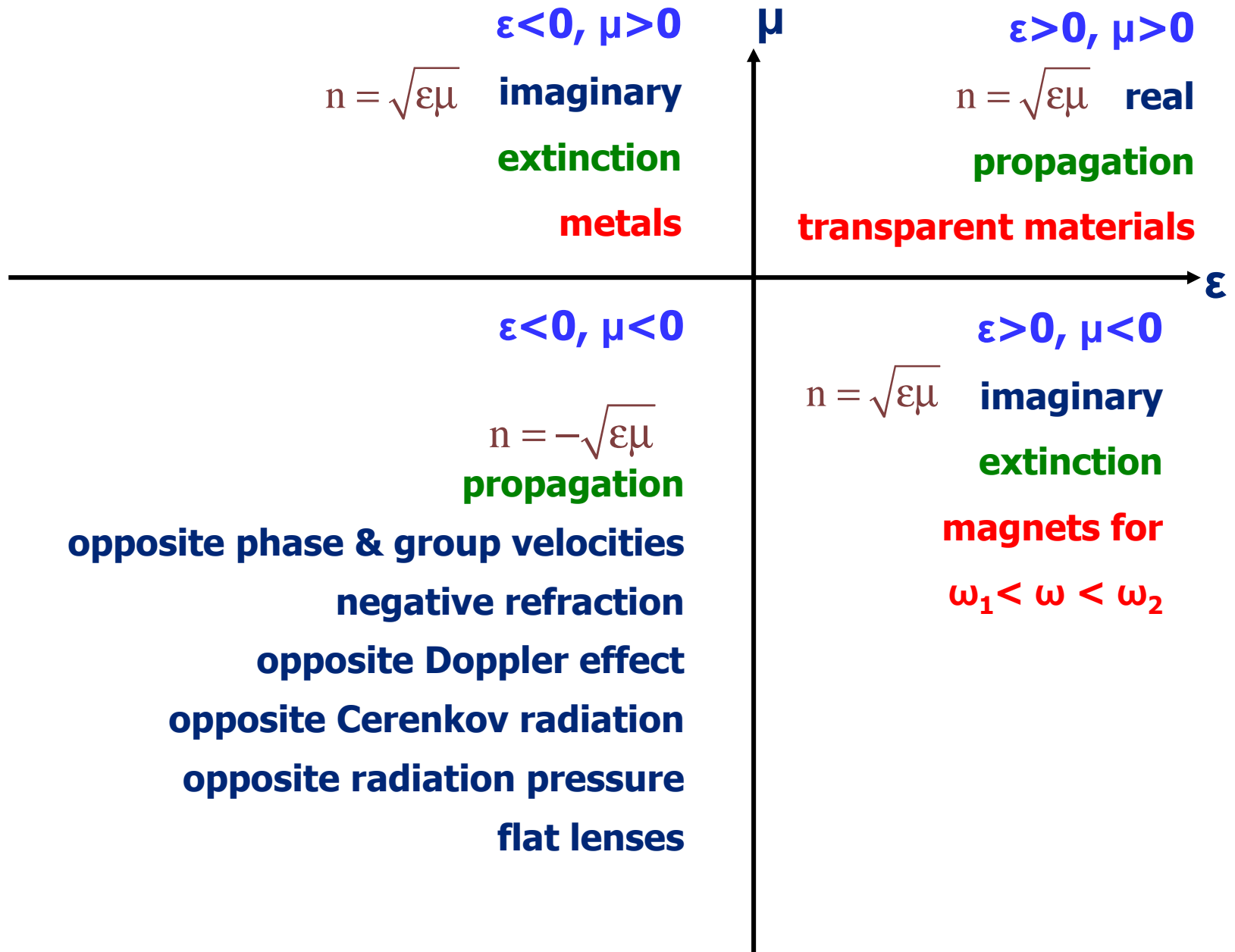


Snell-Descartes' law  
valid with  $n < 0$



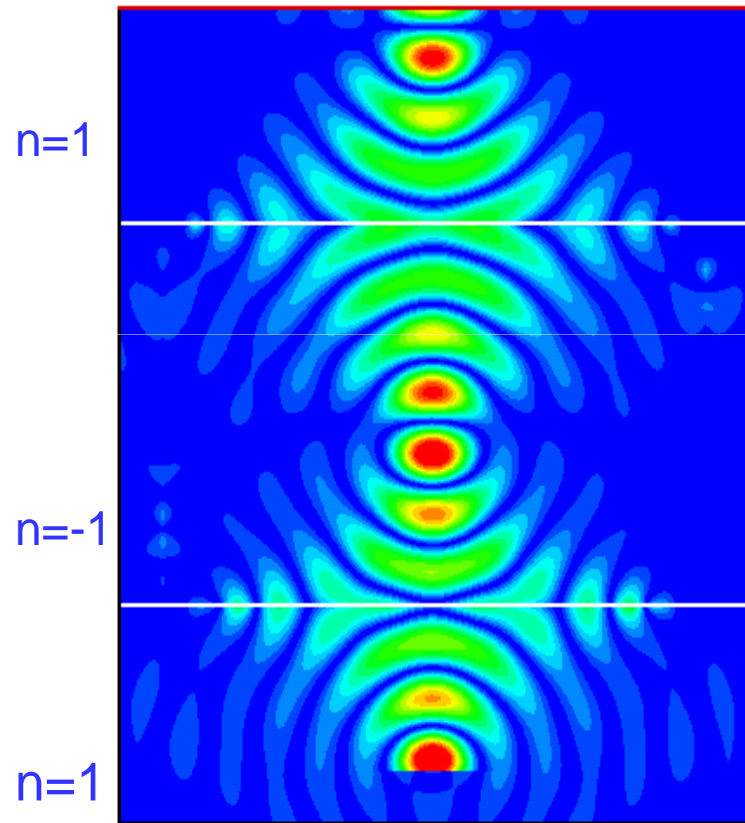
# Negative refraction - Snell's law reversion





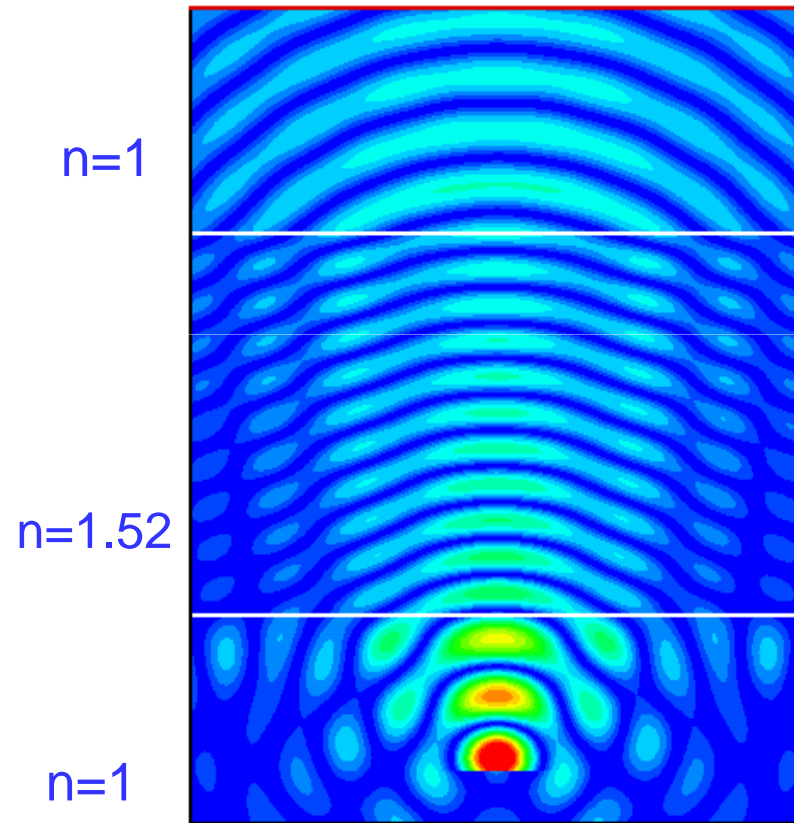
# Left-handed media as flat lenses

Left-handed



Source

Right-handed



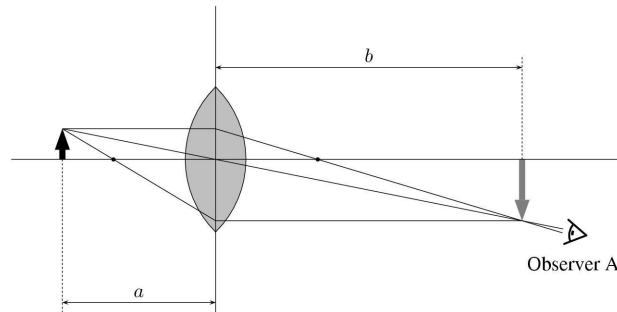
Source

# Left-handed media as flat lenses

Normal lens  
real, inverse 2d-image

$$\frac{1}{a} + \frac{1}{b} = \frac{1}{f}$$

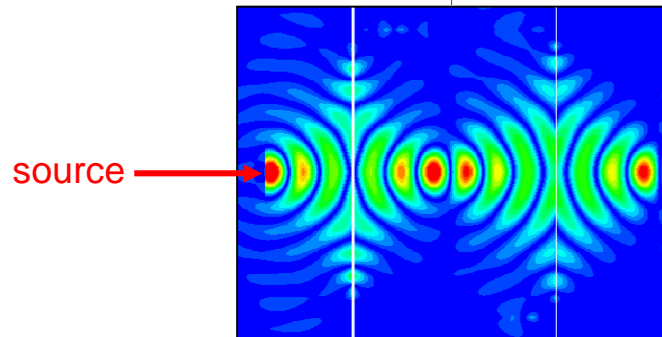
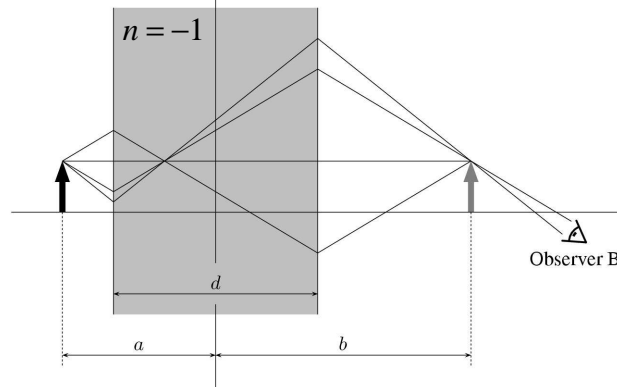
$$\frac{\partial b}{\partial a} = -\left(\frac{b}{a}\right)^2$$



Flat LH lens  
real, (semi-) 3d-image

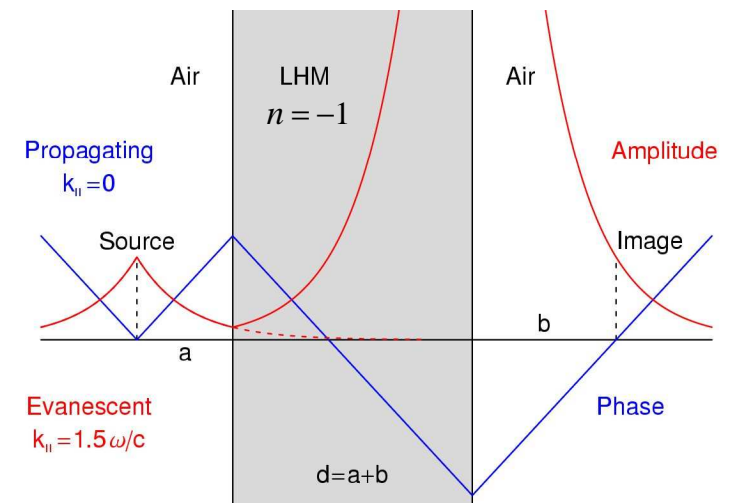
$$a + b = 2d$$

$$\frac{\partial b}{\partial a} = -1$$



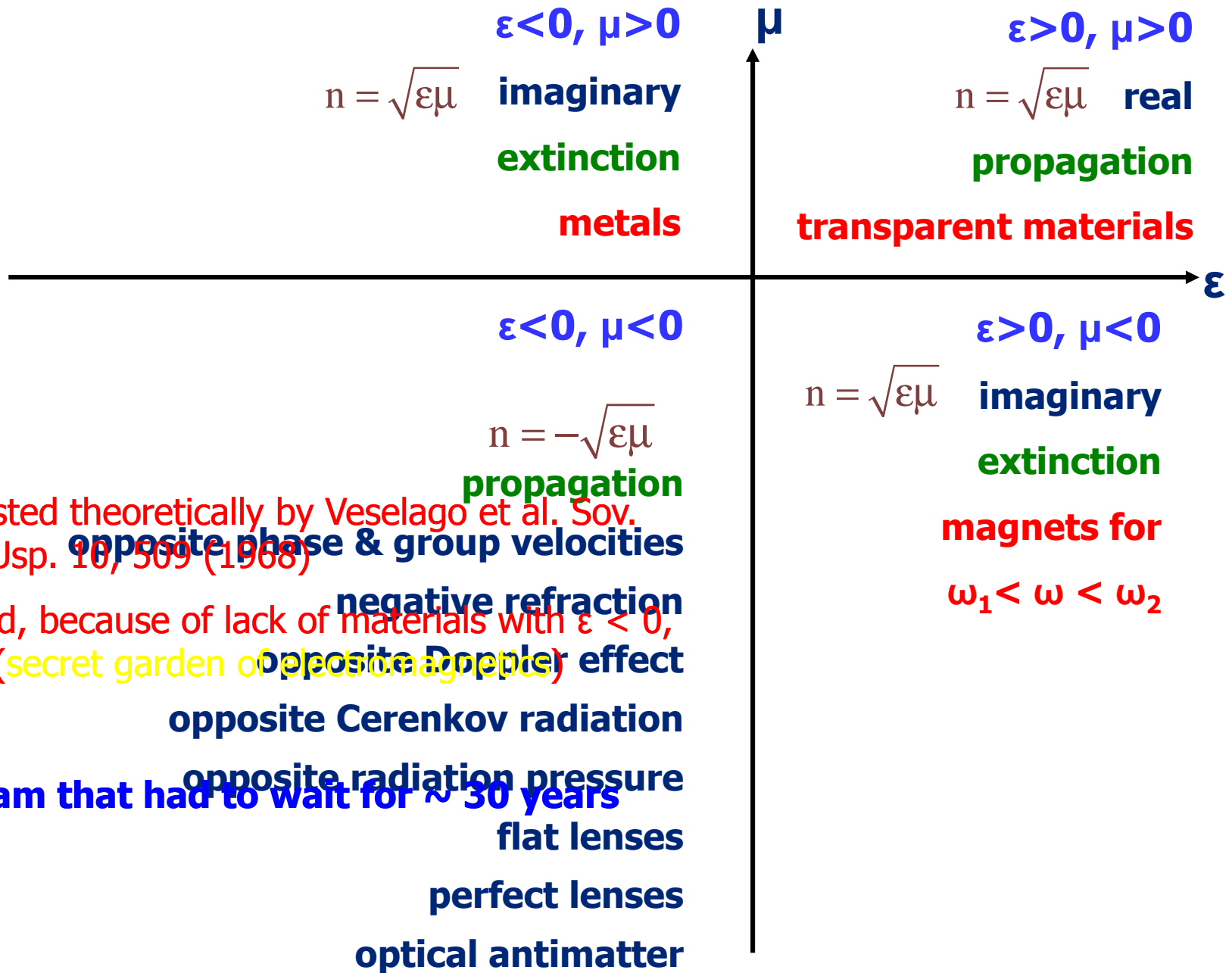
## Super lens

Optical lens = phase restoration  
(for propagating modes, evanescent waves are lost)



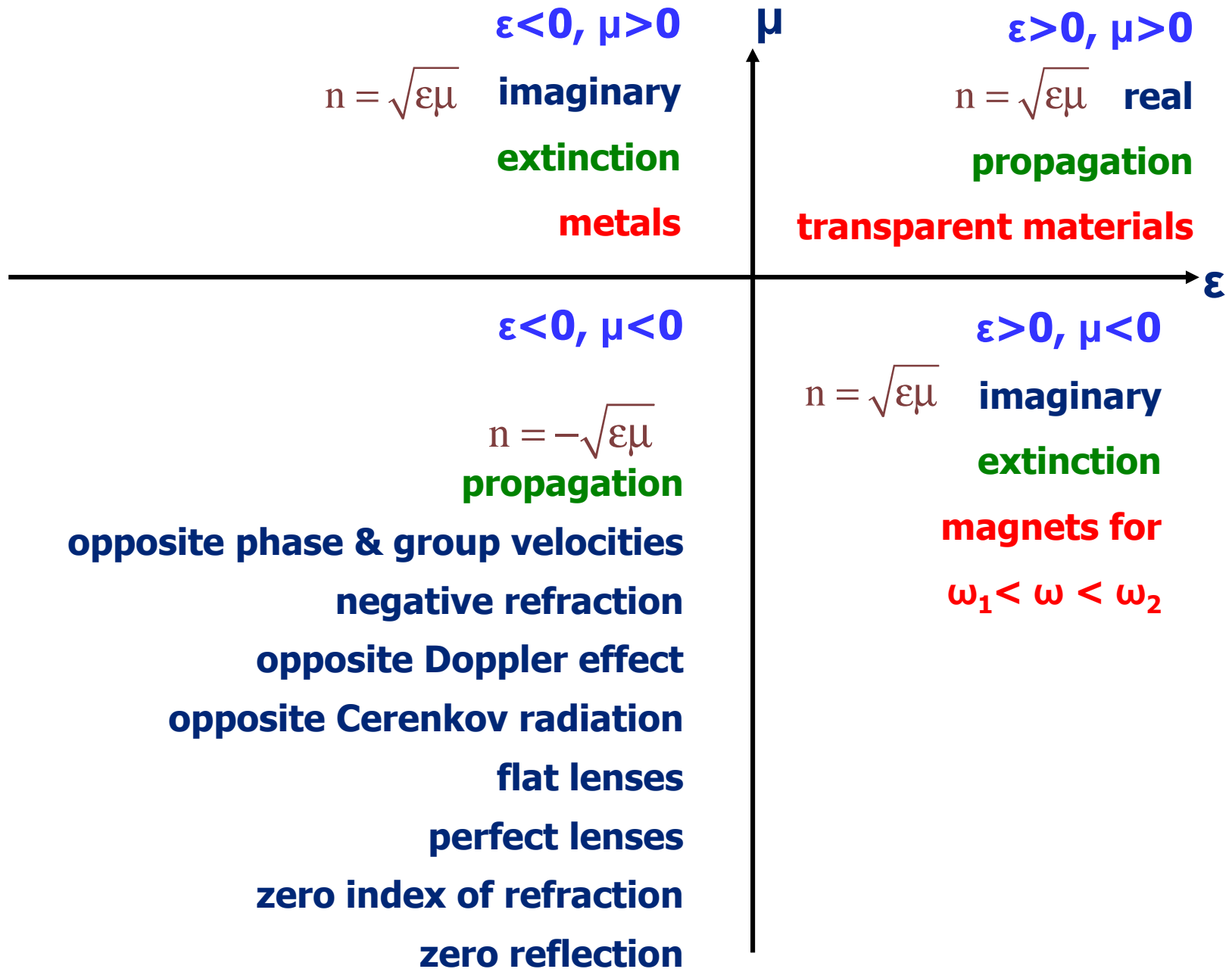
Super-lens = phase + amplitude restoration  
(propagating and evanescent waves are recovered)

$$k_{\perp} = -k'_{\perp}$$

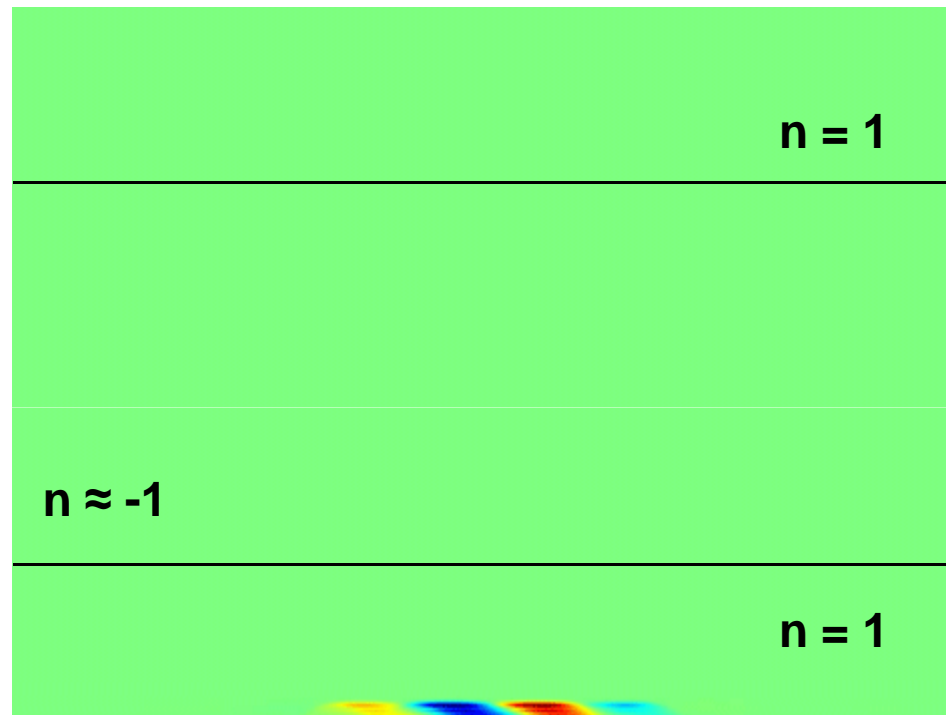


- Suggested theoretically by Veselago et al. Sov. Phys. Usp. 10, 509 (1968)
- Ignored, because of lack of materials with  $\epsilon < 0$ ,  $\mu < 0$  (secret garden of electromagnetics)

**A dream that had to wait for ~ 30 years**



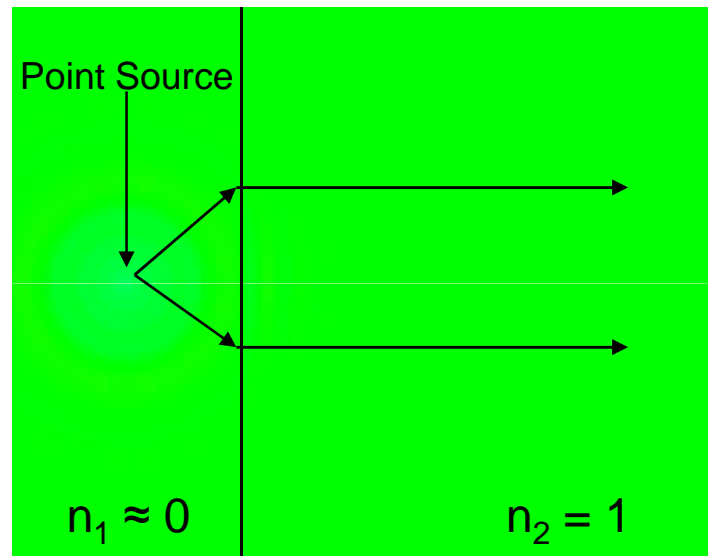
# Negative refraction in metamaterials



$$\epsilon = \mu = -1 + 0.001 i$$

Frequency=30 GHz,  $\lambda=0.01$  m, thickness of metamaterial =  $4 \lambda$

## Metamaterials with zero index of refraction



All the angles  $\theta_2$  should be zero, and therefore perpendicular to the surface



A dream comes true: Prof. Pendry suggests structures with  $\epsilon < 0$  and  $\mu < 0$

## Basic idea behind: RESONANCES

- **For  $\epsilon < 0$ : A wire medium.** It can yield  $\epsilon < 0$  in a tunable frequency range (Pendry et al., PRL 1996). Artificial dielectrics using metals!!
- **For  $\mu < 0$ : Split ring resonators.** They yield  $\mu < 0$  in a tunable frequency range (Pendry et al., IEEE, 1999). **Artificial magnetic materials using non-magnetic metals!!**
- **A combined medium can yield *both*  $\epsilon < 0$  and  $\mu < 0$  simultaneously.**
- A medium with negative index of refraction should be possible by using these suggested structures?



# Frequency dispersion of LH medium

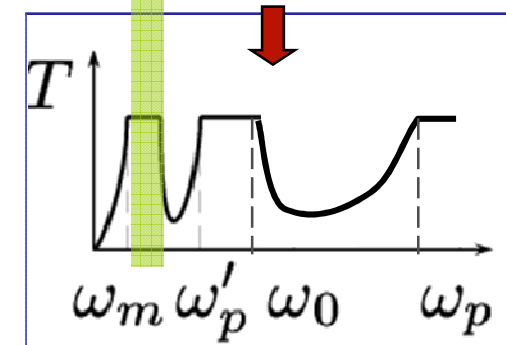
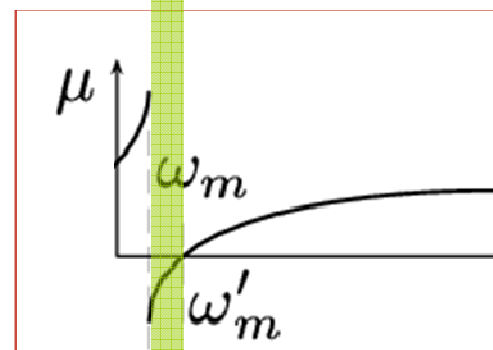
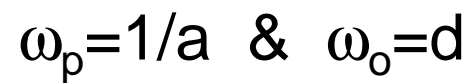
- Energy density in the dispersive medium



- Energy density  $W$  must be **positive** and this requires



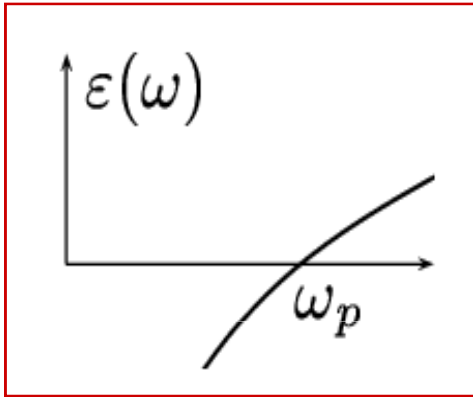
- LH medium is always dispersive
- According to the Kramers-Kronig relations -  
it is always dissipative



PRB 70, 201101 (2004)

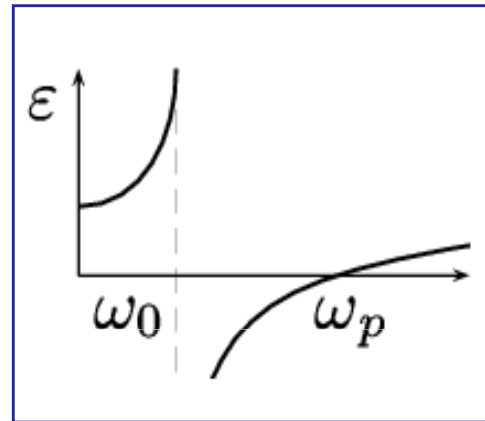
# Cut-wire response

Continuous-wires



$$m \frac{\partial^2 \mathbf{x}}{\partial t^2} = e\mathbf{E}$$

Cut-wires



$$m \frac{\partial^2 \mathbf{x}}{\partial t^2} = e\mathbf{E} - m\omega_0^2 \mathbf{x}$$

$$\mathbf{j} = -i\omega \mathbf{P} = ne\mathbf{v} = ne(-i\omega \mathbf{x}) \Rightarrow \mathbf{P} = ne\mathbf{x} \Rightarrow \mathbf{P} = f(\mathbf{E})$$

# The SRR is like an LC resonator

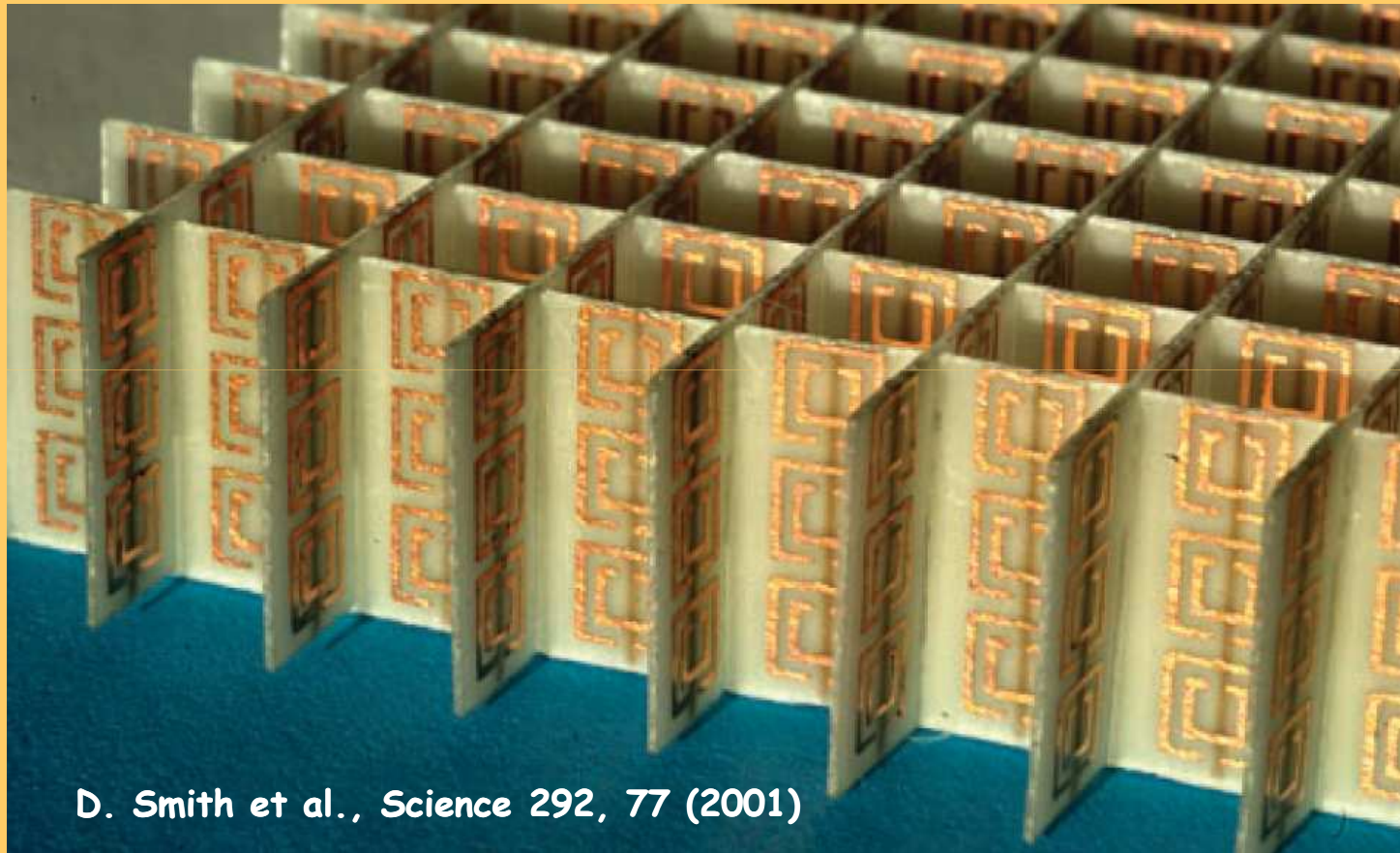


# First Left-Handed Test Structure



UCSD, PRL 84, 4184 (2000)

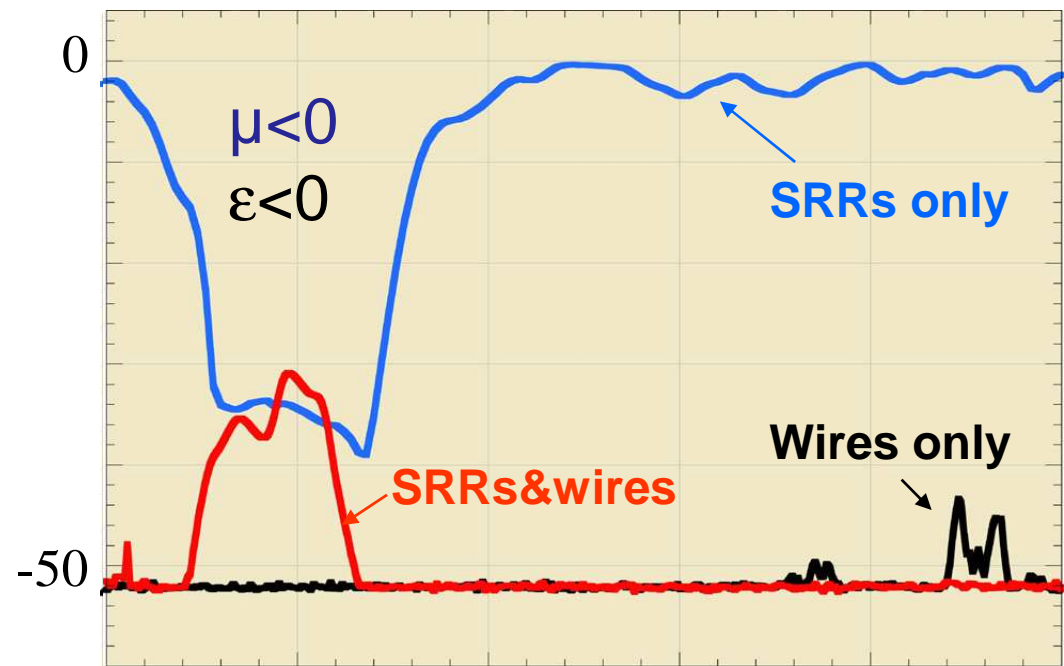
David Smith (UCSD) used Prof. Pendry's suggested structures to demonstrate the first material with a negative index of refraction



D. Smith et al., Science 292, 77 (2001)



# First experimental verification of a NIM

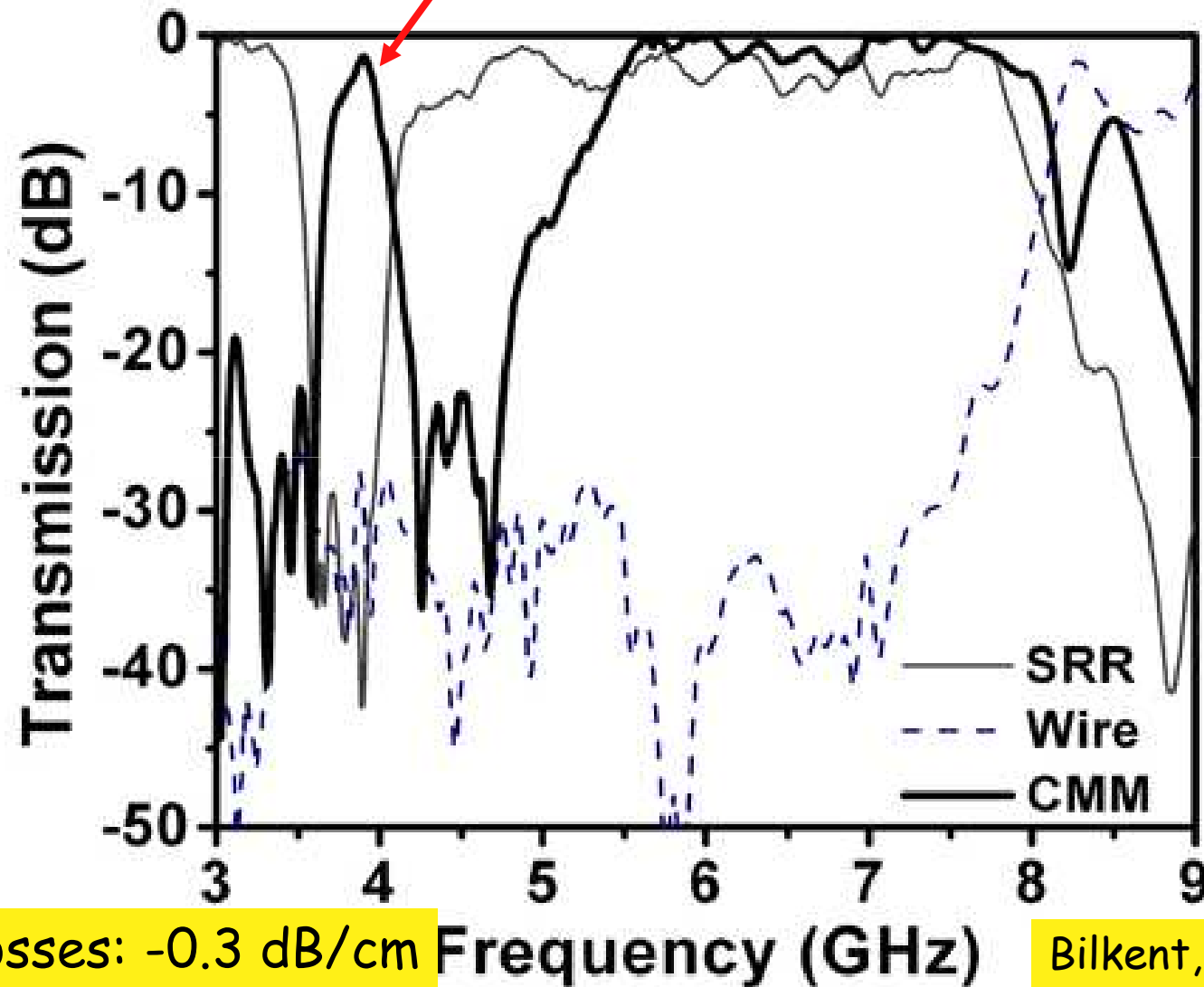


UCSD, PRL 84, 4184 (2000)

Frequency (GHz)



# Best LH peak in a left-handed material



Peak at  
 $f=4$  GHz  
 $\lambda=75$  mm

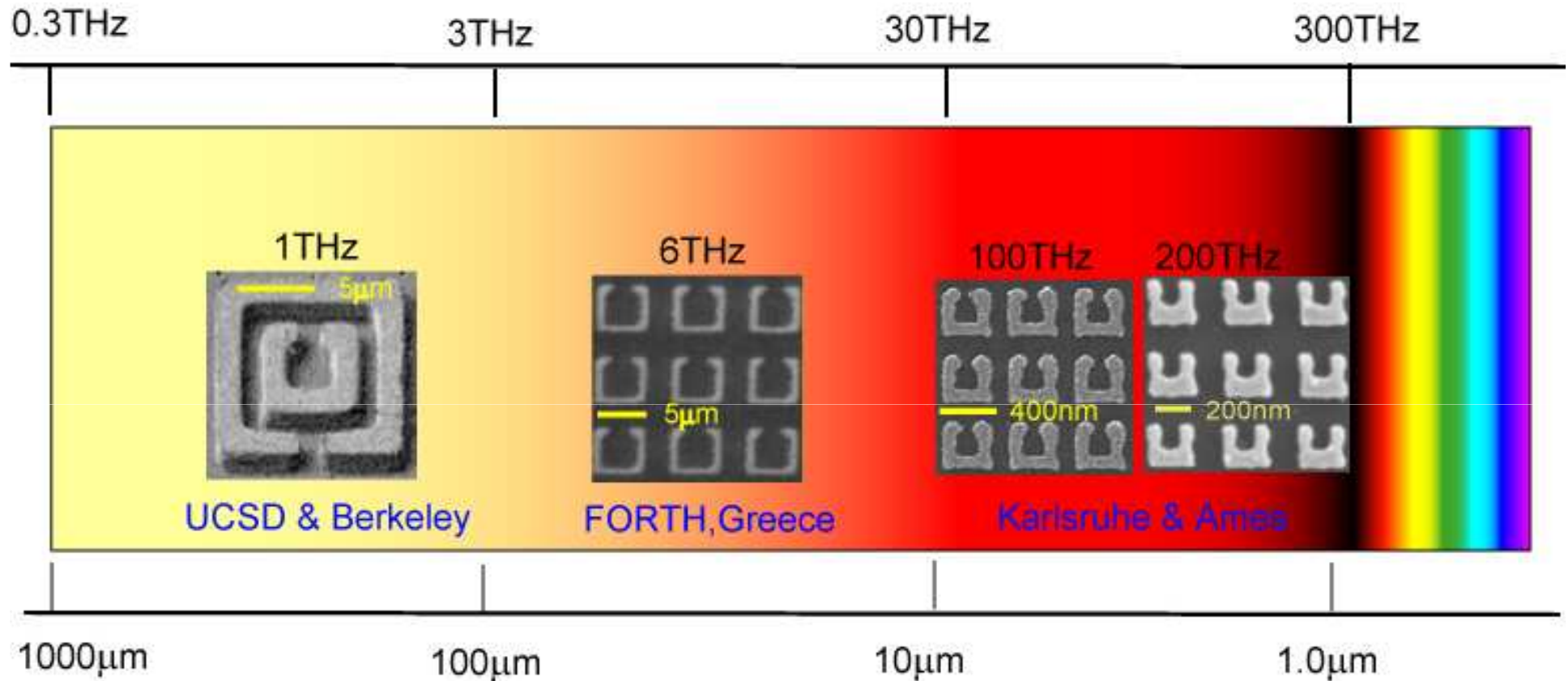
much larger  
than the  
size of SRR

$a=3.6$  mm

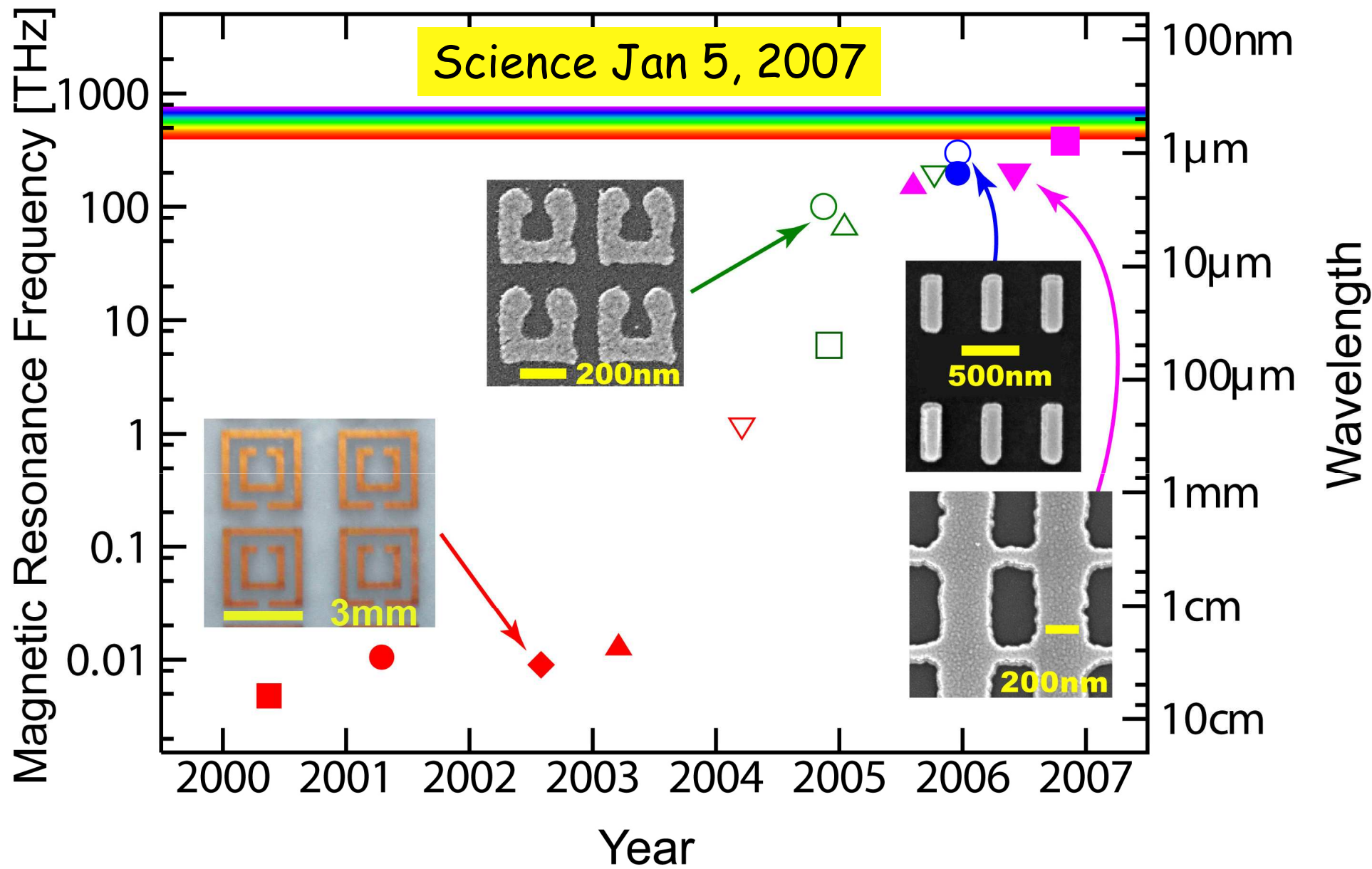
Losses:  $-0.3$  dB/cm

Bilkent, Crete, & Ames  
Optics Lett. 29, 2623 (2004)

# The progress of scaling metamaterials



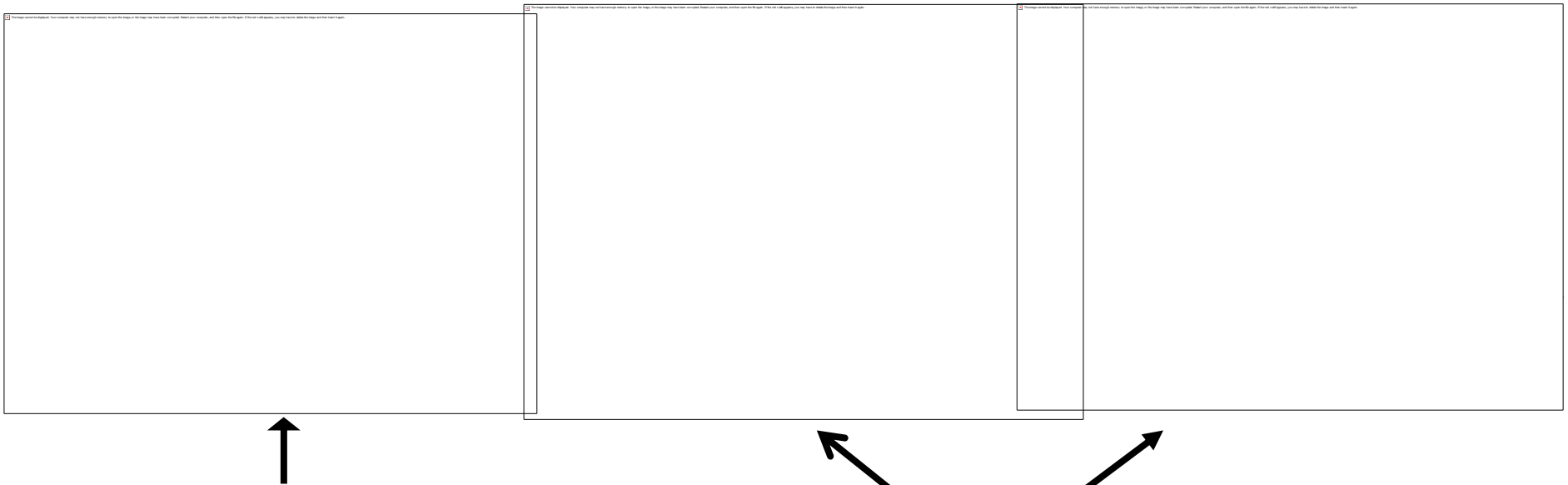
Advances in scaling metamaterials with artificial magnetic response for high-frequency structures has been rapid. The 1, 6 and 100 THz models were fabricated in 2004, and the 200 THz in 2005.



Solid symbol  $n < 0$

Open symbol  $\mu < 0$

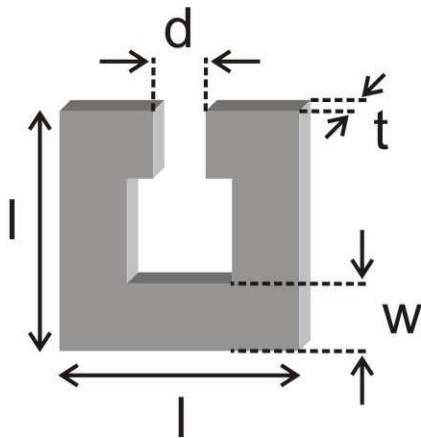
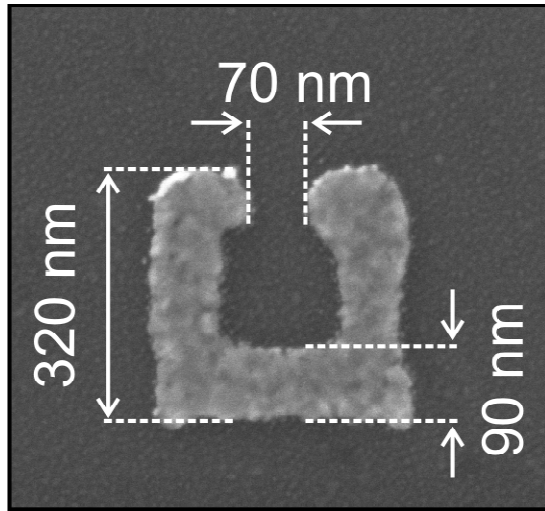
Metamaterials structures shown in these split-ring resonators (SRRs) are amenable to manufacture by common planar lithography.



It has a **magnetic response** **perpendicular** to the plane, which is **difficult** to **detect** by direct incidence measurements.

Use of multilayer processing can be used to fabricate metamaterials that give both **negative  $\epsilon$  and  $\mu$** , as well as  **$n$**  for **perpendicular** propagation.

# Estimating the LC-resonance



$$L = \mu_0 \frac{l^2}{t} = 5.6 \text{ pH}$$

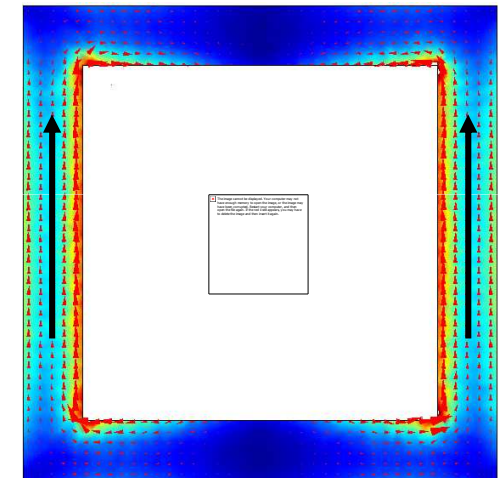
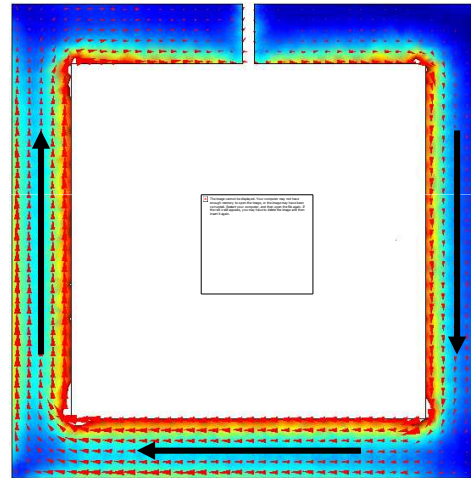
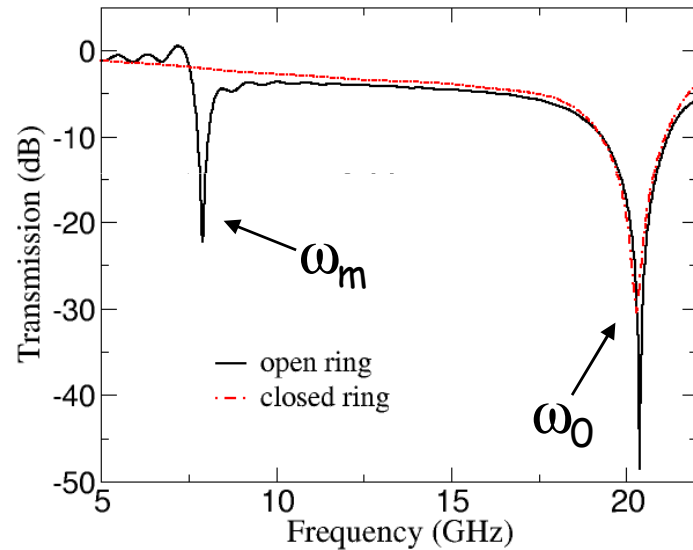
$$C = \epsilon_0 \epsilon_c \frac{wt}{d} = 0.5 \text{ aF}$$

$$\omega_{LC} = \frac{1}{\sqrt{LC}} = 2\pi \cdot 100 \text{ THz}$$

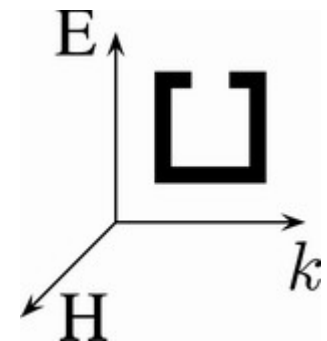
$$\Rightarrow \lambda_{LC} = l \cdot 2\pi \sqrt{\epsilon_c} \sqrt{\frac{w}{d}} \approx 3 \mu\text{m}$$

# Electric response of the SRRs?

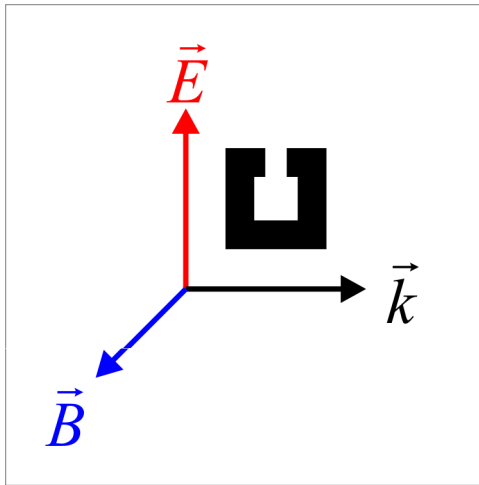
SRRs have also resonant electric response, cut-wire like



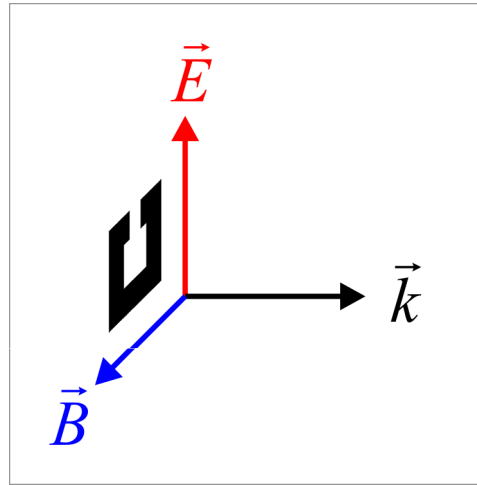
Closing the gap: **Magnetic response off**  
**Electric response unaffected**



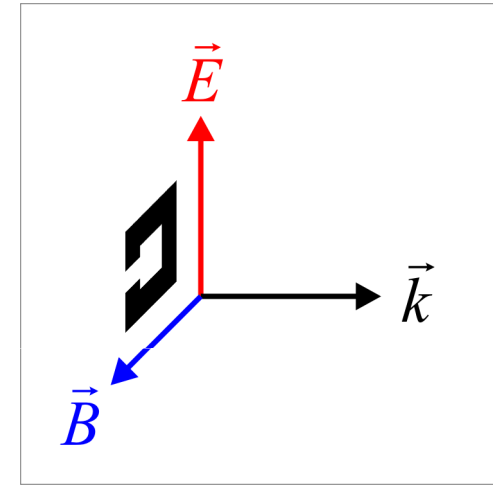
# Polarization dependence



Coupling to  $LC$ -resonance via  $B$ , not accessible for normal incidence.



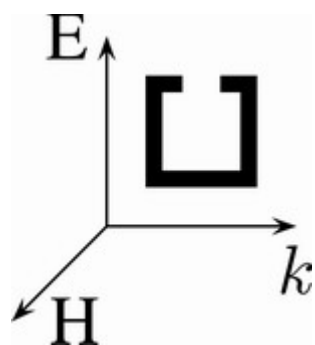
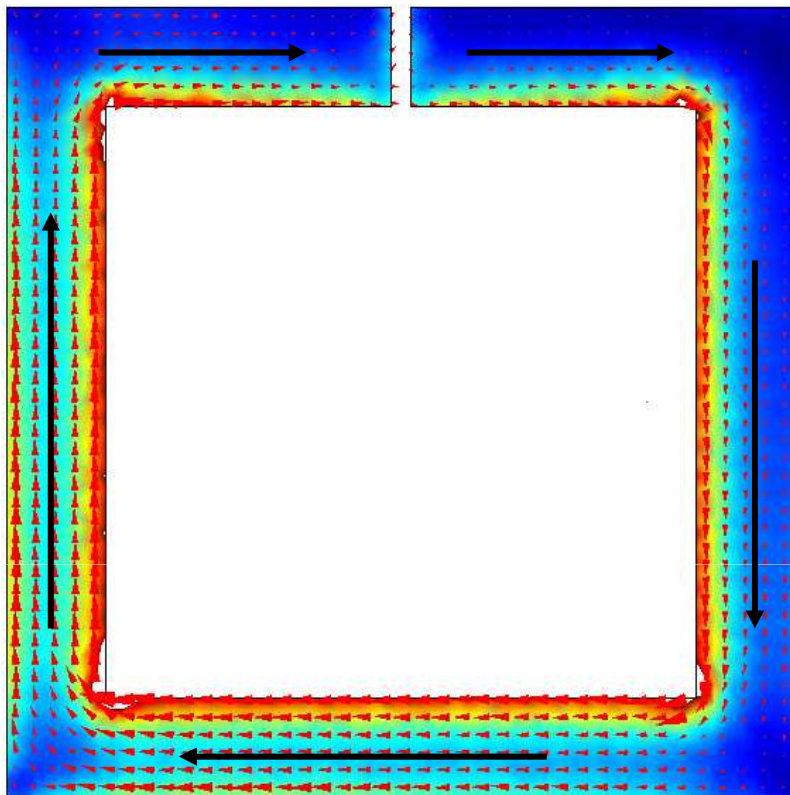
No coupling to  $LC$ -resonance, normal incidence.



Coupling to  $LC$ -resonance via  $E$ , accessible for normal incidence

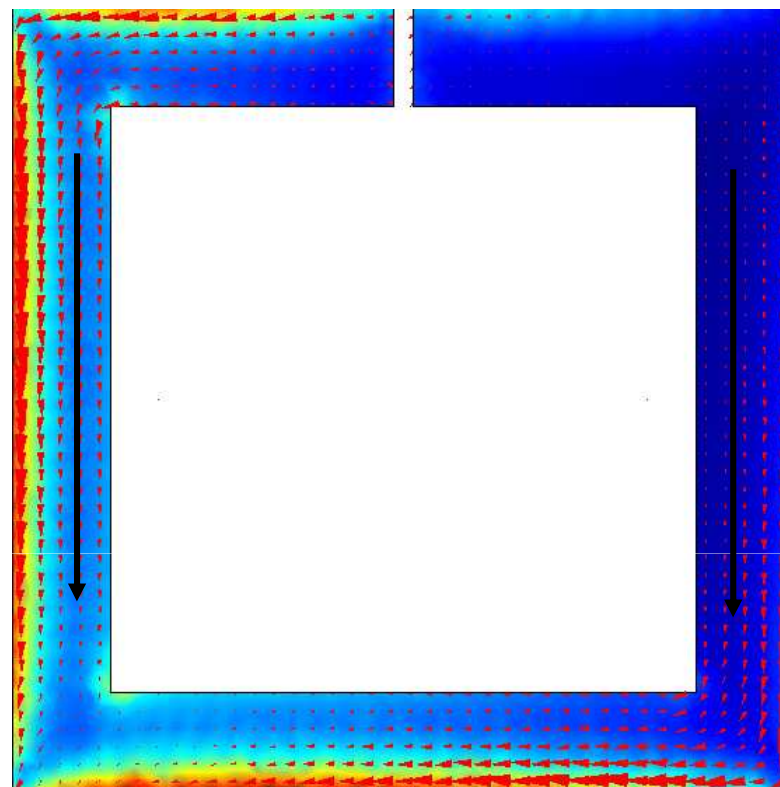


## Magnetic Resonance



27.4 THz

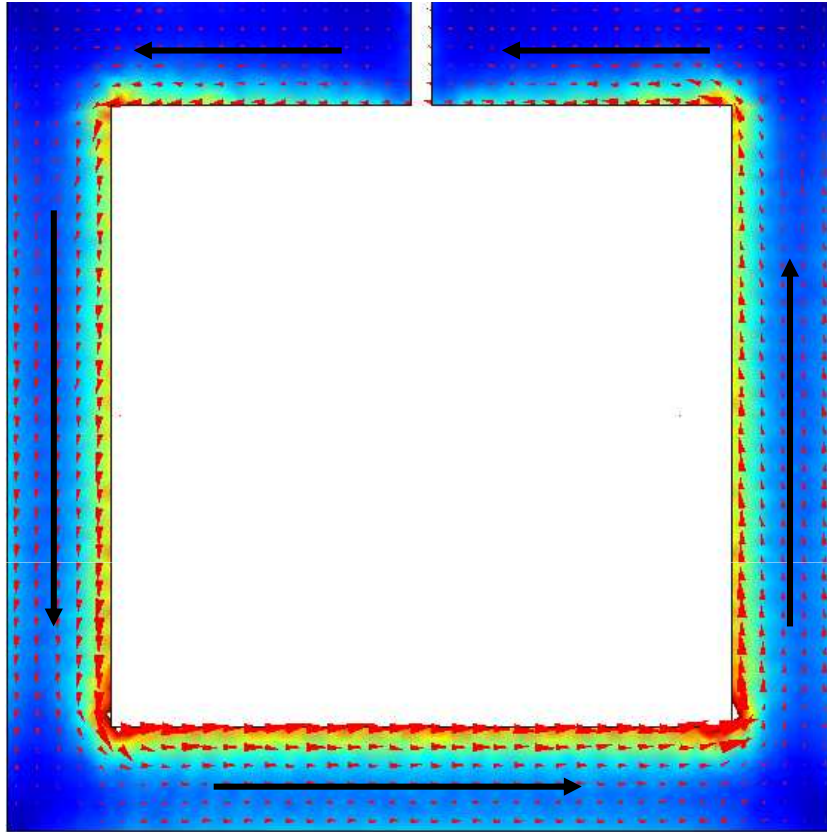
## Electric Resonance



90.3 THz

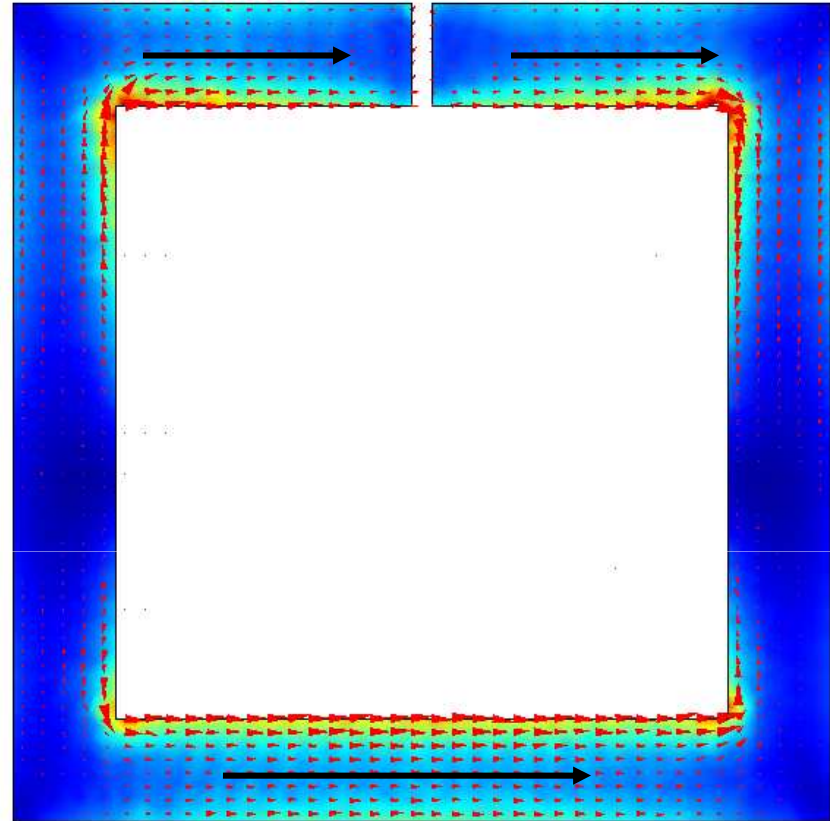


## Magnetic Resonance

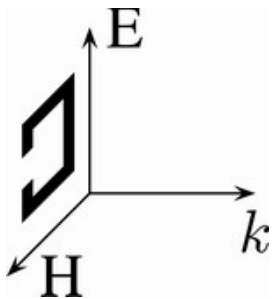


27.4 THz

## Electric Resonance

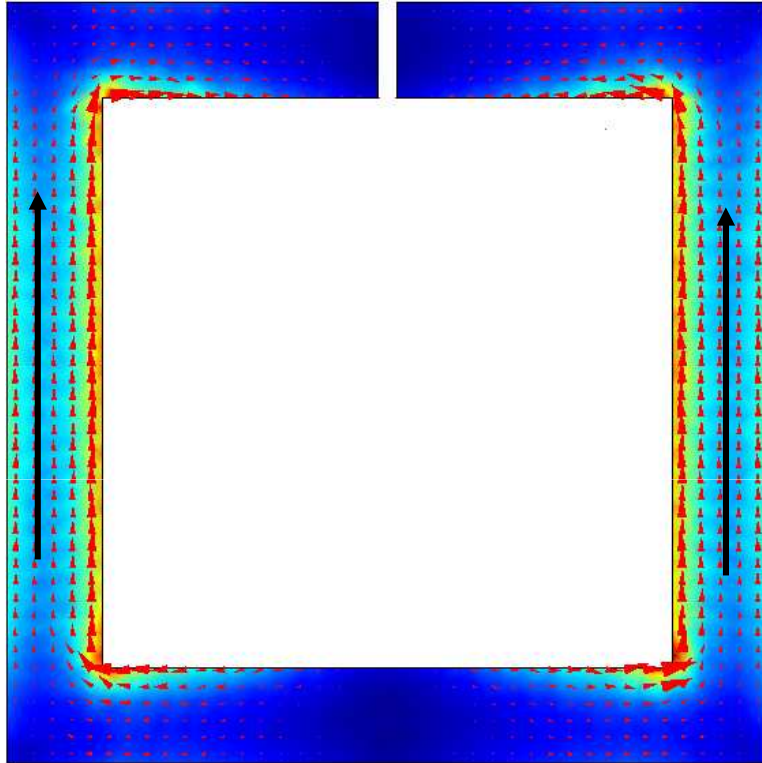


117.9 THz

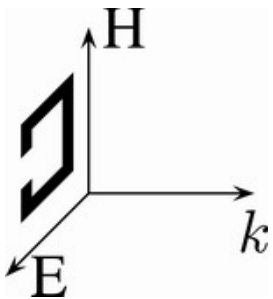


The magnetic resonance frequency is the same for both polarizations  
The electric resonance frequency is higher than the previous case

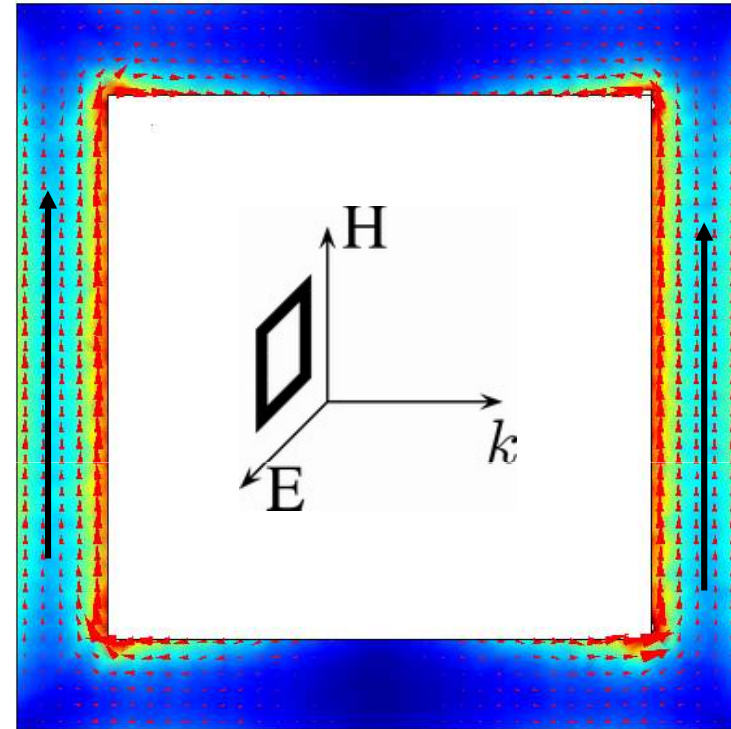
## Electric Resonance



87.4  
THz



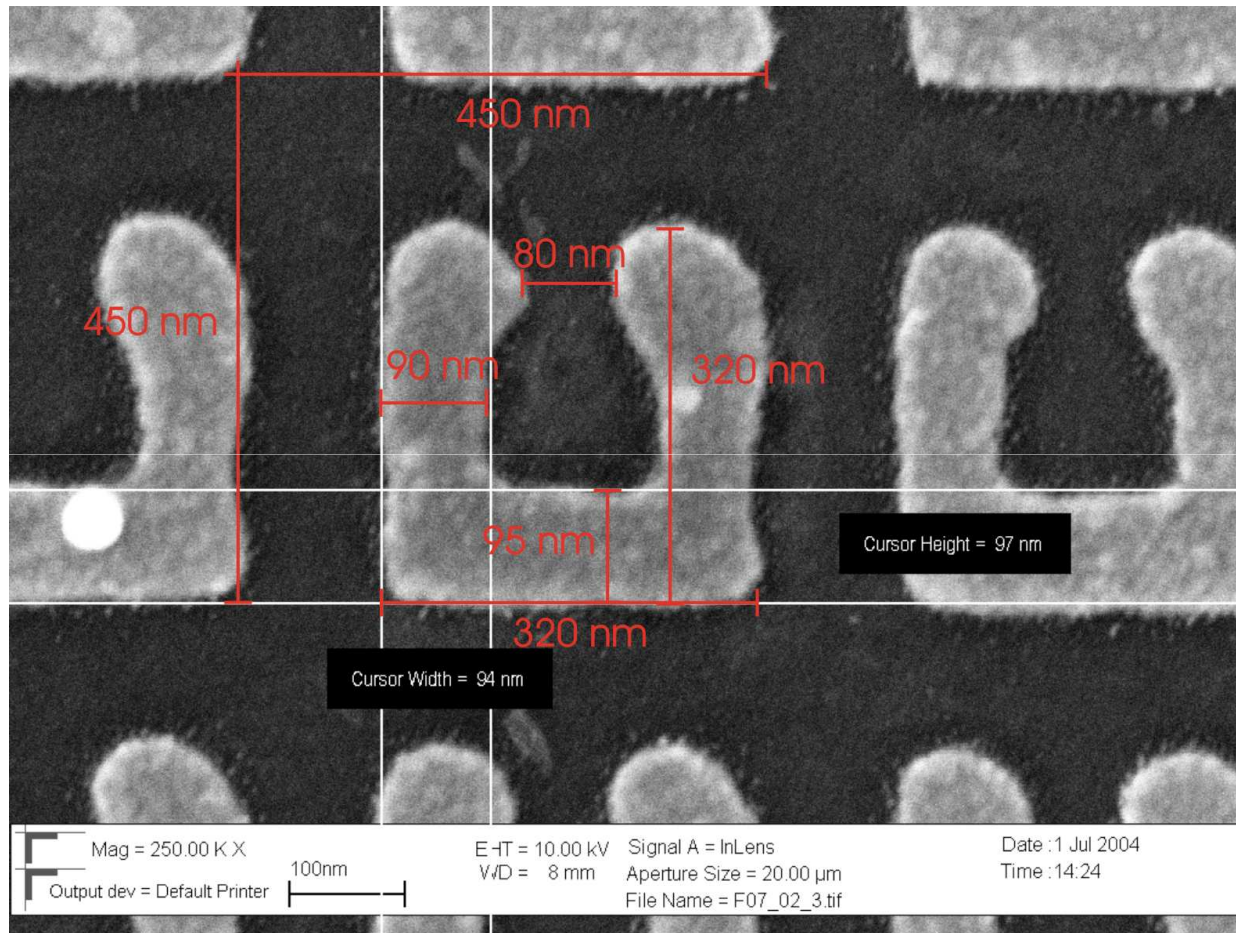
## Electric Resonance



87.4 THz  
**close ring**

Notice that the electric resonance frequency is the same for the **open** and **closed ring** SRR for this incident polarization

# Magnetic response at 100 THz almost optical frequencies

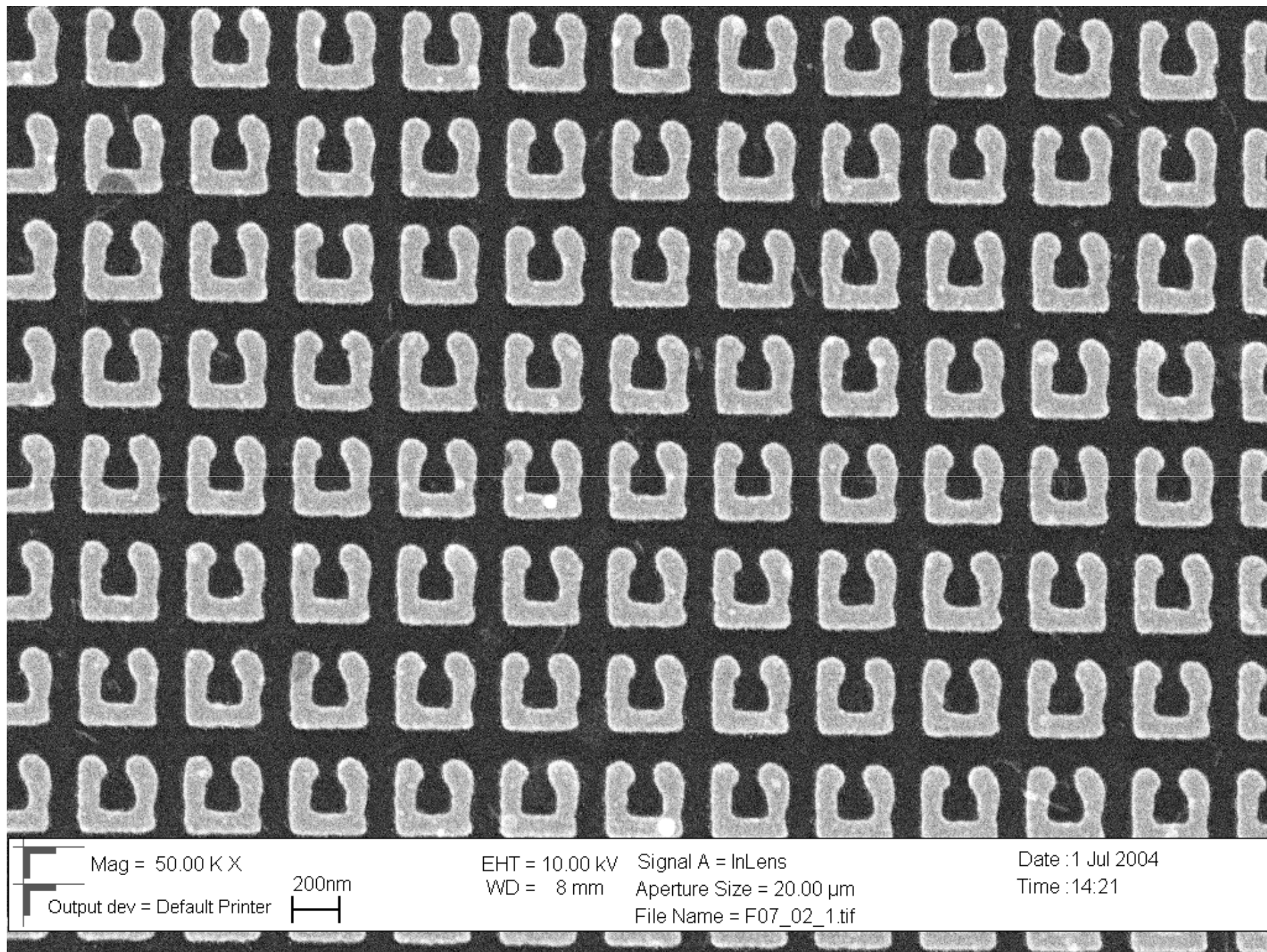


$\lambda \sim 10 \text{ a}$

Science, Nov 19, 2004

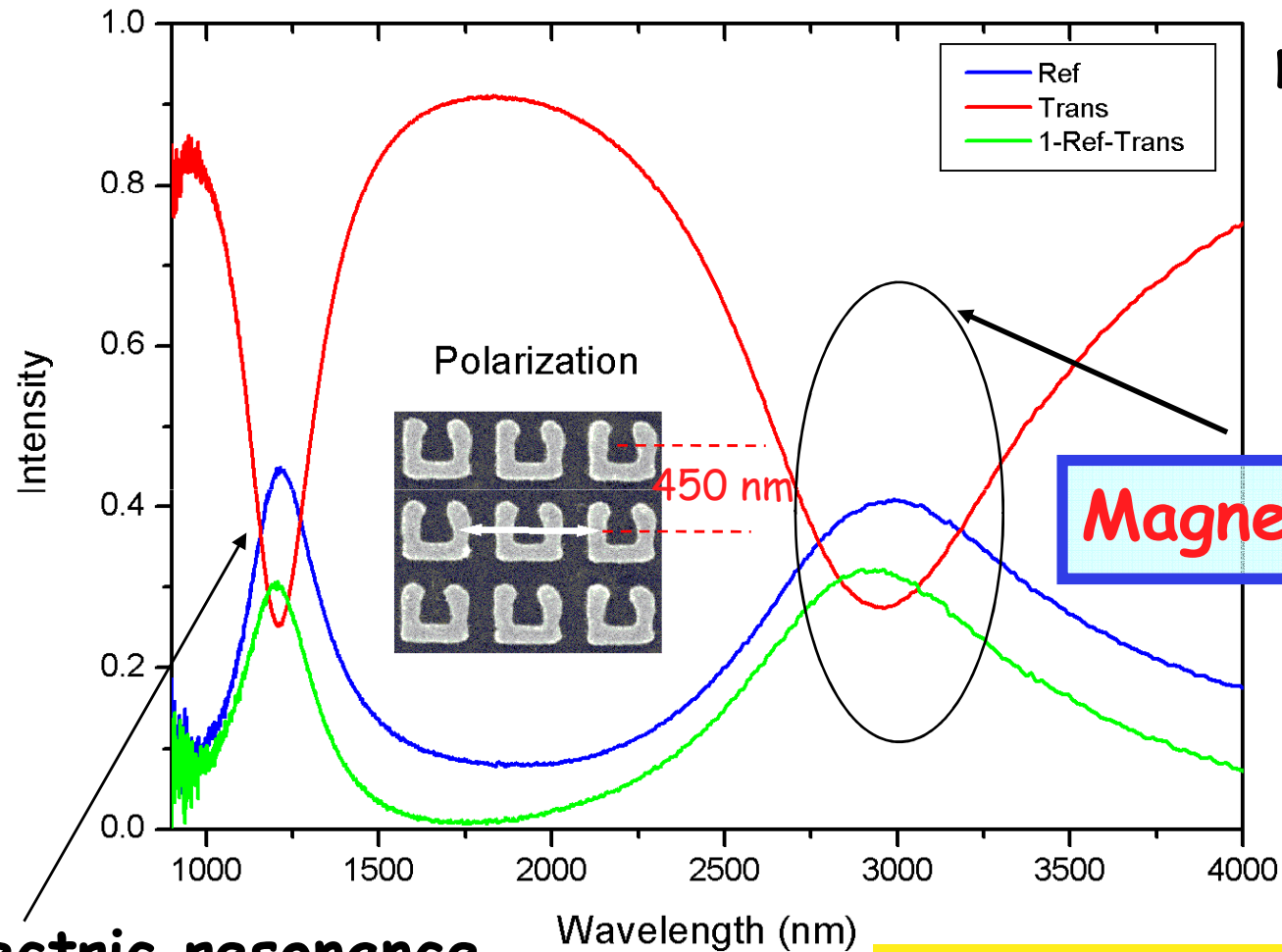
Univ. of Karlsruhe, & Ames Lab



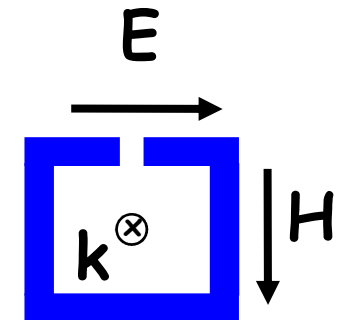


# Magnetic response at 100 THz

Experimental spectra for sample



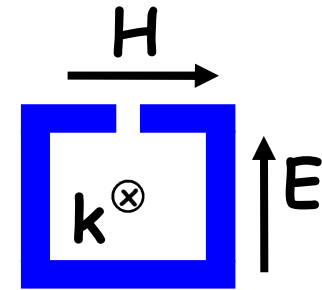
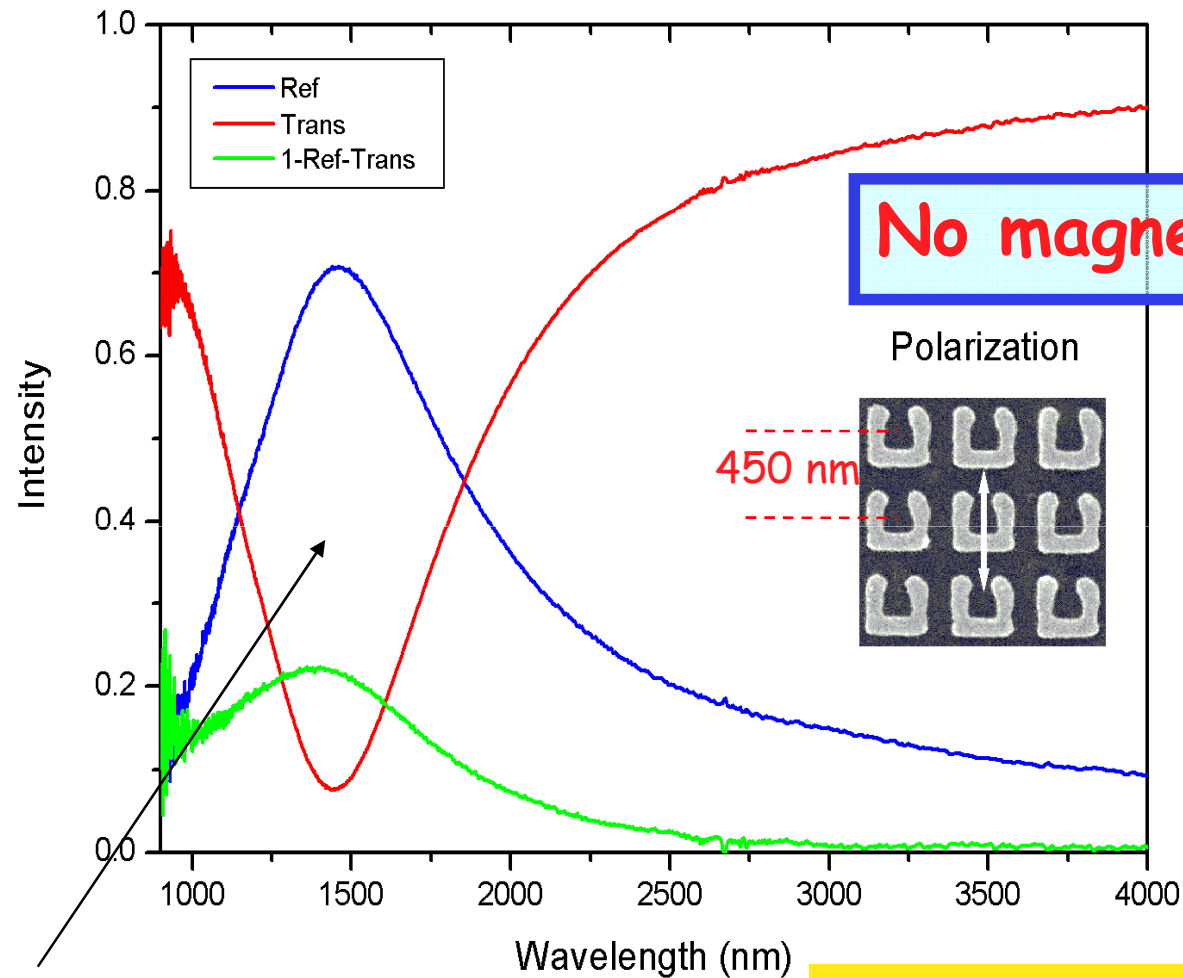
Normal illumination



Science, Nov 19, 2004

Univ. of Karlsruhe and Ames Lab

## Experimental spectra for sample

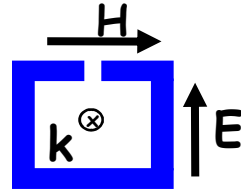
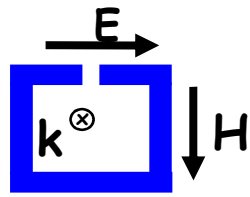


No magnetic resonance

Electric resonance

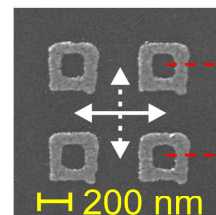
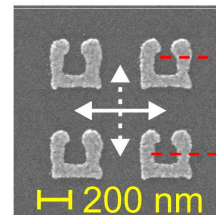
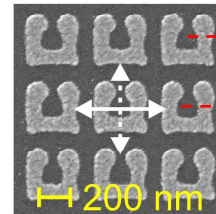
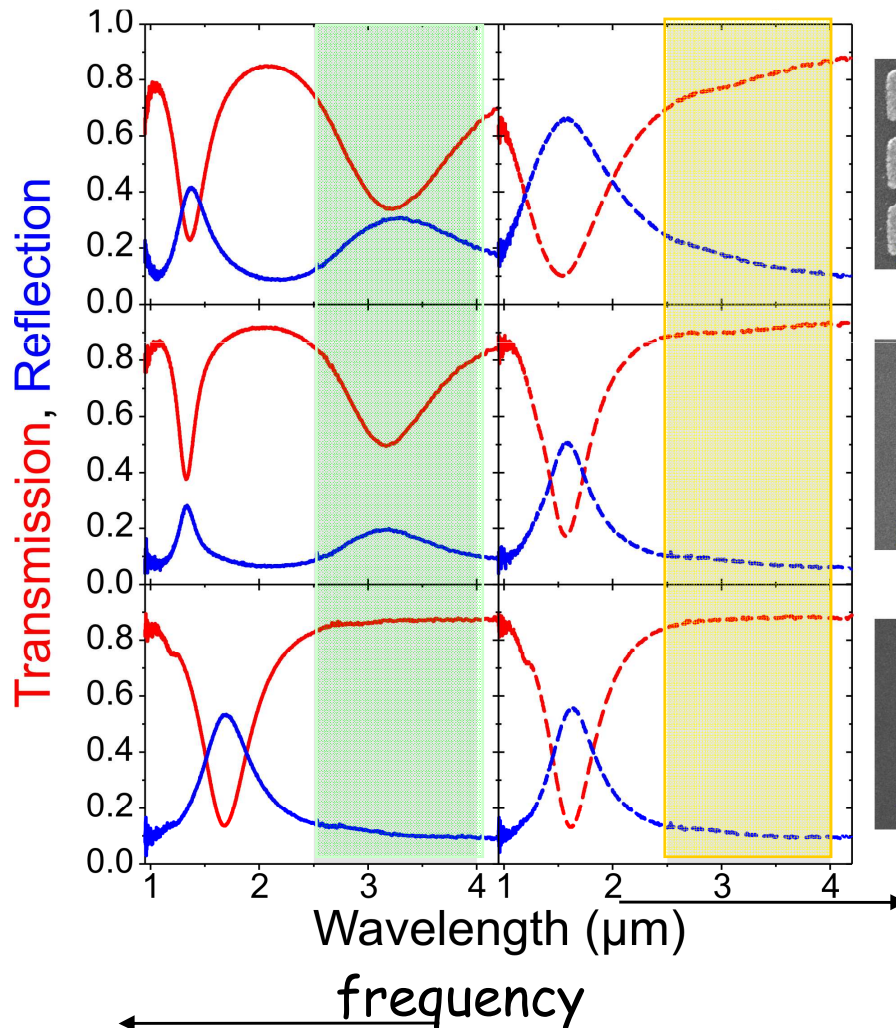
Science, Nov 19, 2004

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Magnetic resonance No magnetic resonance

Magnetic resonance depends on feature size and is independent of the lattice constant



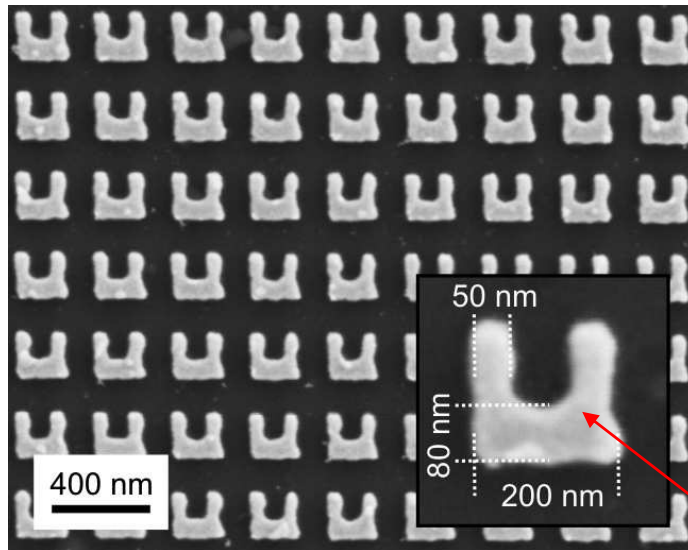
Closed rings  
Resonance disappears

Science, Nov 19, 2004



# Negative $\epsilon$ at telecommunication wavelengths ( $\sim 1.5 \mu\text{m}$ )

Electric coupling to the  
magnetic resonance

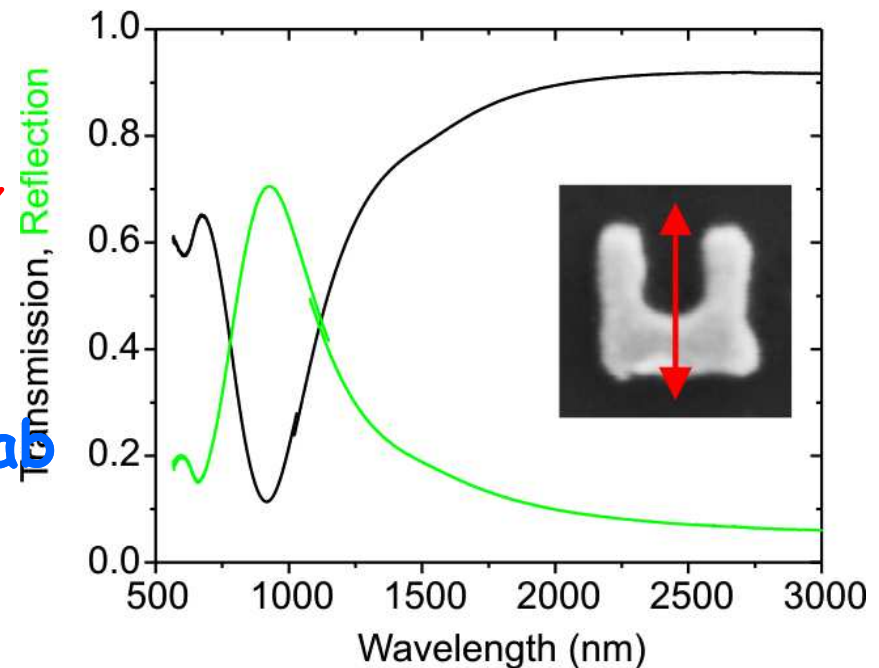
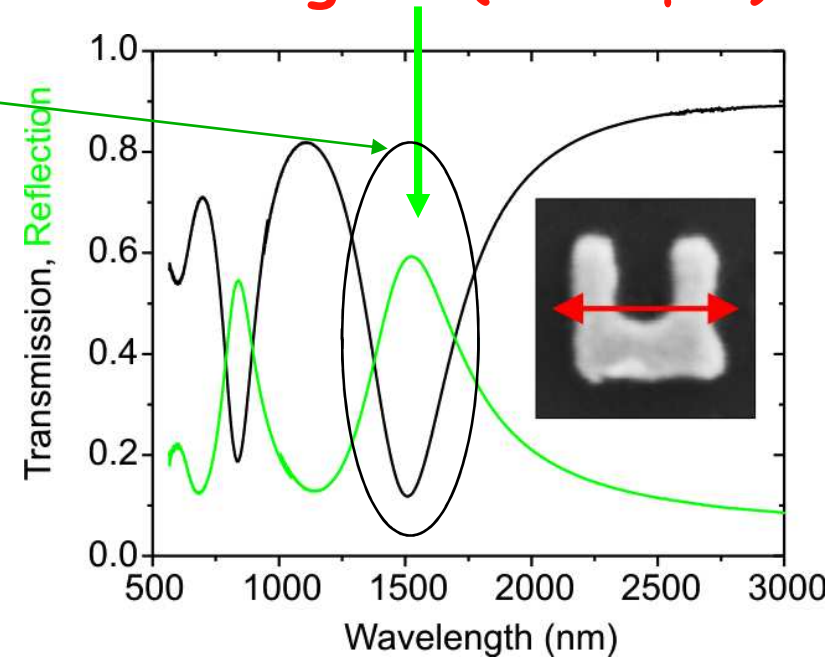


Au

*Fabricated with e-beam lithography*

Univ. of Karlsruhe and Ames Lab

PRL 95, 203901 (2005)





# Going to THz Frequencies

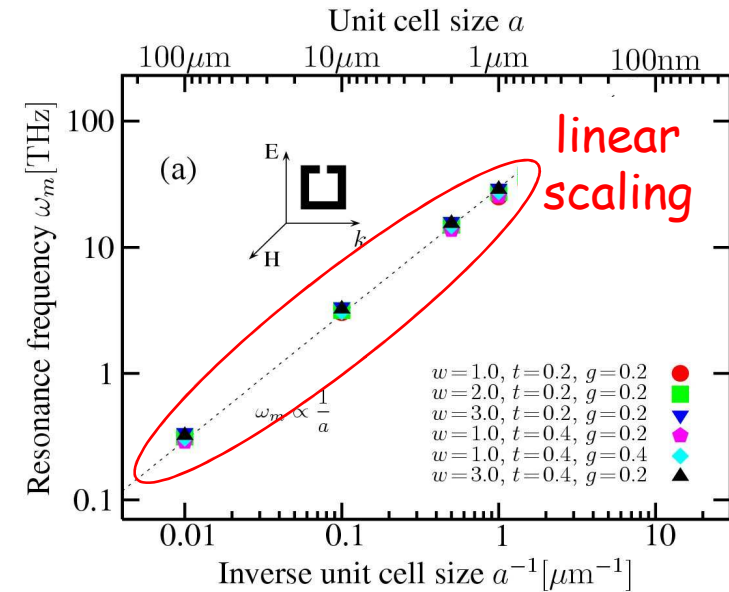
## Idea: geometric scaling

Metals are near-perfect conductors  
LC-resonator

$$C \approx \epsilon_0 \epsilon_{rel} \frac{A}{d} \quad \mu_0 \frac{\pi R^2}{l} \approx L \approx \mu_0 R \left( \log \frac{8R}{r_0} - 2 \right)$$

densely stacked rings                      sparse rings

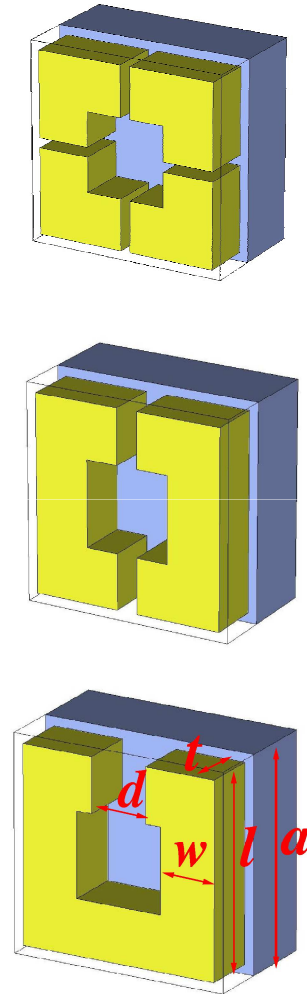
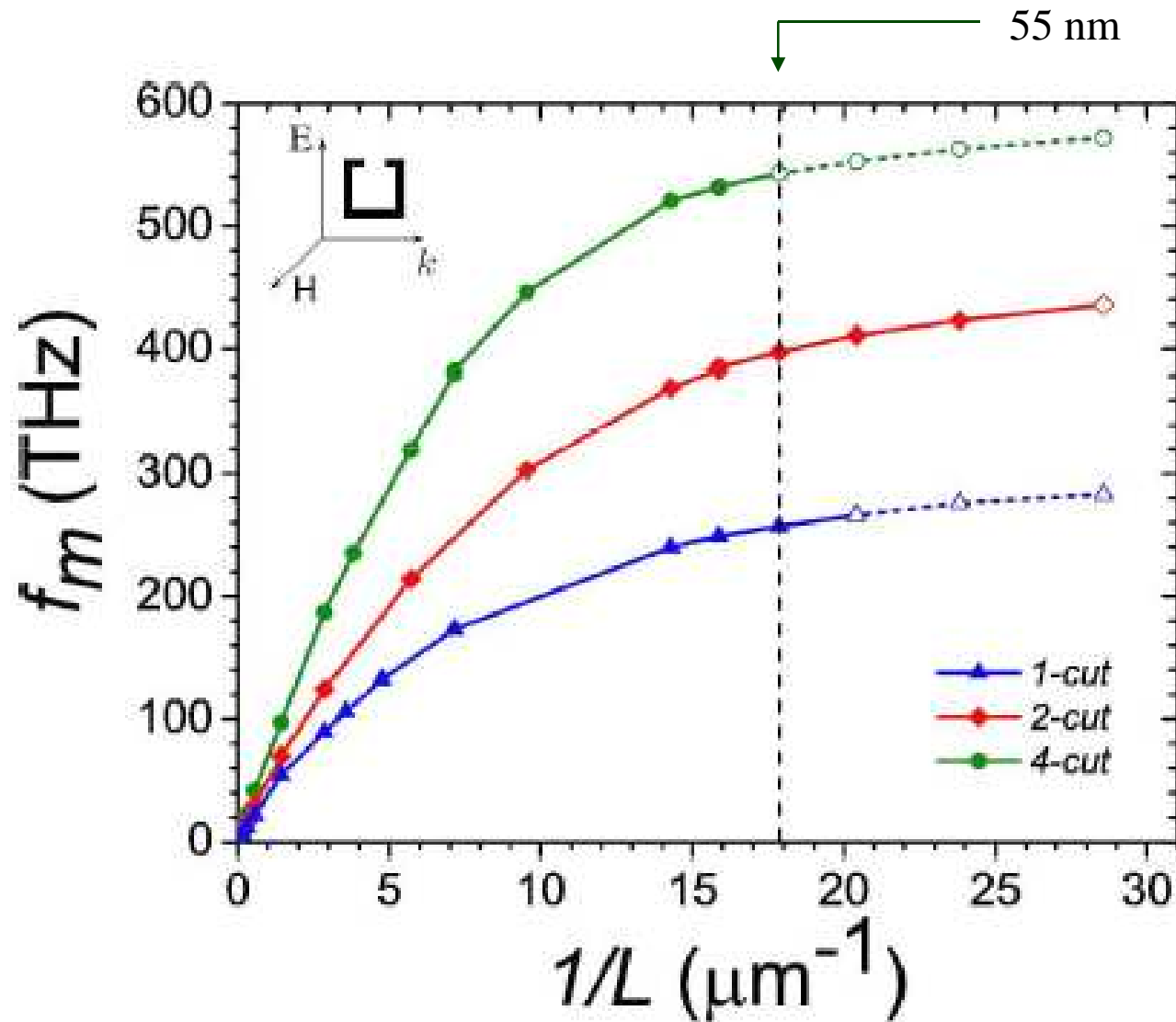
Depends on geometry only



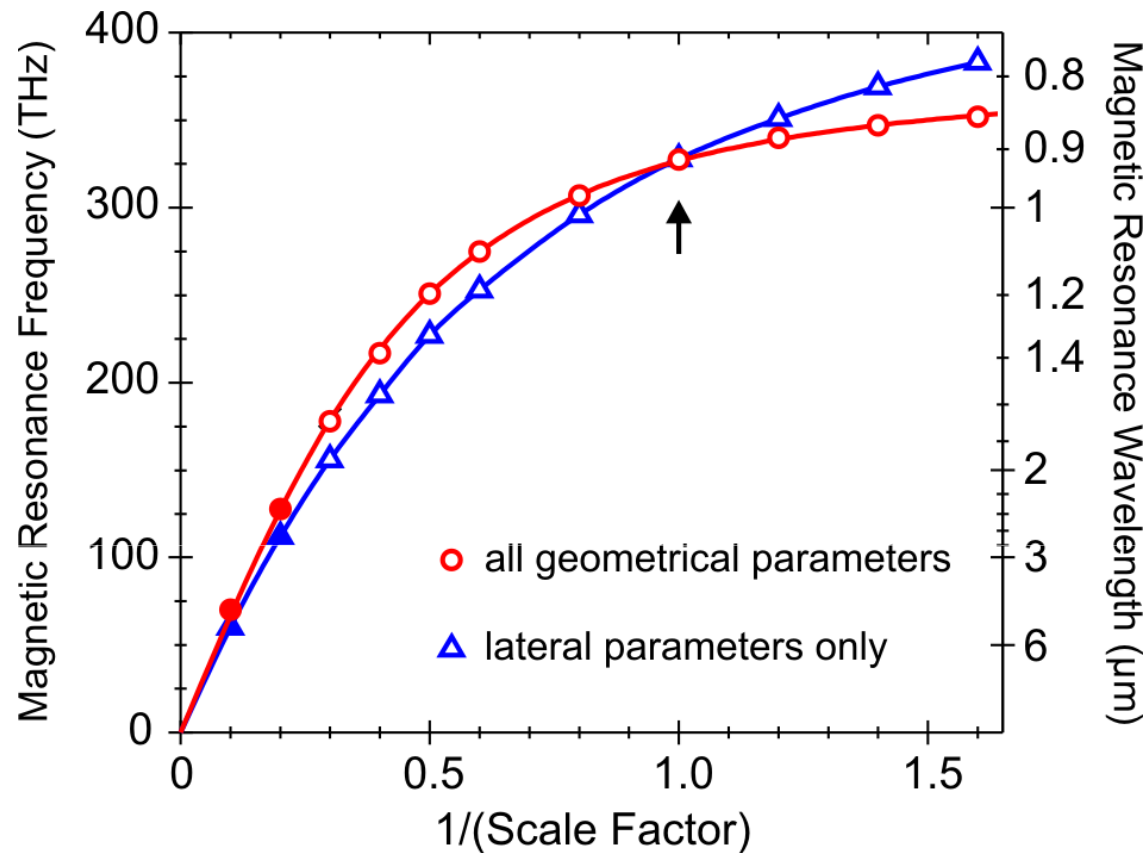
Scale:  $length \rightarrow S \times length \wedge time \rightarrow S \times time$   
Such that speed of light invariant and  $S \rightarrow 0$

$$C \propto S \wedge L \propto S \Rightarrow \omega_m = \frac{1}{\sqrt{LC}} \propto \frac{1}{S}$$

# Upper frequency limit of the SRRs?



# Limits of size scaling



$$\omega_{LC} = \frac{1}{\sqrt{LC}} \propto \frac{1}{\sqrt{size^2 + const}}$$

$E_{kin}$  of electrons!

# Why saturation of $\omega_m$ ?

$$\omega_m = \frac{1}{\sqrt{L_m C}}$$

$$E_m = \frac{1}{2} L_m I^2$$

$$L_m \propto a \quad C \propto a$$

$$\omega_m \propto 1/a \quad (a: \text{u.c. size})$$

**Key point:** Kinetic energy of the electrons becomes comparable to magnetic energy in small scale structures

$$E_e = \frac{1}{2} (n_e V) m_e v_e^2$$

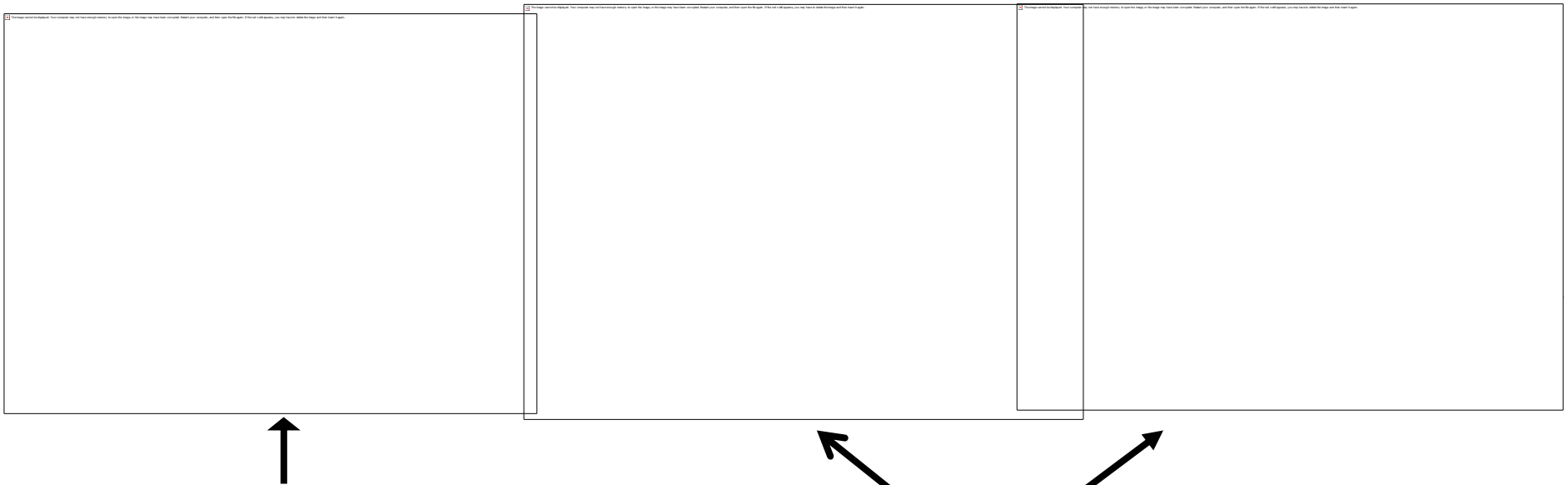
$$= \frac{1}{2} L_e I^2$$

$$v_e = \frac{I}{S e n_e} \quad L_e = \frac{m_e}{n_e e^2} \frac{V}{S^2} \sim \frac{1}{a}$$

$V$ : wire effective volume  
 $S$ : wire effective cross-section  
 $n_e$ :  $e^-$  number density

$$\omega_m = \frac{1}{\sqrt{(L_m + L_e)C}} = \frac{1}{\sqrt{c_1 a^2 + c_2}}$$

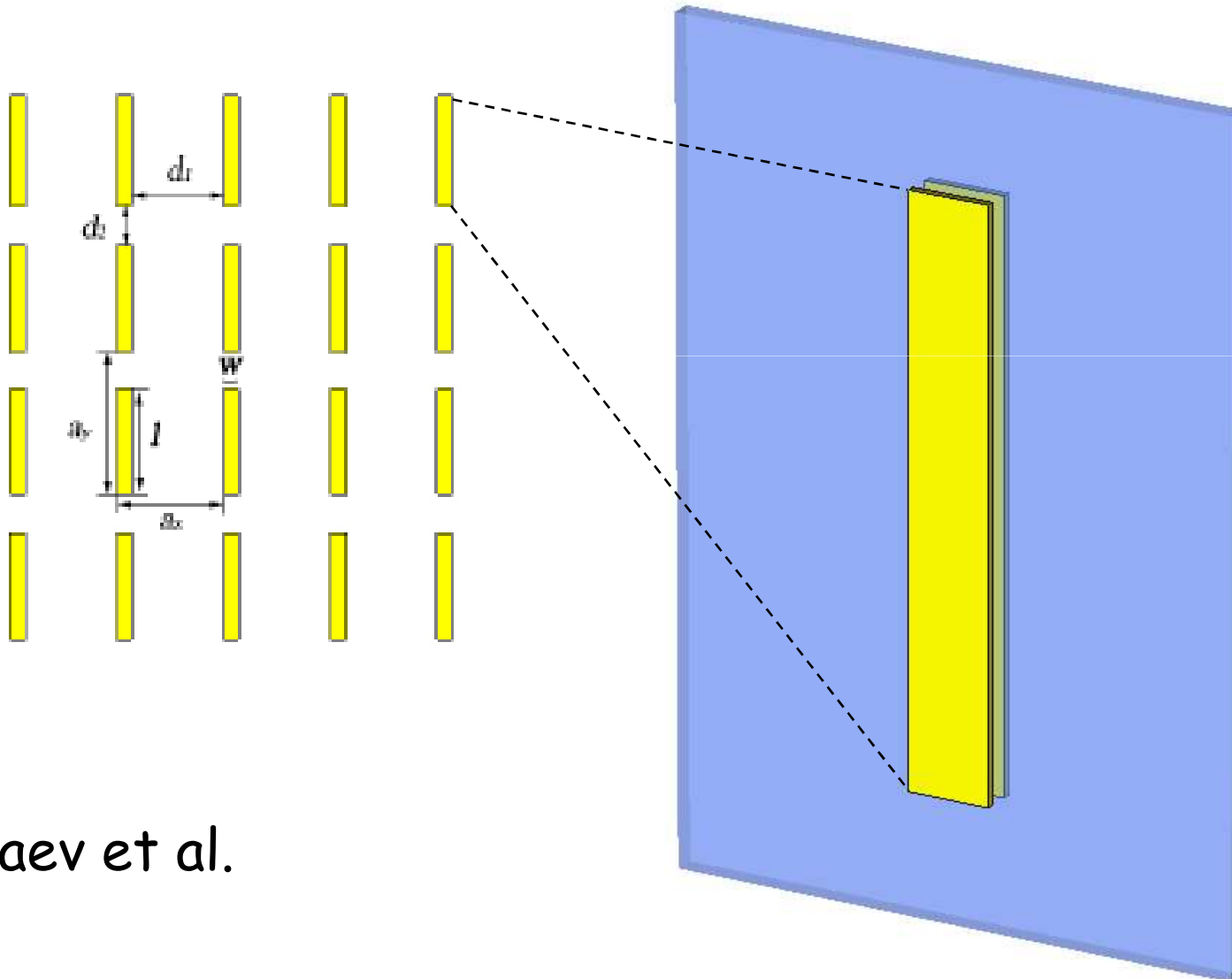
Metamaterials structures shown in these split-ring resonators (SRRs) are amenable to manufacture by common planar lithography.



It has a **magnetic response perpendicular** to the plane, which is **difficult to detect** by direct incidence measurements.

Use of multilayer processing can be used to fabricate metamaterials that give both **negative  $\epsilon$  and  $\mu$** , as well as  **$n$**  for **perpendicular** propagation.

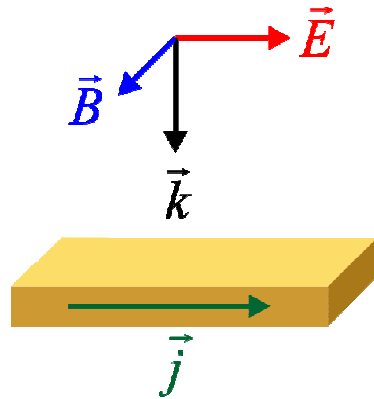
# LHM by Double Layer Cut - Wires



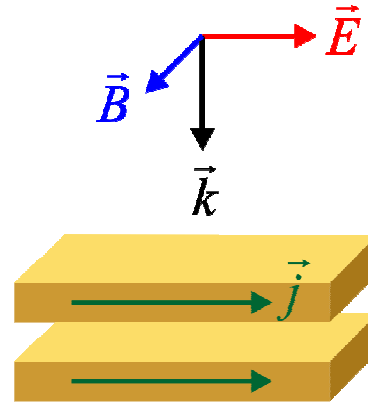
Shalaev et al.

# Magnetic response from cut-wire pairs

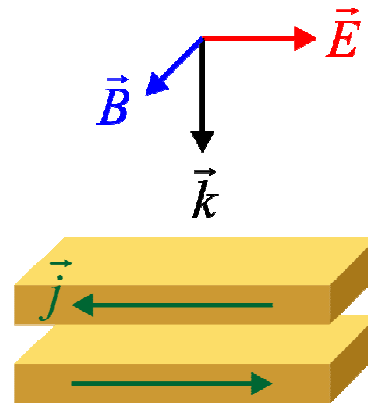
Single cut-wire



Cut-wire pair

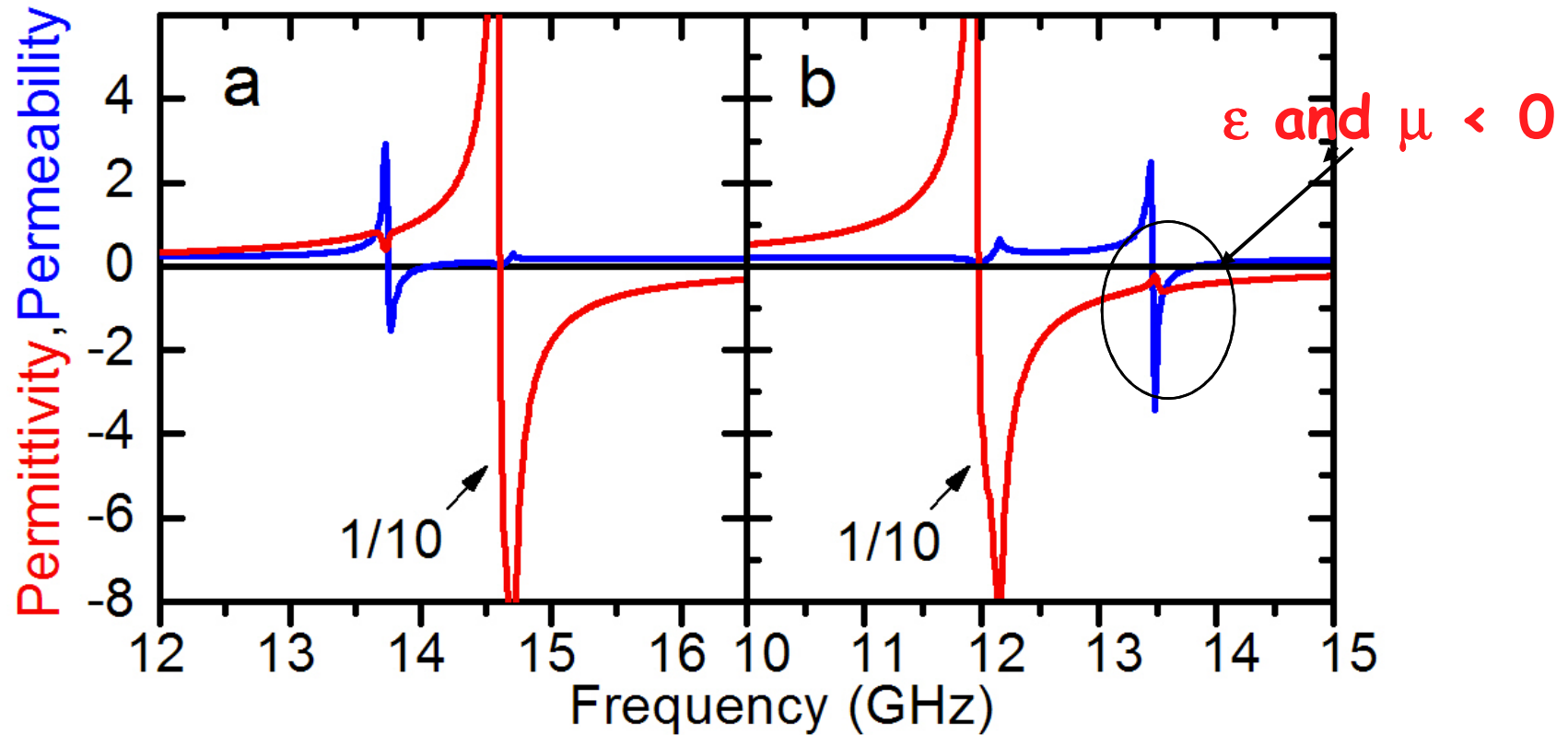


Symmetric mode,  
no magnetic dipole



Anti-symmetric mode,  
magnetic dipole

## Difficulties in obtaining both $\epsilon$ and $\mu < 0$ in cut wire pairs

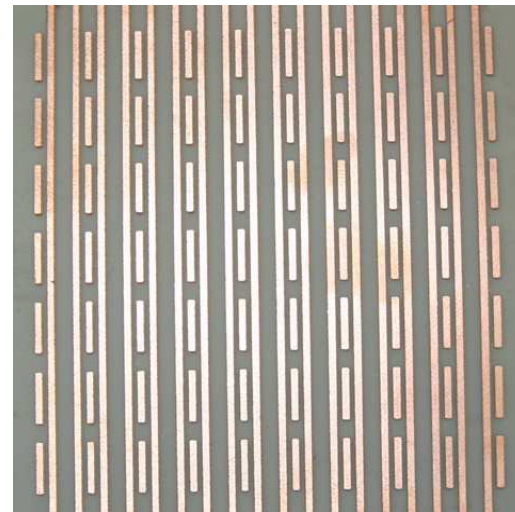
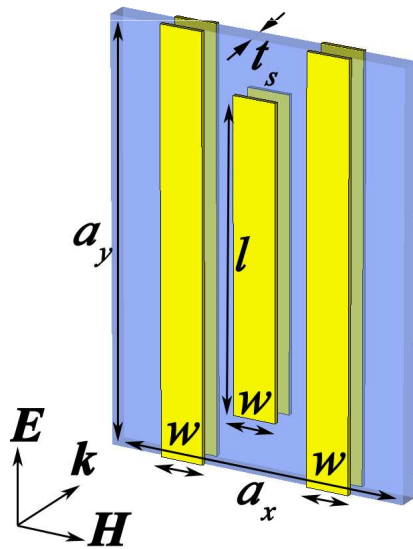
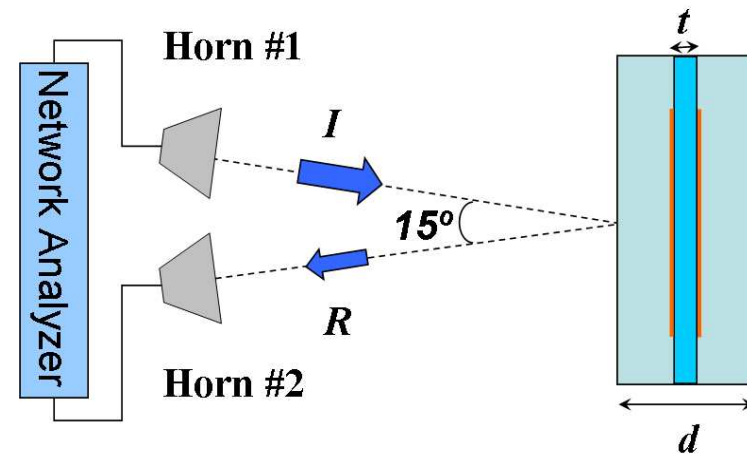
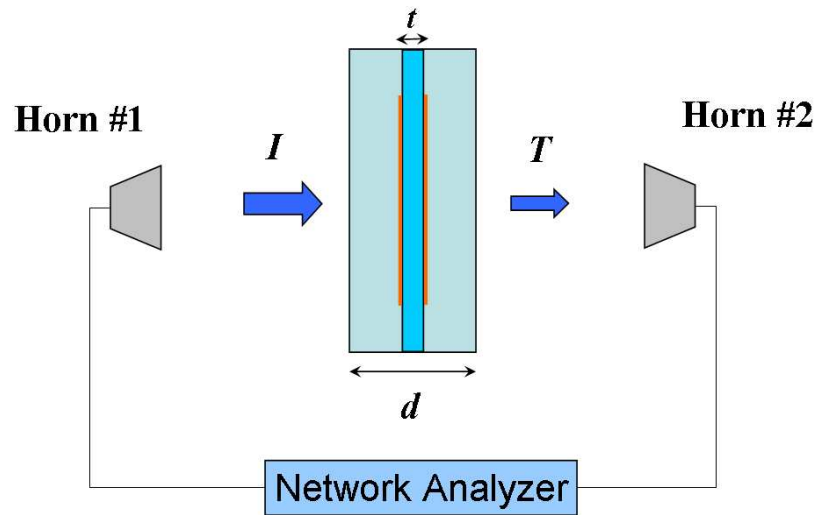


Strong electric resonance and weak magnetic resonance

To get  $\epsilon$  and  $\mu < 0$ , one needs to have  $\omega_e < \omega_m$

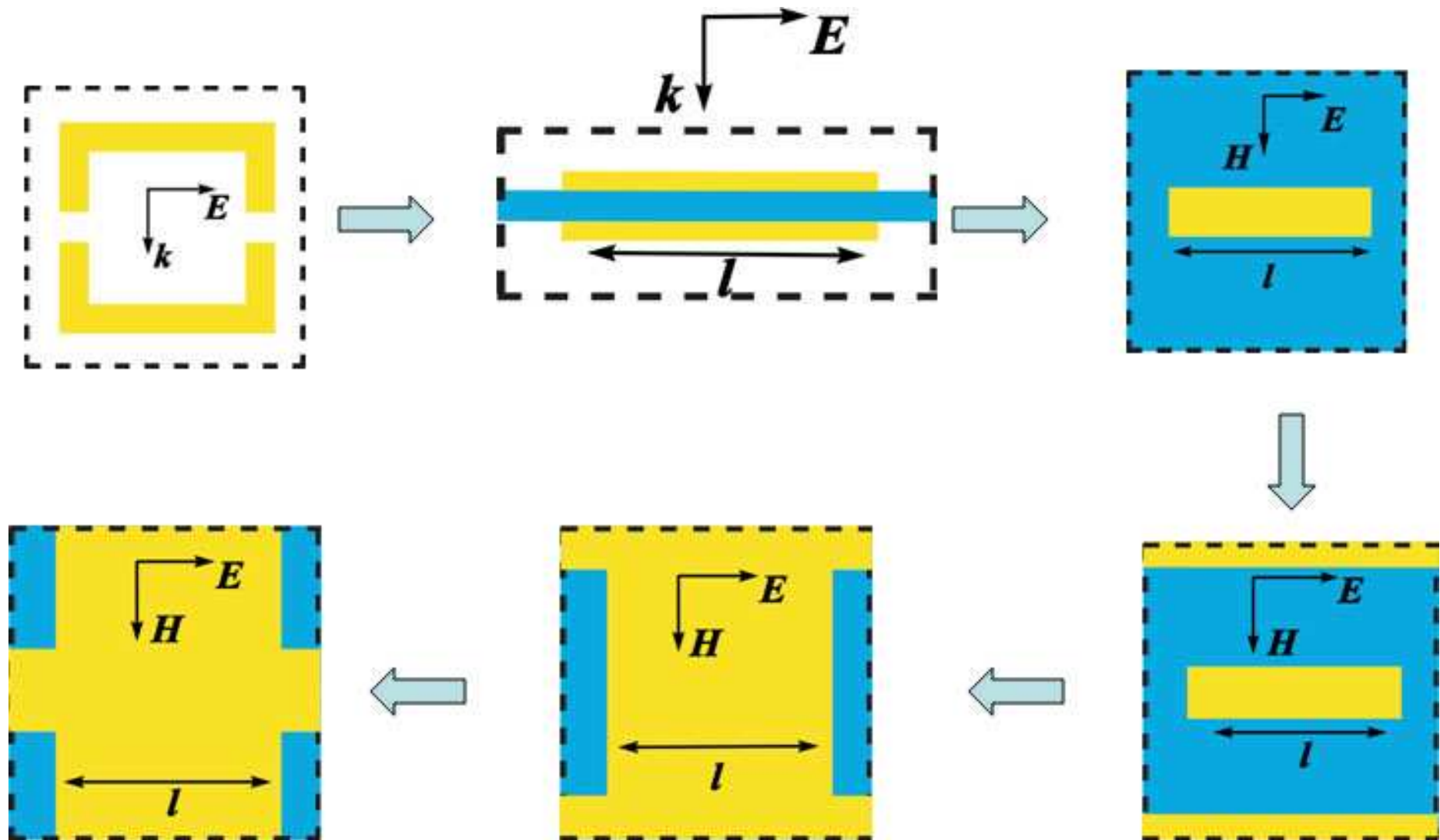
Resonance and anti-resonance effects in  $\epsilon$  and  $\mu$





PRB 73, 041101(R) (2006)

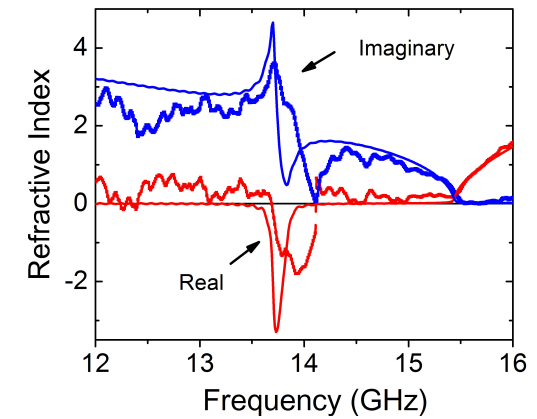
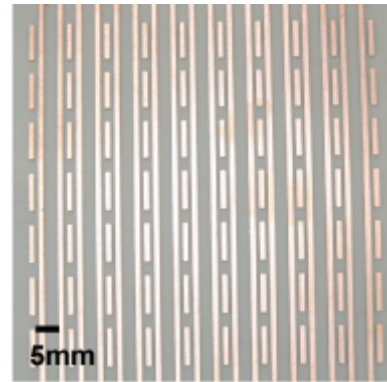
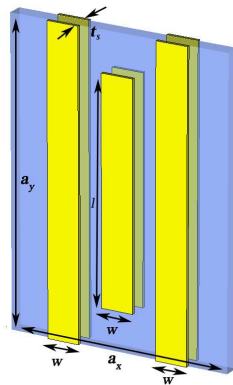
## Development of the 2-cut SRR to the fish-net structure



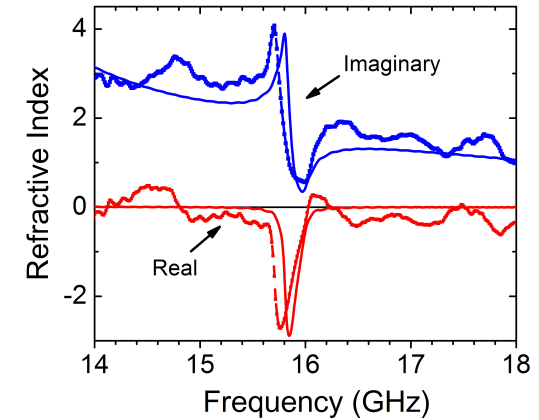
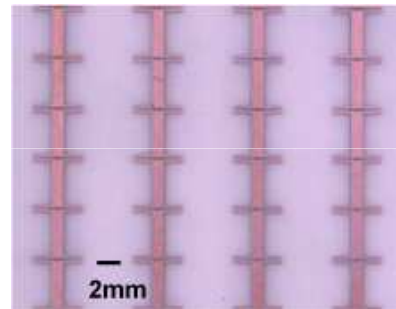
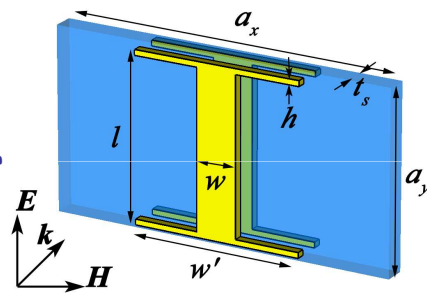
# Negative Index at GHz and THz: Short-wire pairs and Fishnet Structure

Short-wire pair  
+ Continuous wire

PRB **73**, 041101(R) (2006).

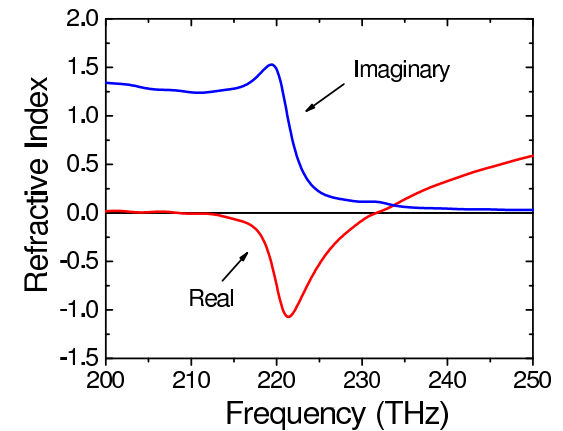
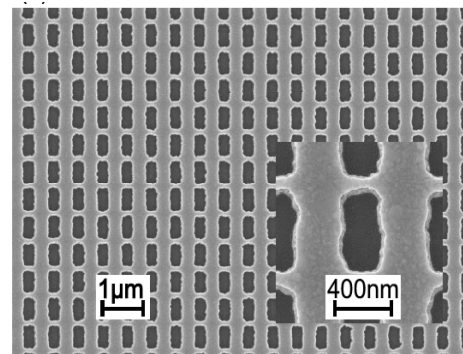
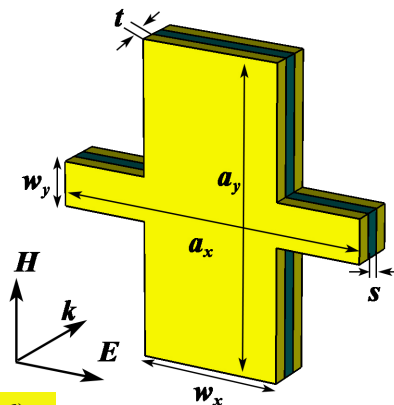


H-shaped  
Short-wire pair



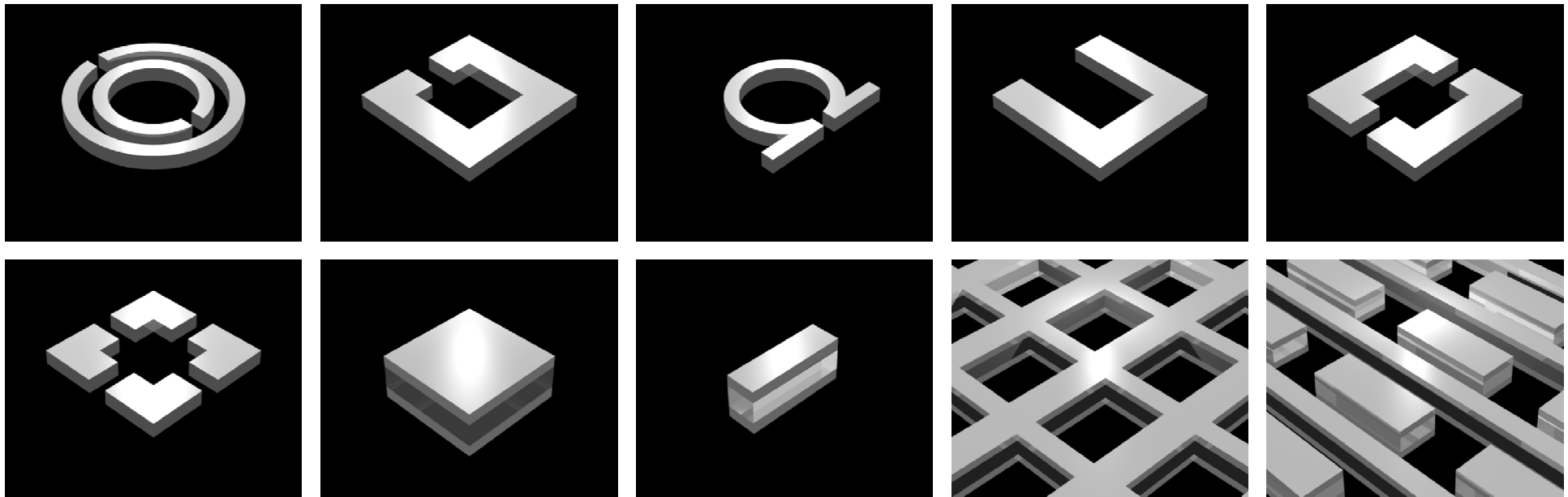
APL **88**, 221103 (2006).

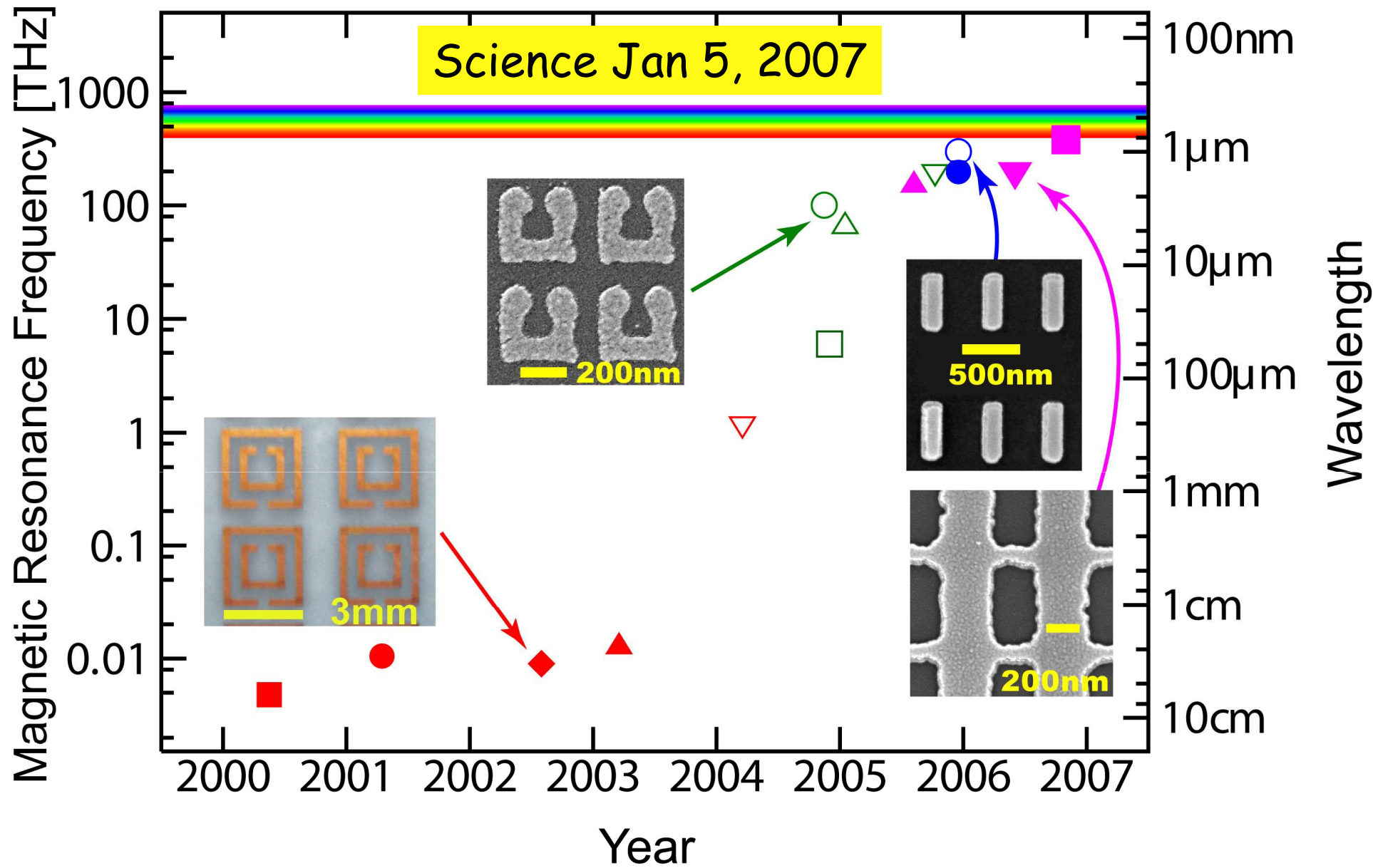
Fishnet



Science **312**, 5775 (2006)  
Opt. Lett. **32**, 53 (2007)

# Different designs used in fabricating LHMs with negative $\mu$ and $n$



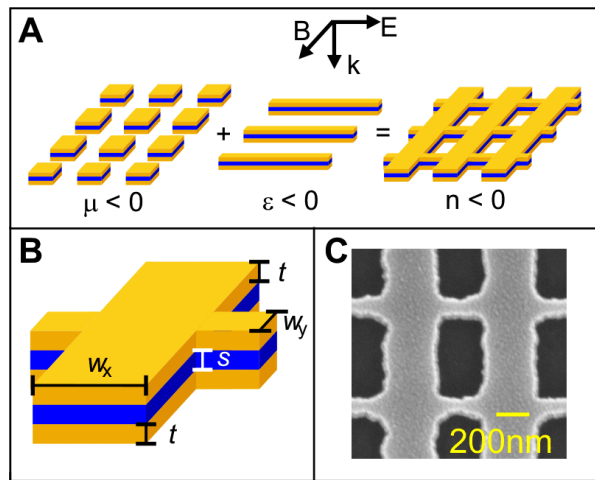


# Metamaterials Used to Alter Light's Path, Speed

Karlsruhe and Ames Lab., designed and fabricated for the first time NIMs at 1.5  $\mu\text{m}$  with low losses! The design is shown below.

Science 312, 892  
(2006)

"Reversing and accelerating the speed of light,"  
(<http://www.ameslab.gov/final/News/2006rel/metamaterials.htm>)



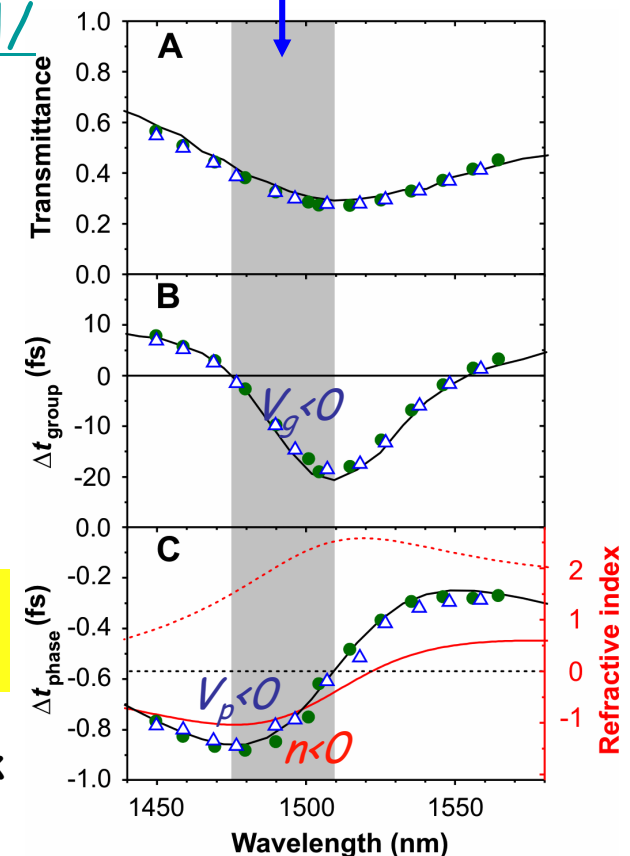
$$v_p = c / n$$

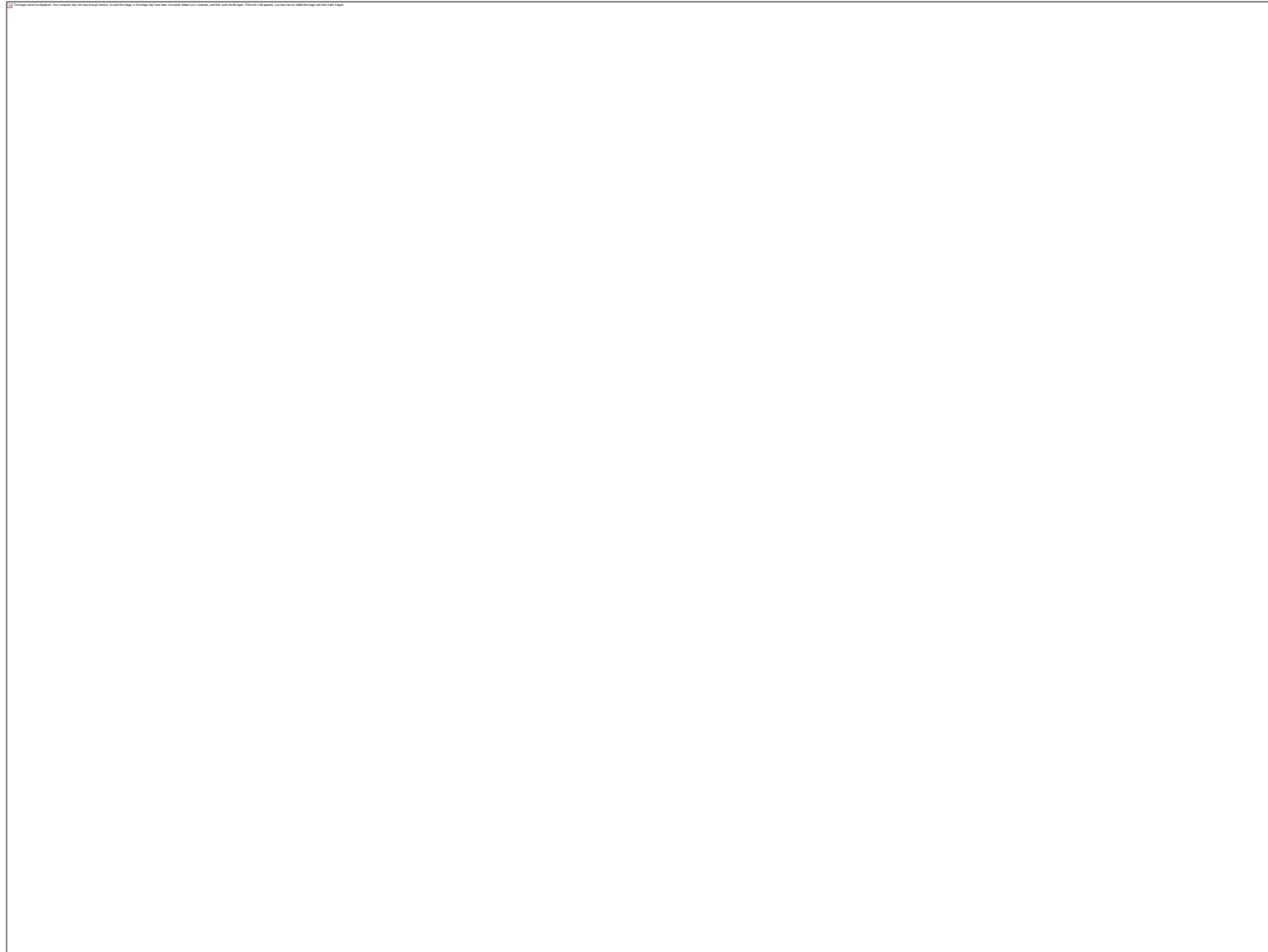
$$v_g = v_p / \left(1 + \frac{\omega}{n} \frac{dn}{d\omega}\right)$$

If  $n$  is negative and dispersive, both  $v_p$  and  $v_g$  can be negative!

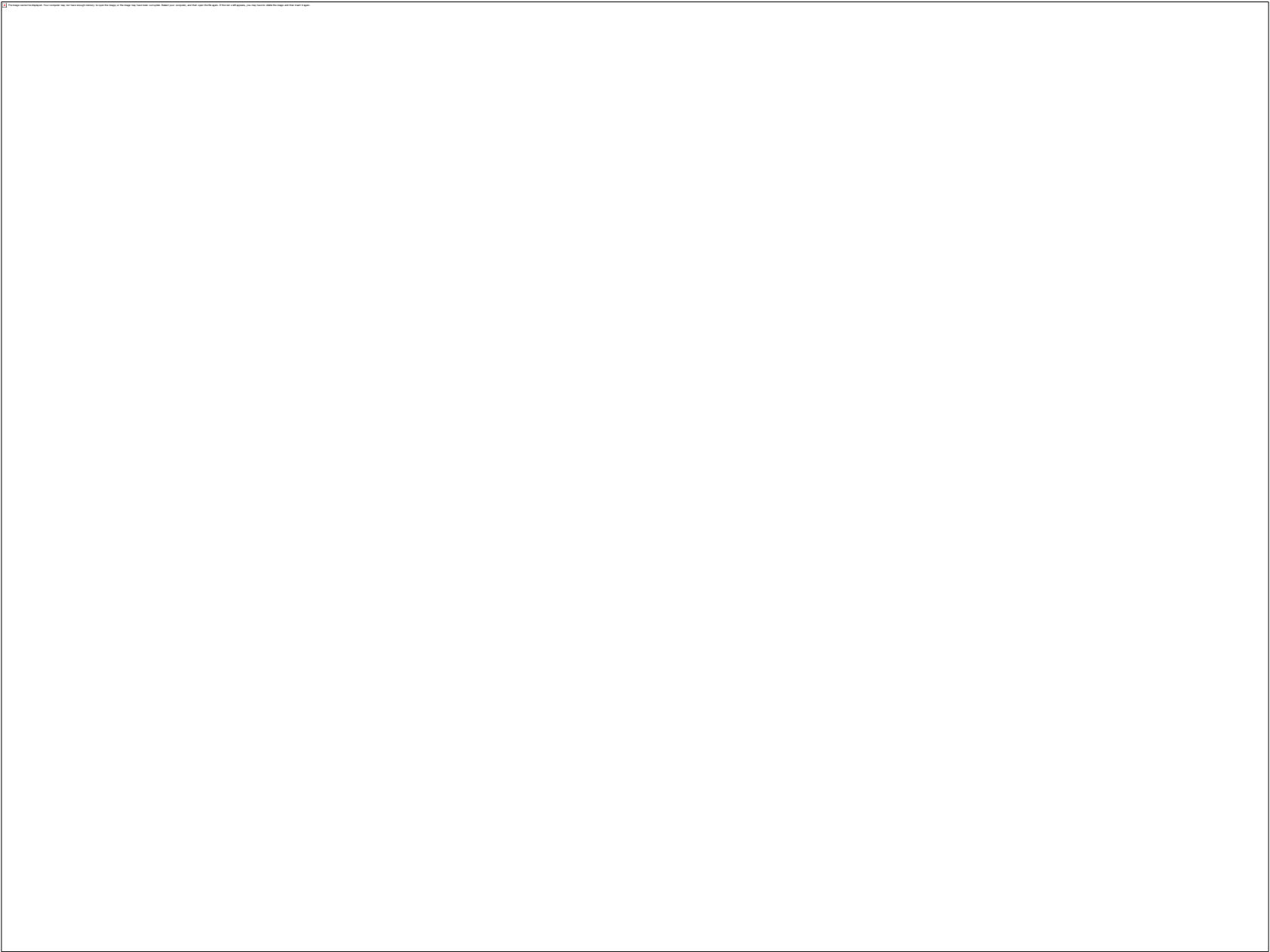
Causality and relativity are ok

Phase and group velocity can be both negative!  
 $n$  is dispersive in NIMs



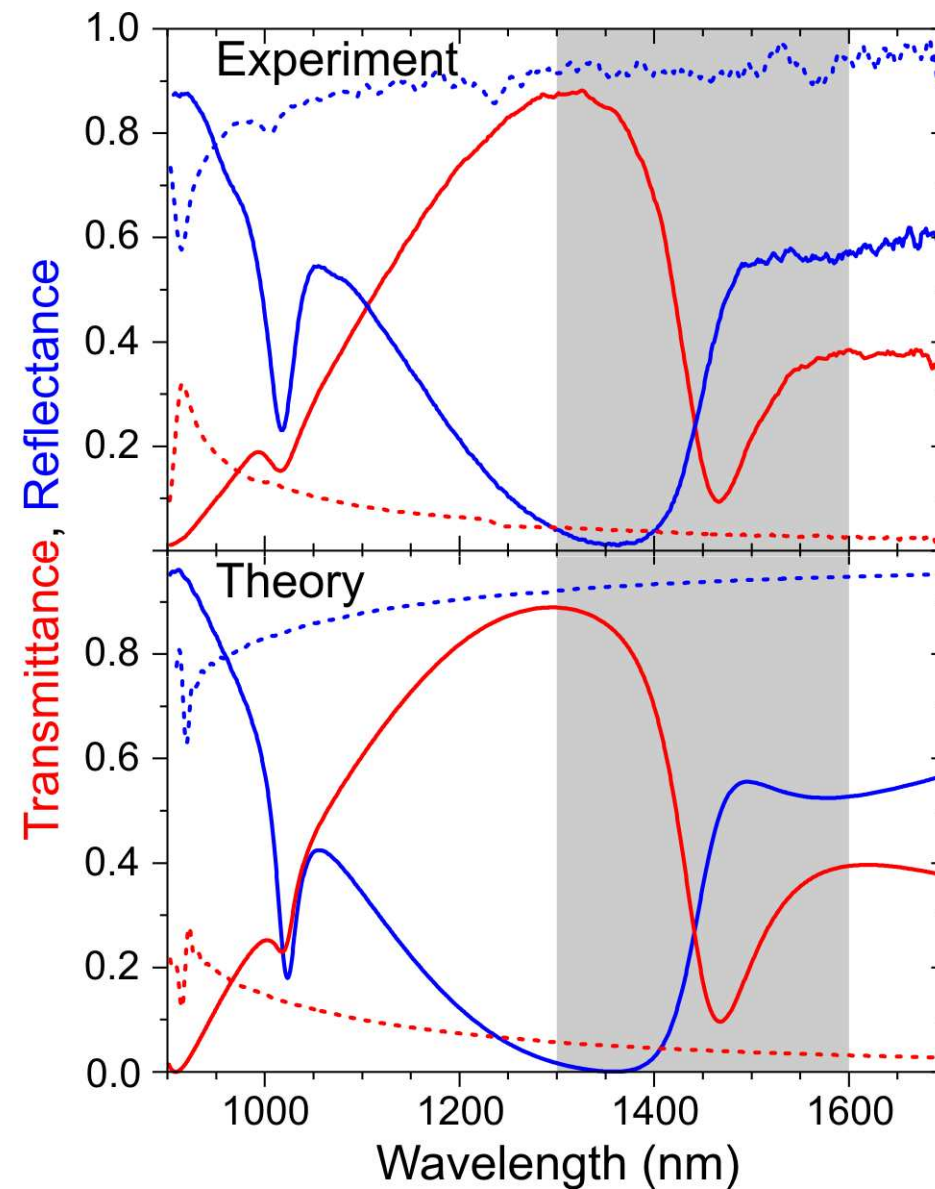
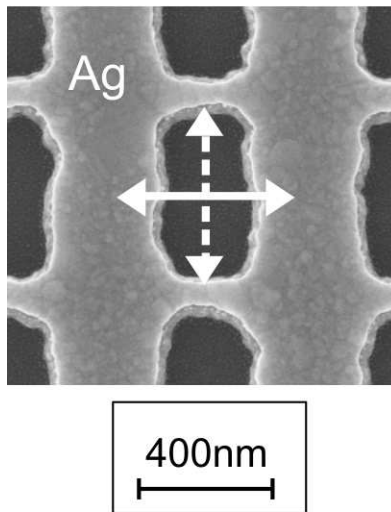


$$\varepsilon=\mu=n=-1$$



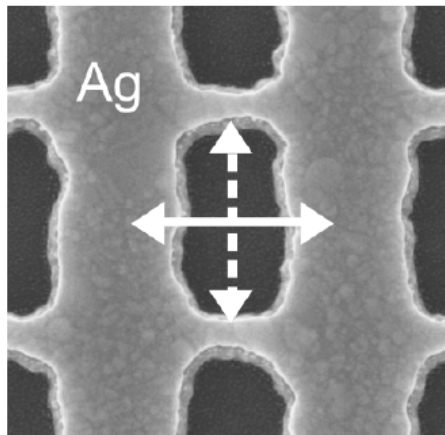


# Low loss negative index metamaterials



G. Dolling *et al.*, Opt. Lett. 31, 1800 (2006)

# Low loss negative index metamaterials

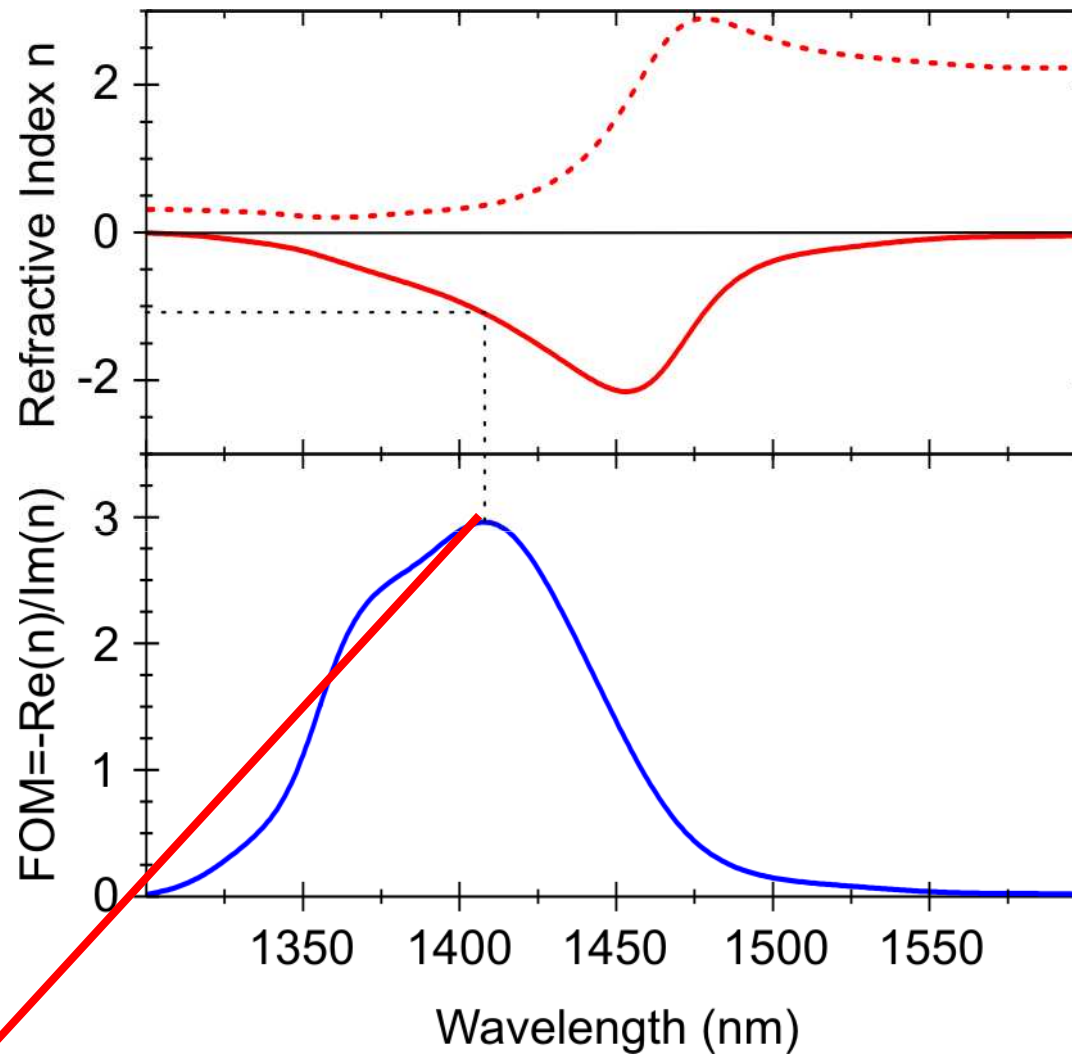


400nm

Decay per period

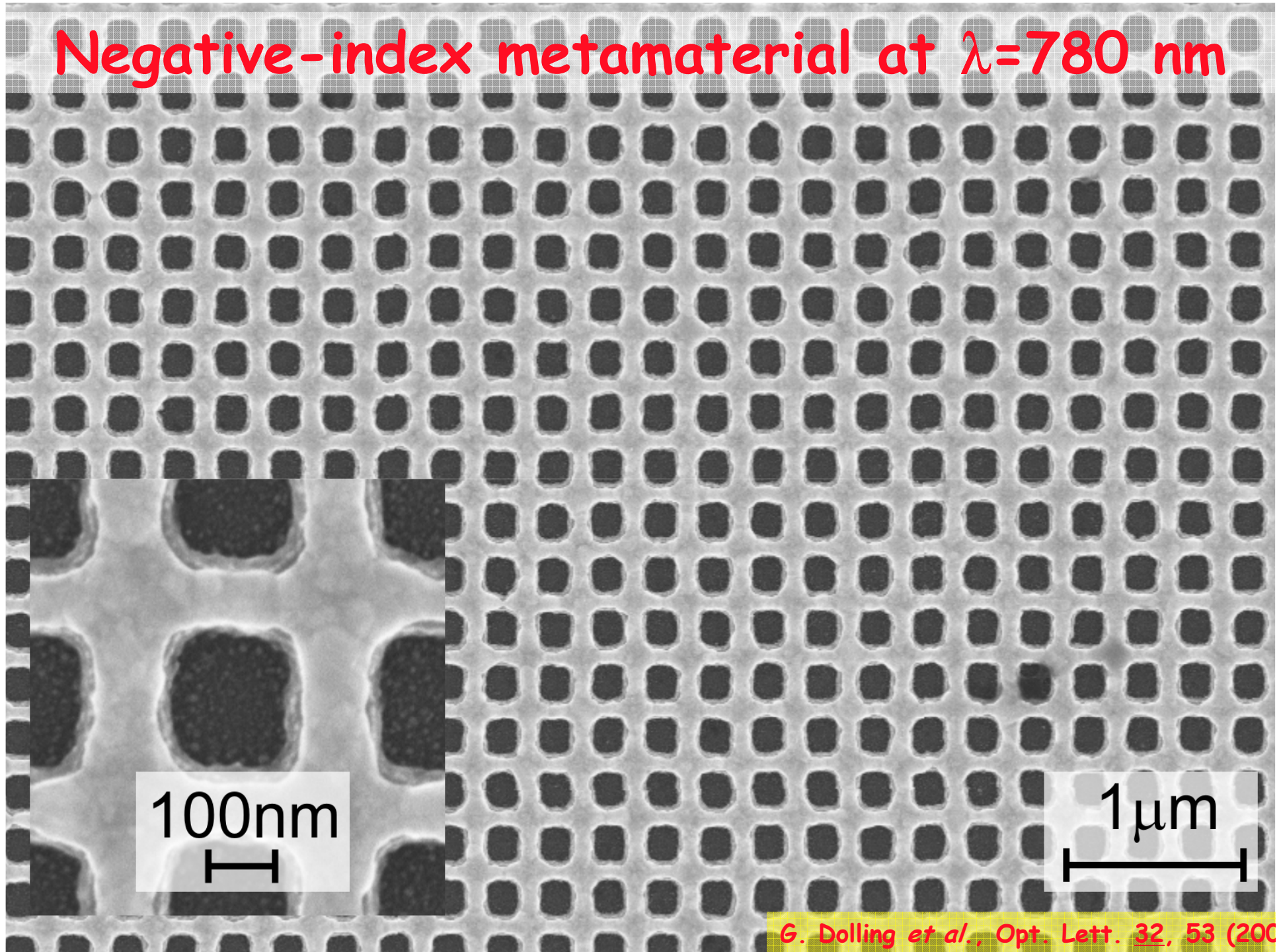


Best system has

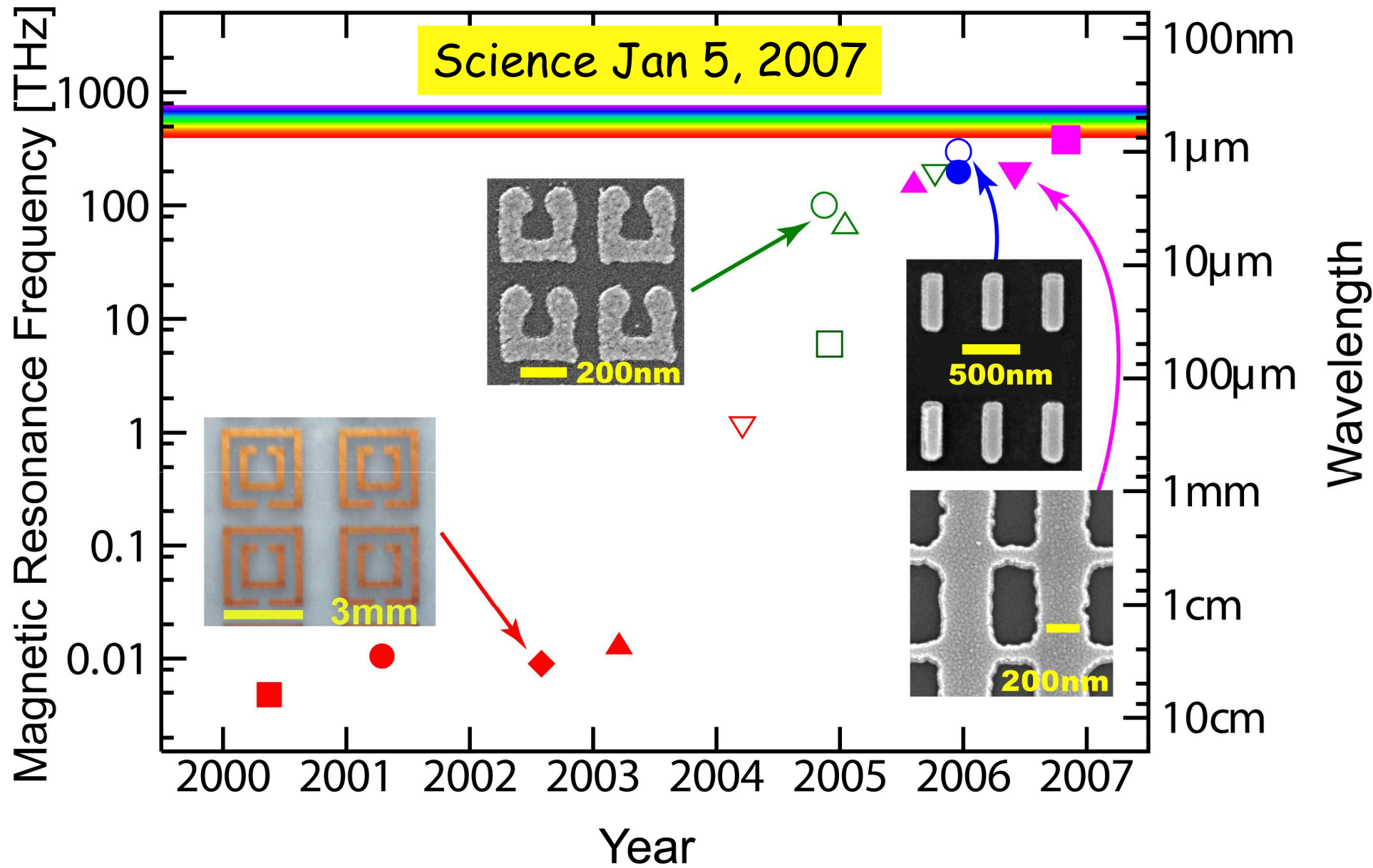


G. Dolling *et al.*, Opt. Lett. 31, 1800 (2006)

# Negative-index metamaterial at $\lambda=780$ nm







## Significant contributions to the development of LHMs by our group:

- Electric response of SRRs
- Electric excitation of the magnetic resonance
- Retrieval calculations for  $\epsilon$ ,  $\mu$
- Closed rings for distinguishing LH from RH peaks
- Negative  $\mu$  at THz and visible.
- Negative  $n$  at 1.5  $\mu$  and 780 nm.
- Upper frequency limit of the SRRs. Diamagnetic response of SRRs.
- Negative  $n$  at GHz and THz. Negative phase and group velocities.
- Designs for **3d isotropic LHMs**.

## Future directions:

- Understanding and reducing losses. Introduce gain to reduce losses.
- Fabrication of 3d LHMs. Direct laser writing. (Karlsruhe)
- Electromagnetic induced transparency. Slow light, low losses.
- Non-linear effects. Chirality effects.
- Anisotropic metamaterials. Pseudo-focusing.
- Applications

**\$\$\$ DOE, DARPA, MURI, NATO, EU, EU-PHOMÉ**

# Conclusions

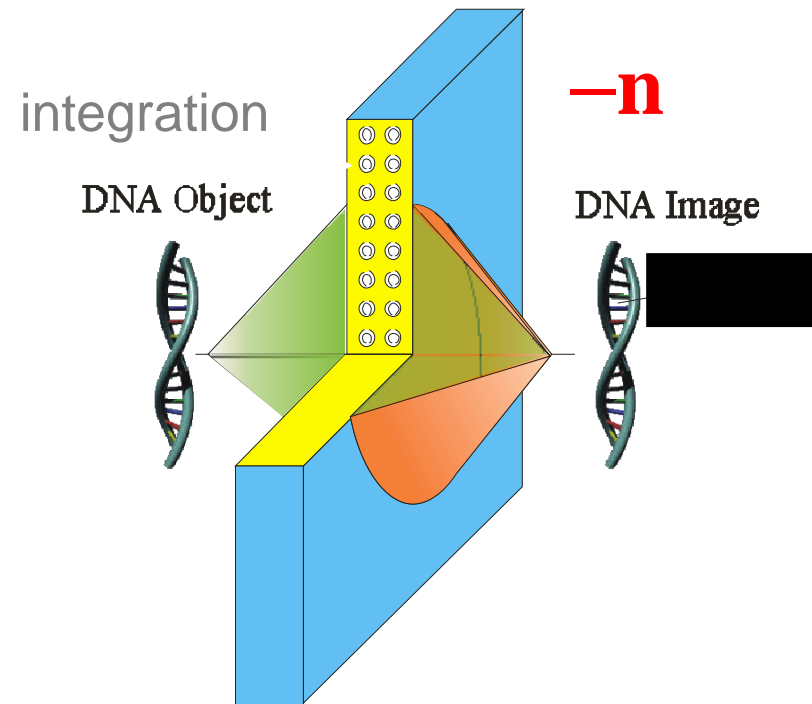
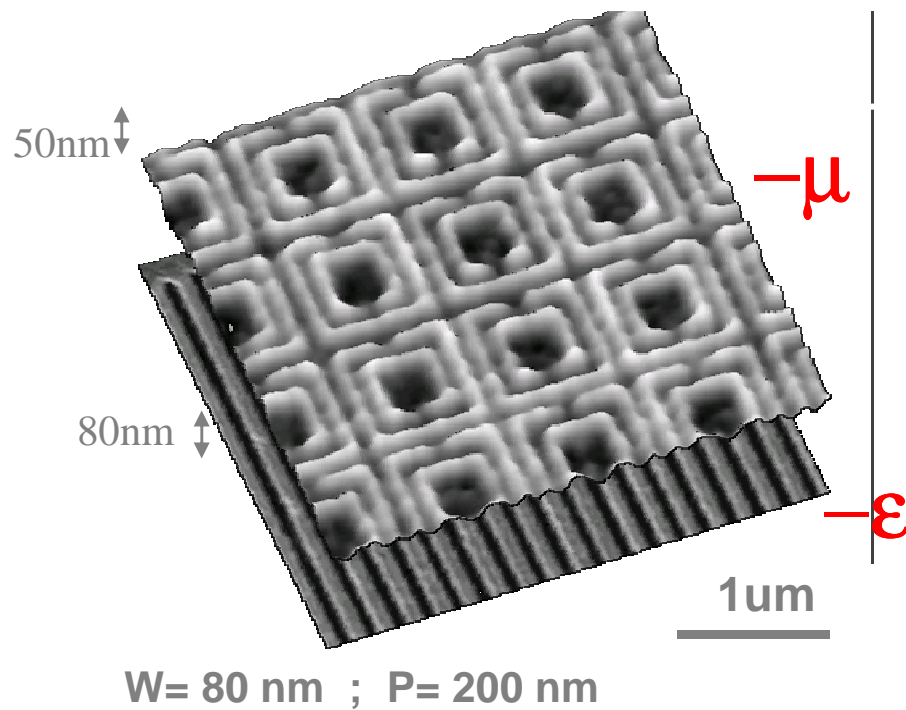
- Our team has been instrumental in creating and developing a new revolutionary field, which extends the realm of electromagnetism and opens up exciting technological applications from the MHz range to optical frequencies
- The realization of negative index materials has opened up the possibility of unprecedented applications and devices.
  - MHz: Artificial magnetic materials for MRI applications
  - GHz: Cellular communications
    - Miniaturized antennas and waveguides
  - Optics: Superlenses with subwavelength imaging

# Nano Plasmonics at Near-IR and Visible

Nanowires and Nano-Rings as  
resonating elements

W= 75 nm ; G= **50 nm** *Ultra Optical Imaging?*

L= 667 nm ; P=717 nm



# Conclusions

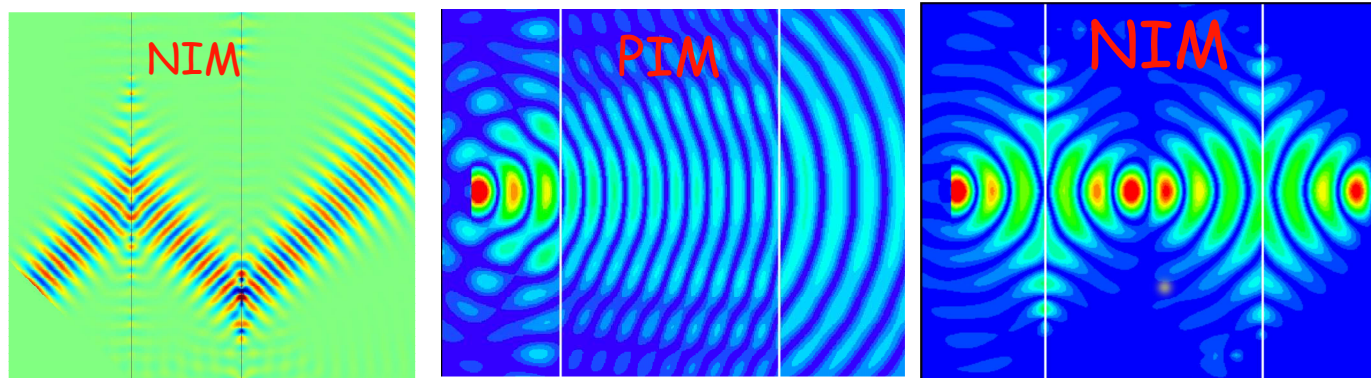
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- The realization of negative index materials has opened up the possibility of unprecedented applications and devices.
  - MHz: Artificial magnetic materials for MRI applications
  - GHz: Cellular communications
    - Miniaturized antennas and waveguides
  - Optics: Superlenses with subwavelength imaging
    - 10 nm VLSI nanolithography using optics
    - smaller integrated circuits
    - Molecular Imaging (Medicine, Biology)
  - Optics: DVDs with 100x capacity

These applications are just a start and more inventions will come from hundreds of research groups working on the newly created area of metamaterials



# Some reviews articles from our group

- 1) Bending Back Light: The Science of Negative Index Materials  
*Optics and Photonics News*, June 2006
- 2) Negative index materials: New frontiers in optics  
*Adv. Mater.* 18, 1941 (2006)
- 3) Photonic metamaterials: Magnetism at optical frequencies,  
*IEEE J. of Selected Topics in Quant. Electr.* 12, 1097 (2006)
- 4) Negative Refractive Index at Optical Wavelengths  
*Science* 315, 47 (2007)



Functional cloaking device !

