

# AntiHydrogen Trapped

Francis Robicheaux  
Auburn University

Collaborators: ALPHA

Jim Hanson, **Turker Topcu**, **Michael Wall**, **Chris Norton**, **Christine Taylor**, **Michele Zhang**, **Jennifer Hurt**, **Patrick Carpenter**, **Patrick Donnan**

Support from DOE and NSF

# Goal

Why are we excited about holding onto atomic antimatter?

Why is it hard to do?

# P = Parity

Laws of physics should(?) be same under the operation:

$$P : \begin{pmatrix} x \\ y \\ z \end{pmatrix} \mapsto - \begin{pmatrix} x \\ y \\ z \end{pmatrix}$$

Scalar  $a \rightarrow a$

$$\text{Example: } P(\vec{r} \bullet \vec{r}) = (-\vec{r}) \bullet (-\vec{r}) = \vec{r} \bullet \vec{r}$$

Vector  $\mathbf{r} \rightarrow -\mathbf{r}$

Axial Vector  $\mathbf{C} \rightarrow \mathbf{C}$

$$\text{Example: } P(\vec{r} \times \vec{p}) = (-\vec{r}) \times (-\vec{p}) = \vec{r} \times \vec{p}$$

# Classical Physics Under P

Newton's Equations: Unchanged

$$P \vec{F} = - \vec{F} = P (M \vec{a}) = - M \vec{a}$$

Maxwell's Equations: Unchanged

Vectors:  $\mathbf{E}$ ,  $\mathbf{J}$

Axial Vector:  $\mathbf{B}$

Scalar:  $\rho_t$

$$\nabla \cdot \mathbf{E} = \frac{\rho_t}{\epsilon_0}$$

$$\nabla \cdot \mathbf{B} = 0$$

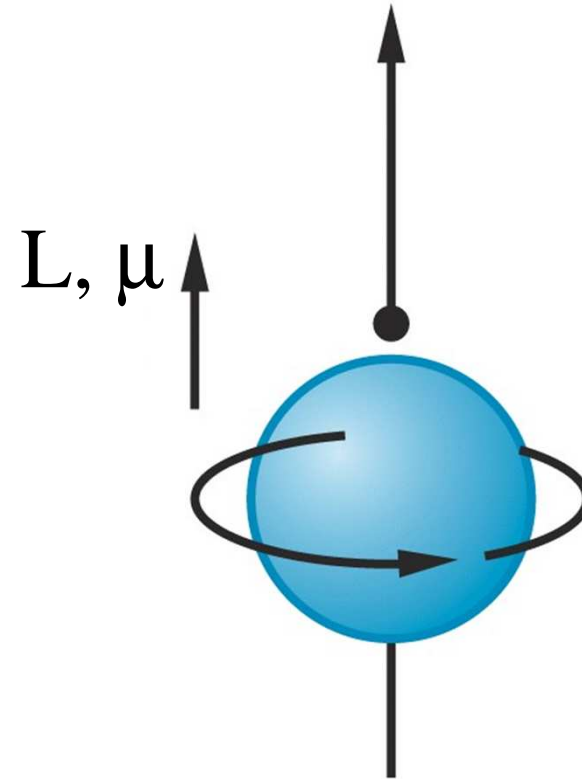
$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t}$$

$$\nabla \times \mathbf{B} = \mu_0 \mathbf{J} + \mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t}$$

# P Violation

Lee and Yang gave arguments for possibility of parity violation in weak decays

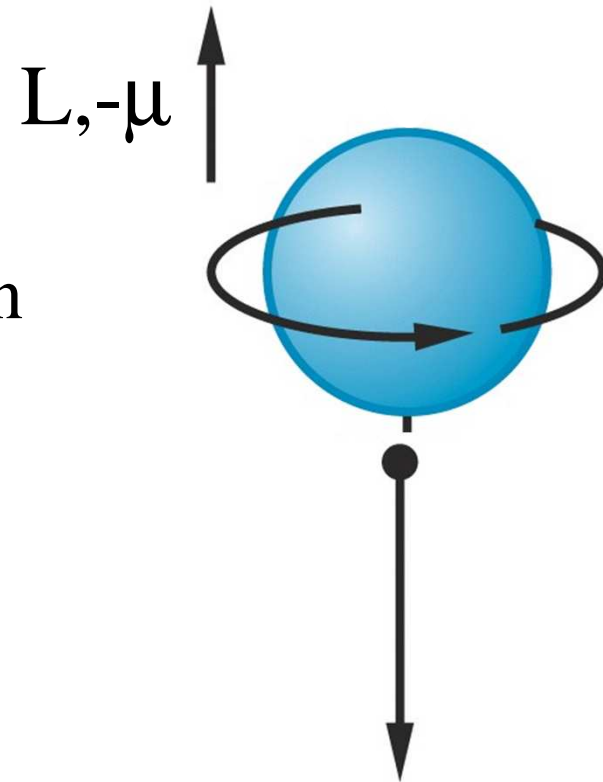
C.S. Wu found signal in  $\beta$  decay of  $^{60}\text{Co}$  by orienting magnetic moment in large magnetic field and low temperature.



# CP OK!

C = Charge conjugation  
transformation ( $Q \rightarrow -Q$ )

Charge conjugation transformation  
and parity inversion works for  
 $^{60}\text{Co}$ .

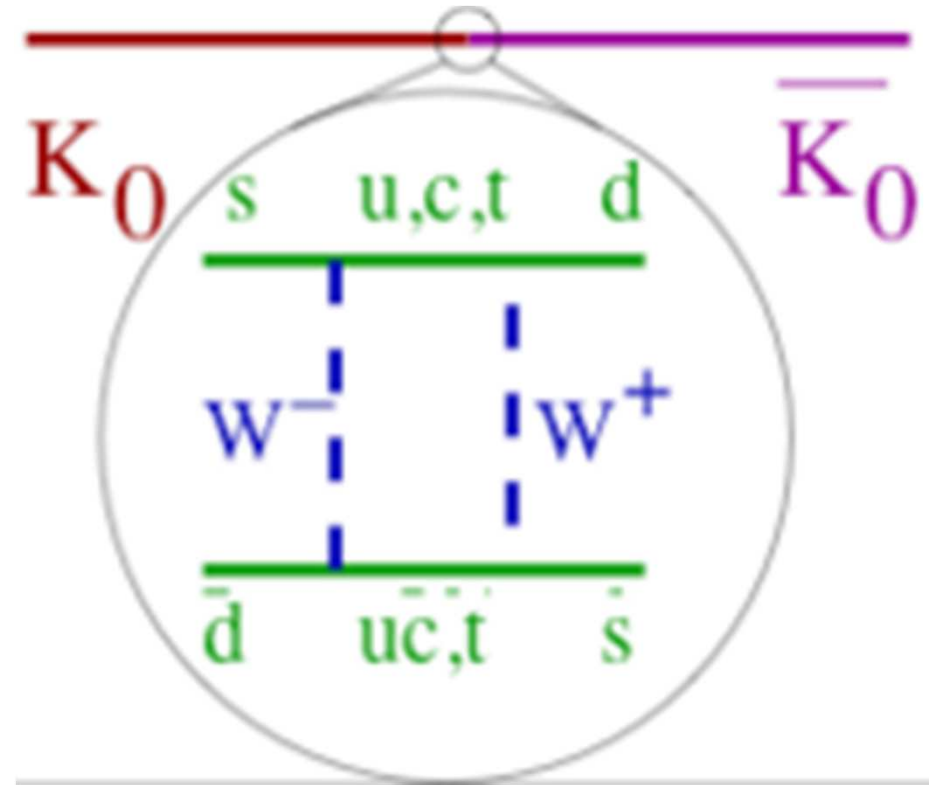


# CP not OK!

$K_0$  and its antiparticle mix through weak interaction to make  $K_L$  and  $K_S$ .

$K_L$  always decays into  $3 \pi$   
 $CP = -1$

$K_S$  always decays into  $2 \pi$   
 $CP = +1$



$K_L$  sometimes decays into  $2 \pi$ . This is different CP symmetry. One phase in Standard Model. (2008 Nobel)

# CPT OK?



Any **local** quantum field theory,  
obeying **Lorentz invariance** and  
usual **spin-statistics** connection



**CPT Invariance**

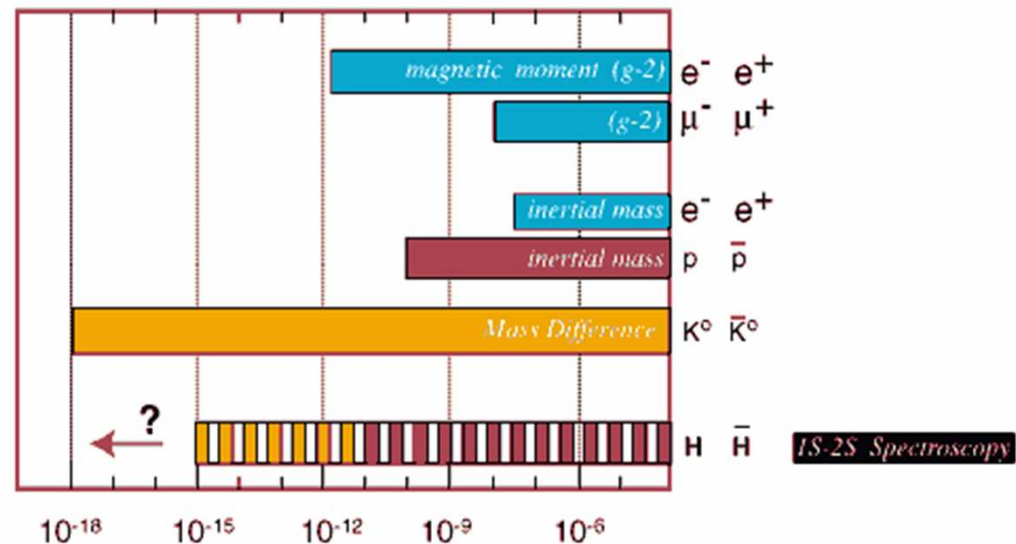
G. Lüders, Ann. Phys. 2, 1-15 (1957)  
(also: W. Pauli, J. Schwinger)

**MATTER = ANTIMATTER** at **any** level of precision

T is the  
time  
reversal  
operation.

(note the logarithmic scale)

## The most precise CPT Tests



# Anti-hydrogen

Hydrogen atom = 1 proton & 1 electron

Anti-hydrogen atom = 1 anti-proton & 1 positron

**Don't let the anti-matter touch matter.**

Positron finds an electron, they disappear and gamma rays come out.

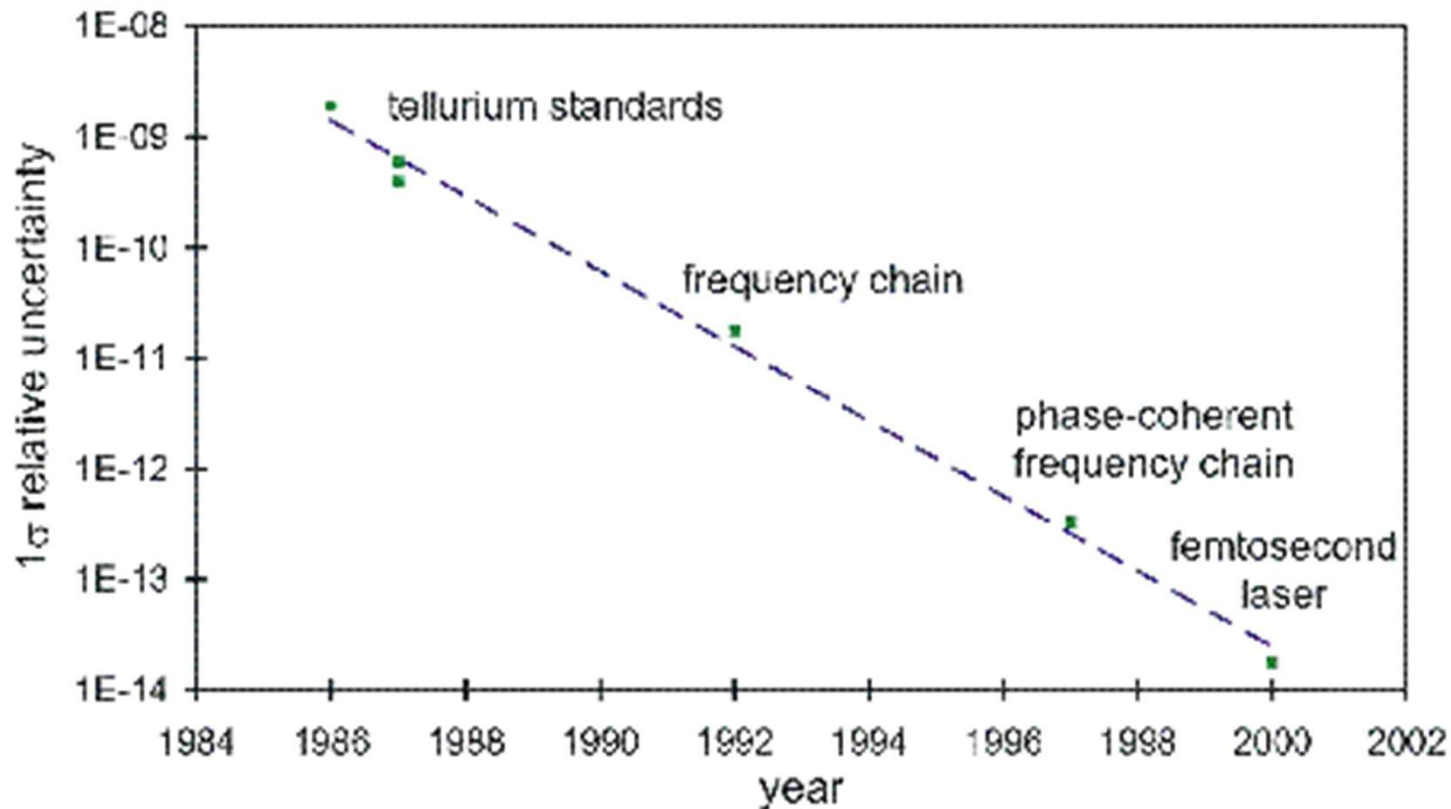
Antiproton finds a nucleus, antiproton “disappears” and a bunch of stuff comes out.

If you have a good vacuum, you could hold antihydrogen for a loooooooooooooooooong time.

# H Spectroscopy

H is excited from 1s to 2s by 2 photons with  $\lambda \sim 243$  nm

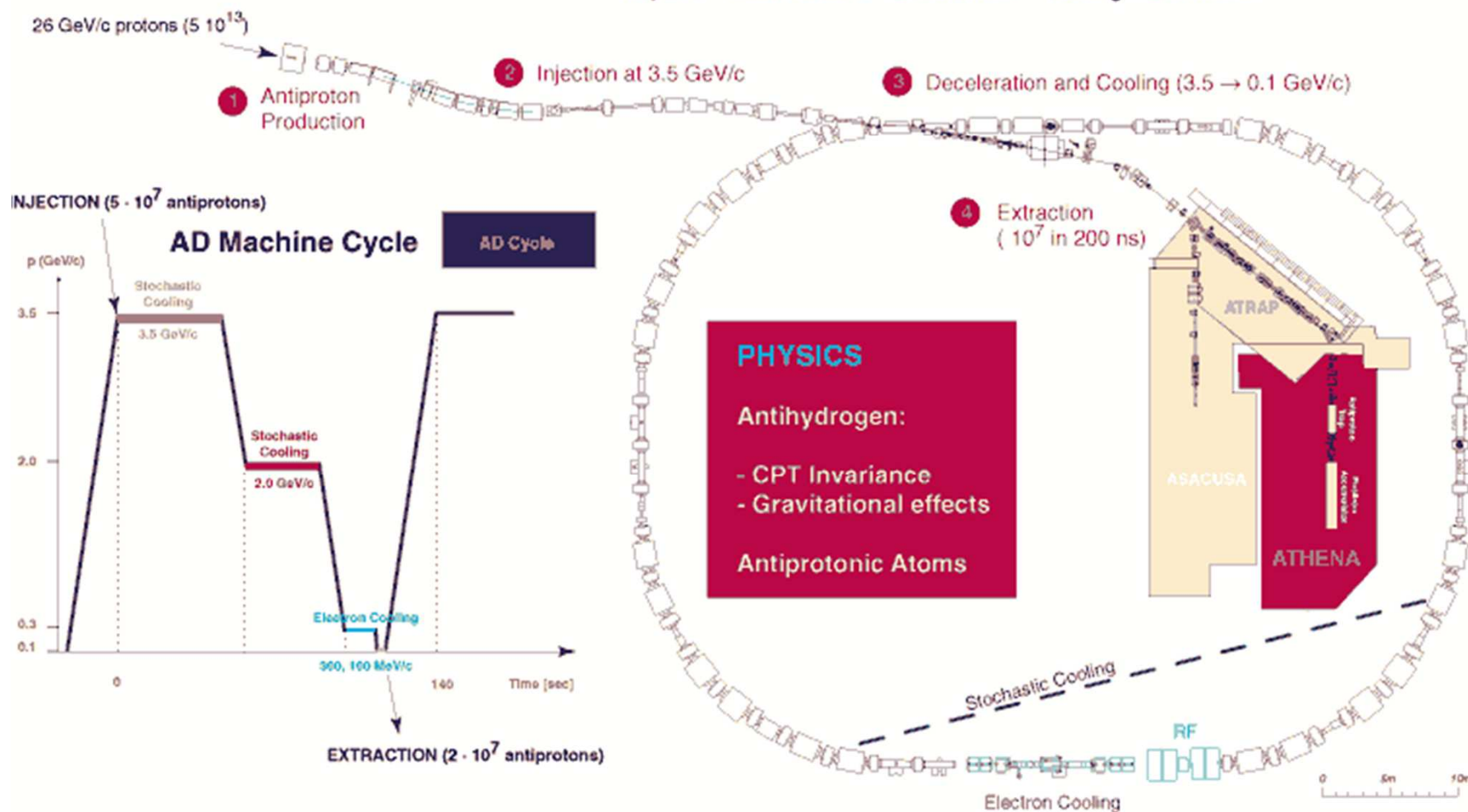
$$f(1S-2S) = 2\,466\,061\,413\,187\,074(34) \text{ Hz}$$



# Antiproton Decelerator

AD replaces (AC + AA + PS + LEAR)

Capture Accumulation Deceleration Cooling + Extraction



# ALPHA Collaboration

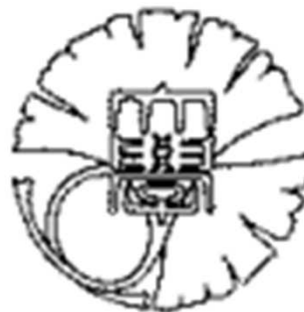
<http://alpha.web.cern.ch/alpha>



THE UNIVERSITY  
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UNIVERSITY  
*Sesquicentennial*



# Making Antihydrogen

Movie

Sept 18, 2002

## Production and detection of cold antihydrogen atoms

M. Amoretti\*, C. Amsler†, G. Bonomi‡§, A. Bouchta‡, P. Bowe||, C. Carraro\*, C. L. Cesar¶, M. Charlton#, M. J. T. Collier#, M. Doser‡, V. Filippini☆, K. S. Fine‡, A. Fontana☆☆, M. C. Fujiwara††, R. Funakoshi††, P. Genova☆☆, J. S. Hangst||, R. S. Hayano††, M. H. Holzschneider‡, L. V. Jørgensen#, V. Lagomarsino\*‡‡, R. Landua‡, D. Lindelöf†, E. Lodi Rizzini§☆, M. Macri\*, N. Madsen†, G. Manuzio\*‡‡, M. Marchesotti☆, P. Montagna☆☆, H. Pruys†, C. Regenfus†, P. Riedler‡, J. Rochet†#, A. Rotondi☆☆, G. Rouleau‡#, G. Testera\*, A. Variola\*, T. L. Watson# & D. P. van der Werf#

\* Istituto Nazionale di Fisica Nucleare, Sezione di Genova, and ‡‡ Dipartimento di Fisica, Università di Genova, 16146 Genova, Italy

† Physik-Institut, Zürich University, CH-8057 Zürich, Switzerland

‡ EP Division, CERN, CH-1211 Geneva 23, Switzerland

§ Dipartimento di Chimica e Fisica per l'Ingegneria e per i Materiali, Università di Brescia, 25123 Brescia, Italy

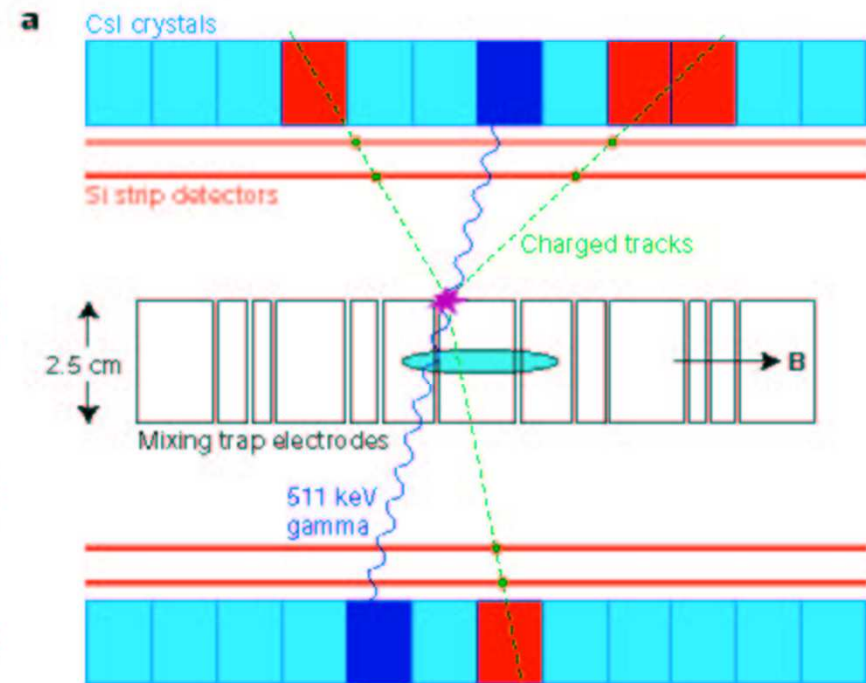
|| Department of Physics and Astronomy, University of Aarhus, DK-8000 Aarhus C, Denmark

¶ Instituto de Física, Universidade Federal do Rio de Janeiro, Rio de Janeiro 21945-970, and Centro Federal de Educação Tecnológica do Ceará, Fortaleza 60040-531, Brazil

# Department of Physics, University of Wales Swansea, Swansea SA2 8PP, UK

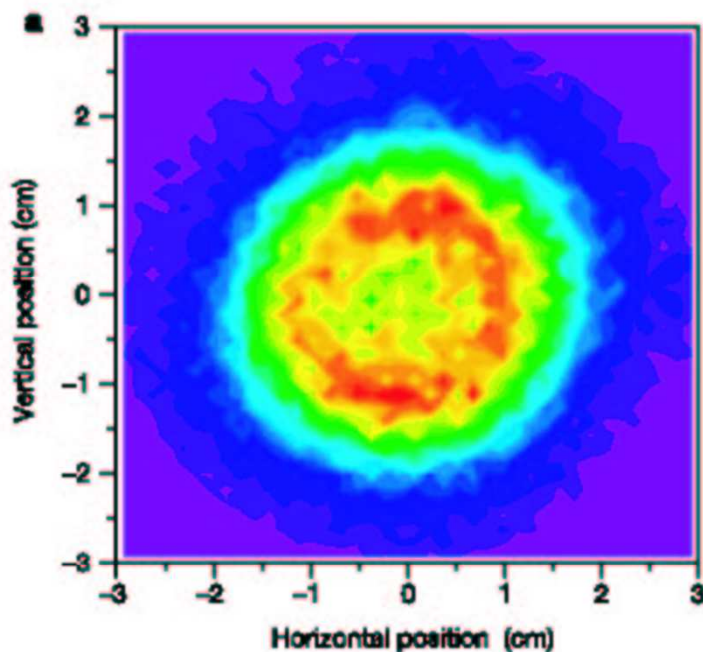
☆ Istituto Nazionale di Fisica Nucleare, Sezione di Pavia, and \*\* Dipartimento di Fisica Nucleare e Teorica, Università di Pavia, 27100 Pavia, Italy

†† Department of Physics, University of Tokyo, Tokyo 113-0033, Japan

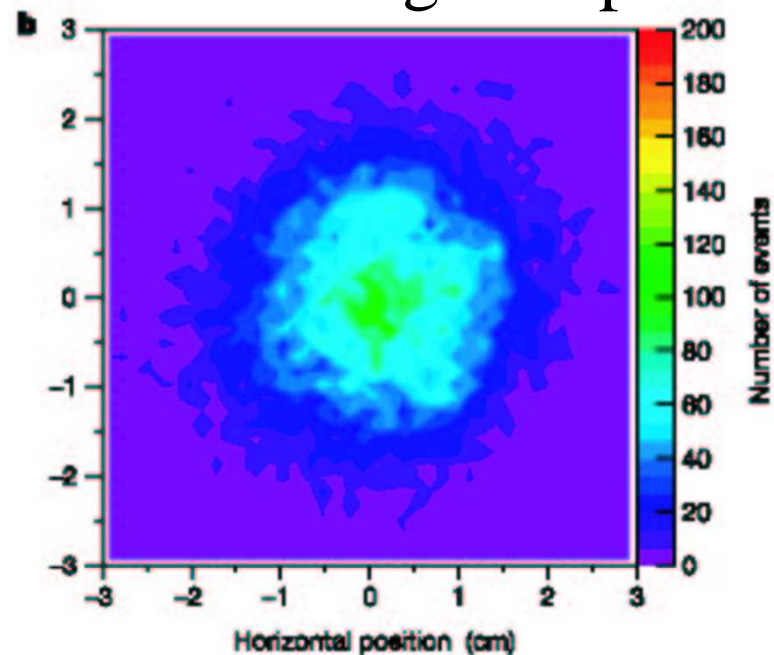


# Signal

Cold Mixing  
Hits on walls



Hot Mixing  
Hits on gas impurities



**Figure 3** Colour contour plots of the distribution (obtained by projecting into the plane perpendicular to the magnetic field) of the vertex positions of reconstructed events.

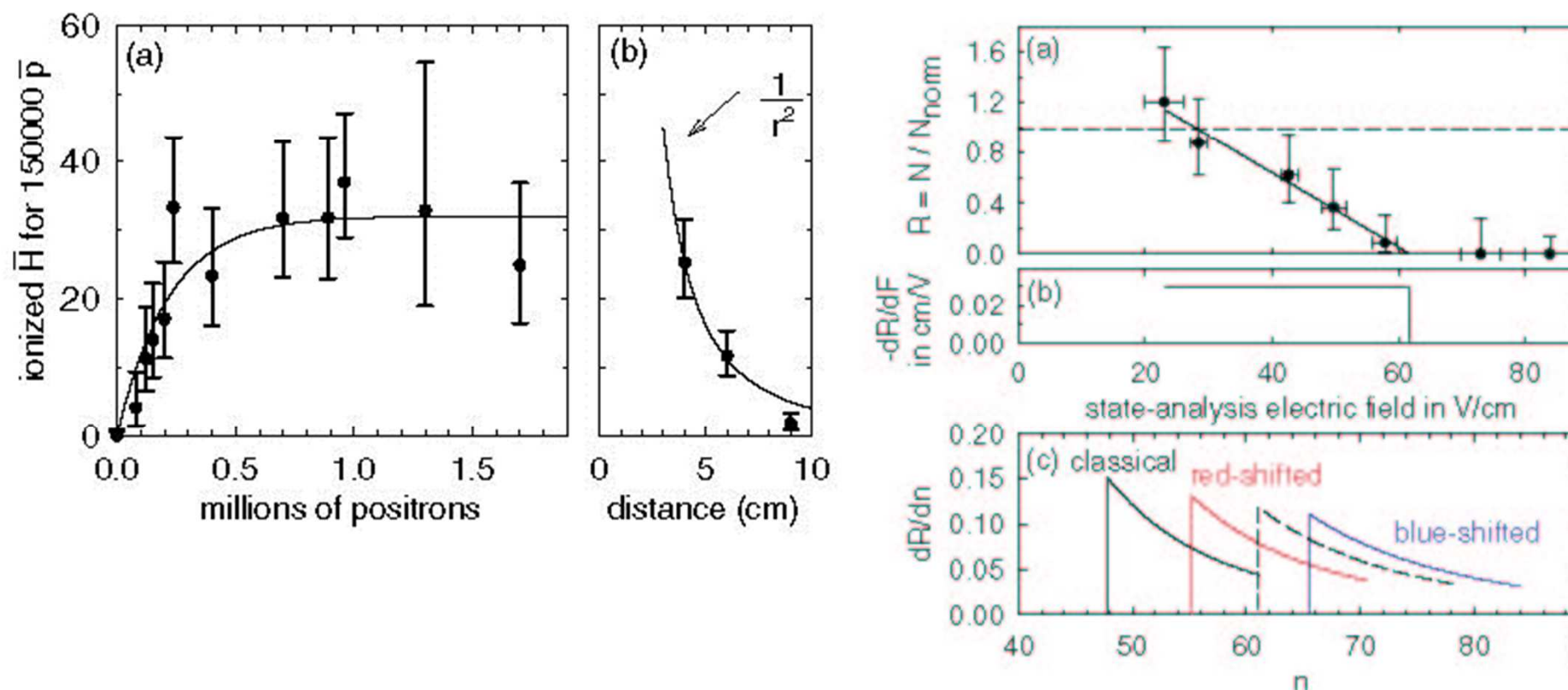
**a**, Cold mixing. All reconstructed antiproton annihilation vertices from the mixing region are plotted—no crystal cuts are applied. The trap inner radius is 1.25 cm. The

annihilations are centred on a slightly smaller radius, in agreement with our Monte Carlo simulations. (Some events appear to be outside of the trap radius owing to vertex reconstruction errors.) **b**, The same plot as above, but for hot mixing. These data are normalized to represent the same number of mixing cycles (165) as those in **a**.

# Background-Free Observation of Cold Antihydrogen with Field-Ionization Analysis of Its States

G. Gabrielse,<sup>1,\*</sup> N. S. Bowden,<sup>1</sup> P. Oxley,<sup>1</sup> A. Speck,<sup>1</sup> C. H. Storry,<sup>1</sup> J. N. Tan,<sup>1</sup> M. Wessels,<sup>1</sup> D. Grzonka,<sup>2</sup> W. Oelert,<sup>2</sup> G. Schepers,<sup>2</sup> T. Seifick,<sup>2</sup> J. Walz,<sup>3</sup> H. Pittner,<sup>4</sup> T. W. Hänsch,<sup>4,5</sup> and E. A. Hessels<sup>6</sup>

(ATRAP Collaboration)



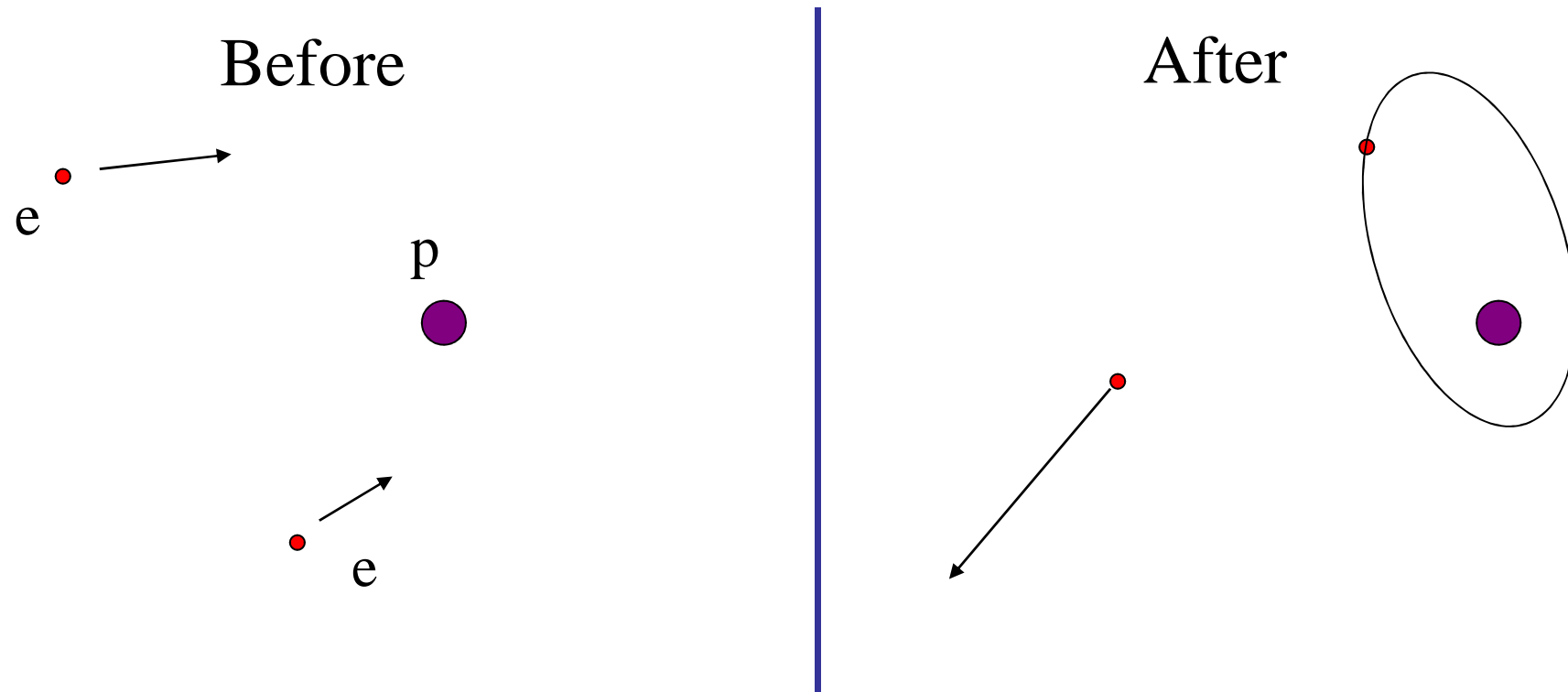
## Three-body recombination for protons moving in a strong magnetic field

F. Robicheaux\* and James D. Hanson

*Department of Physics, Auburn University, Auburn, Alabama 36849-5311, USA*

(Received 23 July 2003; published 27 January 2004)

Using a classical Monte Carlo method, we have computed the three-body recombination (two free electrons and a proton scattering into one free electron and a hydrogen atom:  $e + e + p \rightarrow \text{H} + e$ ) in strong magnetic fields. The proton is allowed its full motion whereas the motion of the electron is given by the guiding center approximation. We investigate recombination for temperatures and fields similar to those used in recent ex-



# Results: Fast H?

$$\text{Rate} = C n^2 v b^5$$

$n$  = density

$$v = (k_B T/m_e)^{1/2}$$

$$b = k e^2/k_B T$$

TABLE II. Three-body recombination coefficient as a function of proton speed along the  $B$  field.  $B=5.4$  T,  $T_e=4$  K. The speed of the proton is given in units of electron thermal speed  $V_0 = \sqrt{2k_B T_e/m_e} = 1.1 \times 10^4$  m/s. The energy of a proton with speed  $V_z$  is also given.

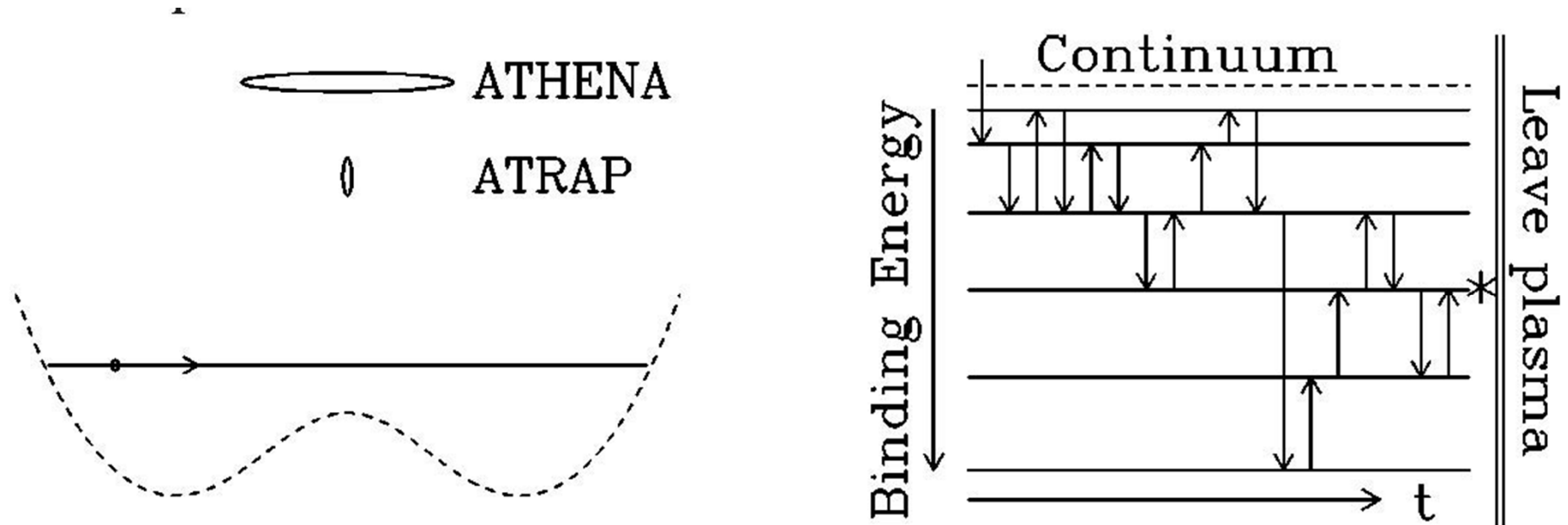
$V_z/V_0$	$C$	$E$ (eV)
0.000	0.100	0.00
0.167	0.081	0.04
0.333	0.051	0.14
0.500	0.031	0.32
0.667	0.018	0.56
0.833	0.011	0.88
1.000	0.008	1.27

## Simulations of antihydrogen formation

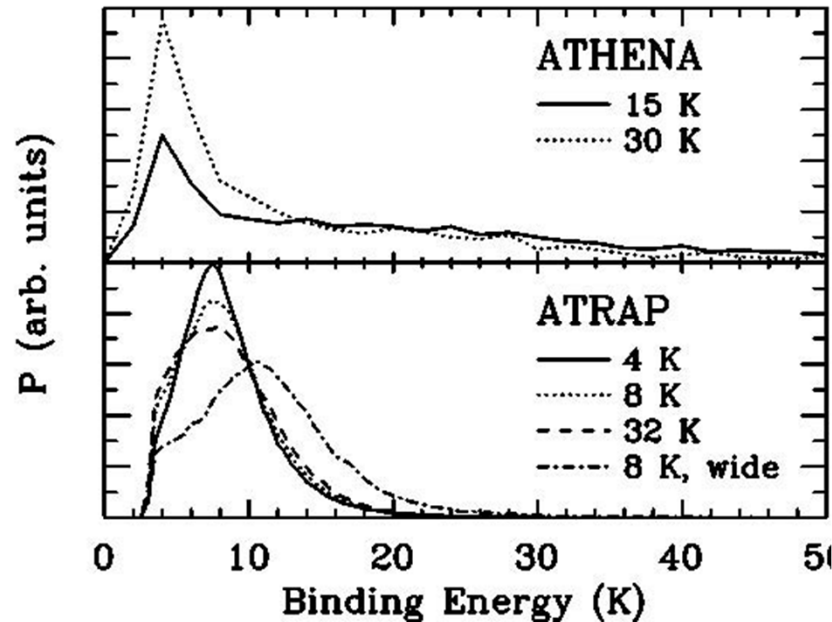
F. Robicheaux\*

*Department of Physics, Auburn University, Alabama 36849-5311, USA*

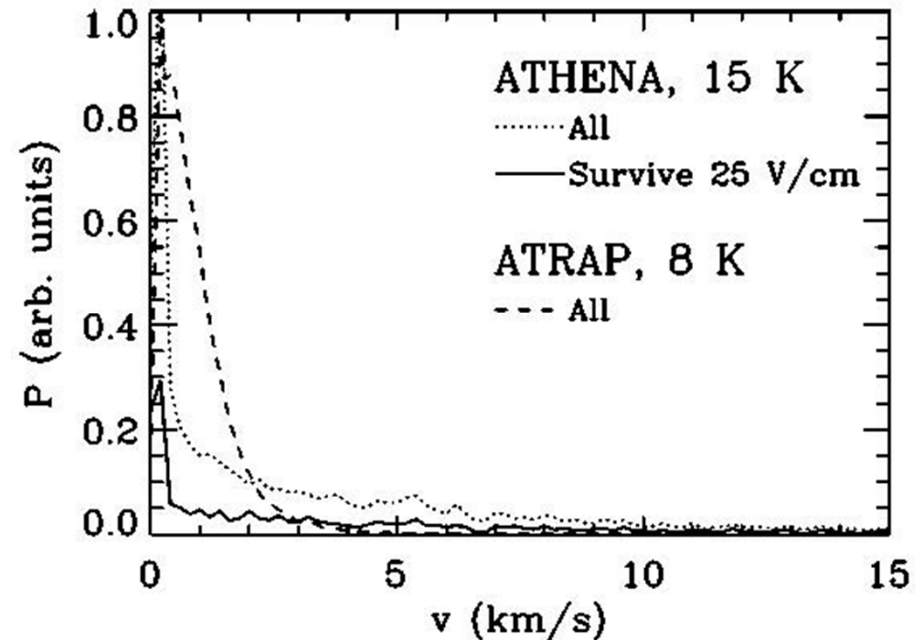
(Received 31 October 2003; published 20 August 2004)



# Results: Small BE & Fast!



Binding energy  
much less than  
expected



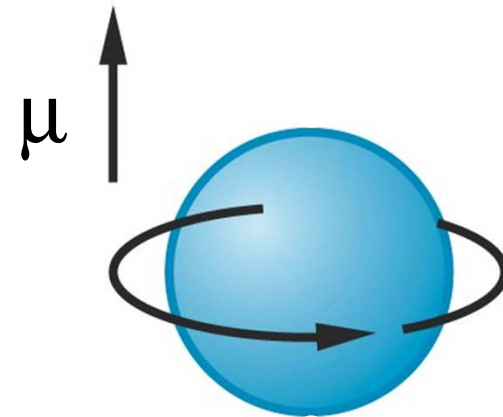
Speed of Hbar much  
larger than thermal  
(250 m/s) or trap  
depth (110 m/s).

# Trapping Neutral Particles

$$PE = - \vec{\mu} \bullet \vec{B}$$

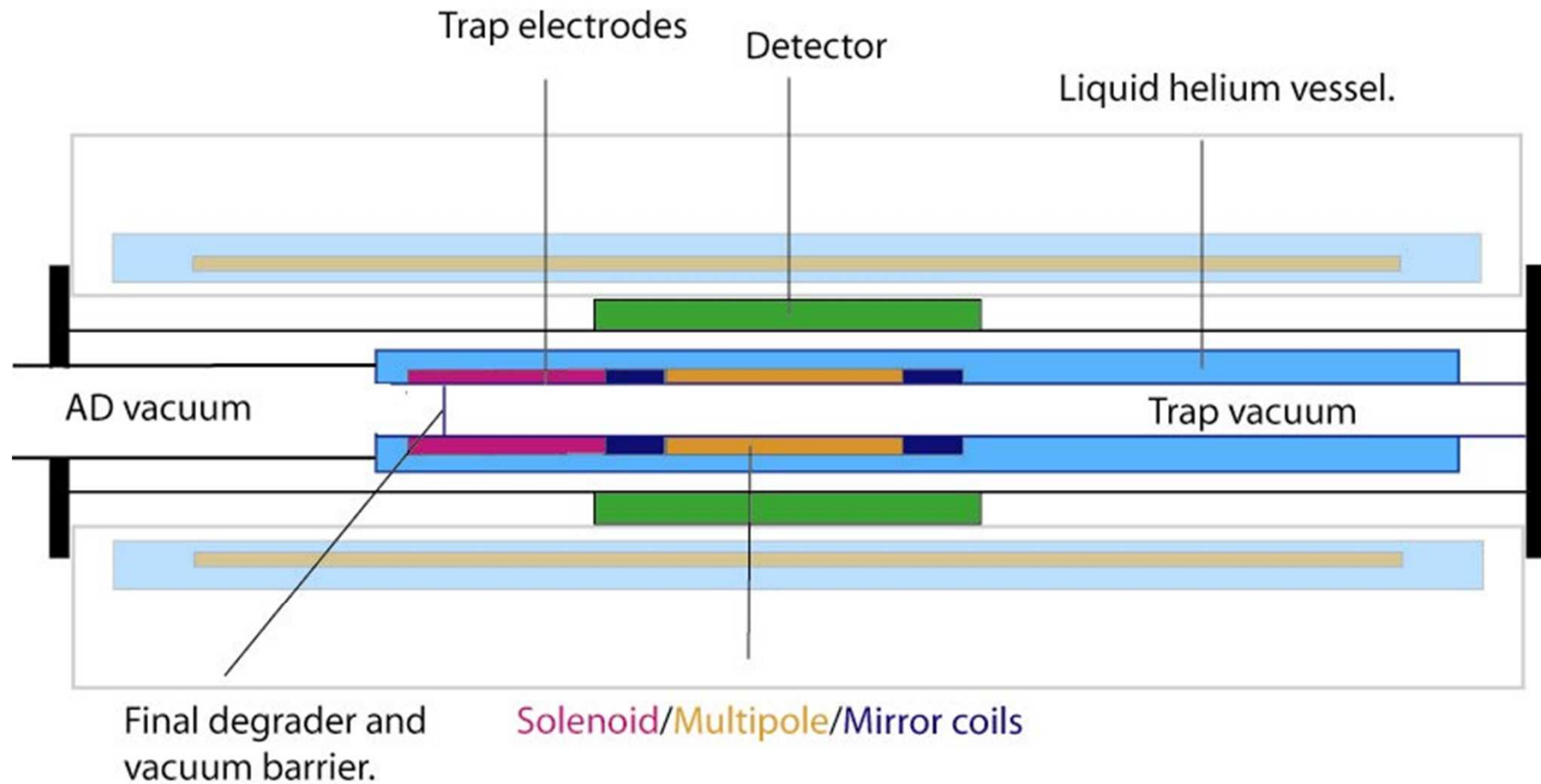
$$\vec{\mu} \bullet \hat{B} \cong \text{constant}$$

$$PE = - \left( \vec{\mu} \bullet \hat{B} \right) B(x, y, z)$$

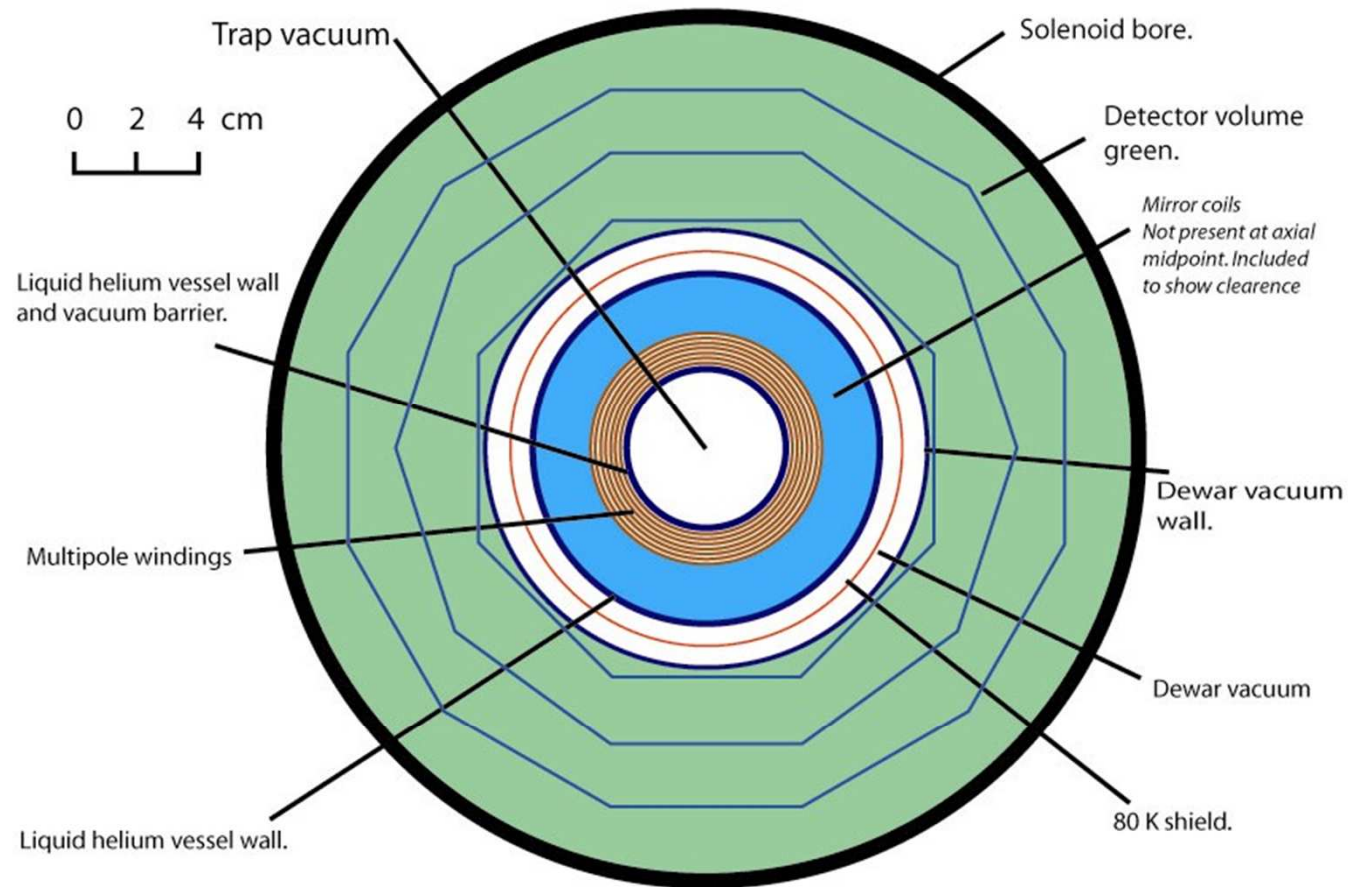


$\mu$  for ground state  $\sim 2/3$  K/T

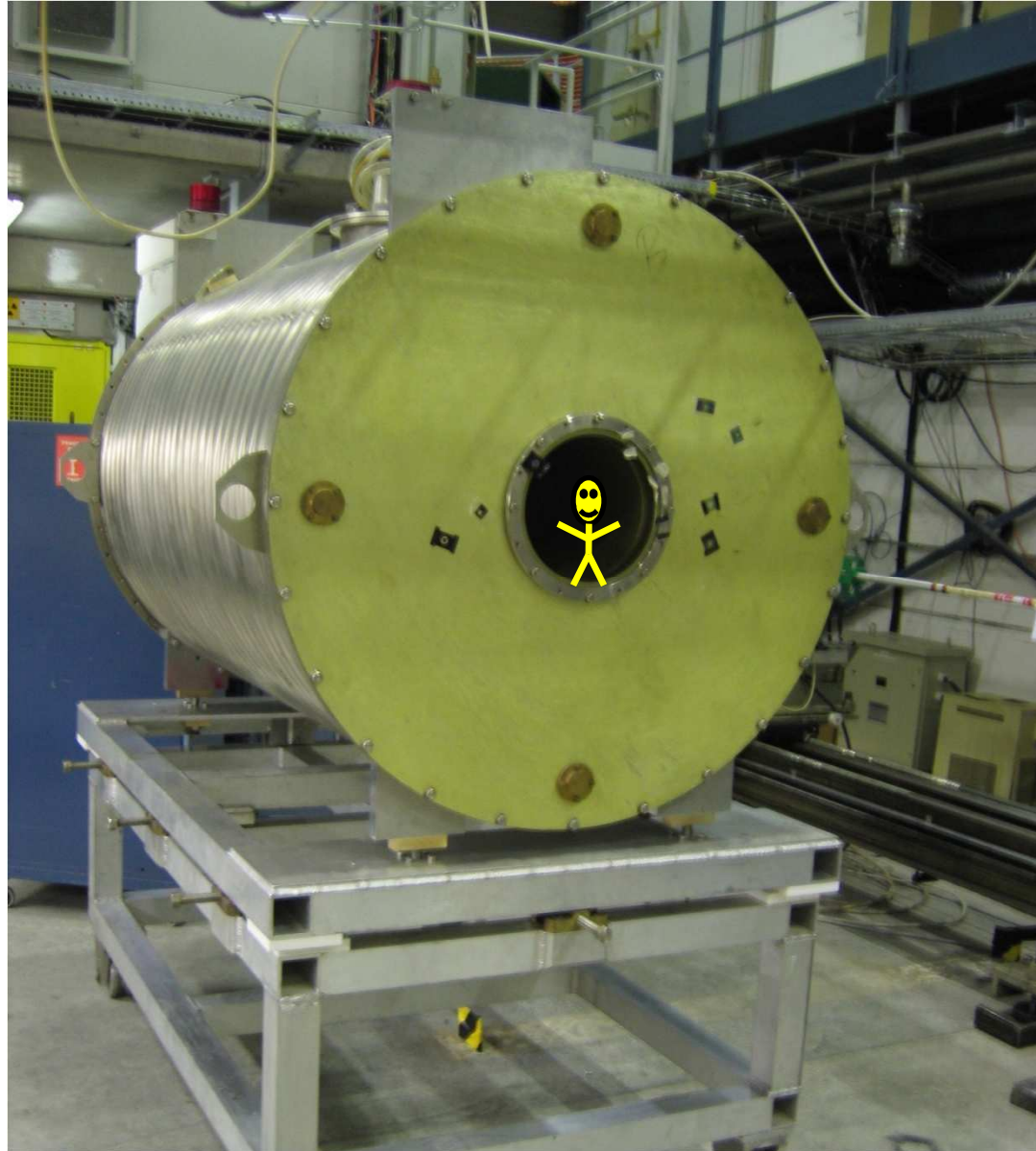
# Antiparticle Trap



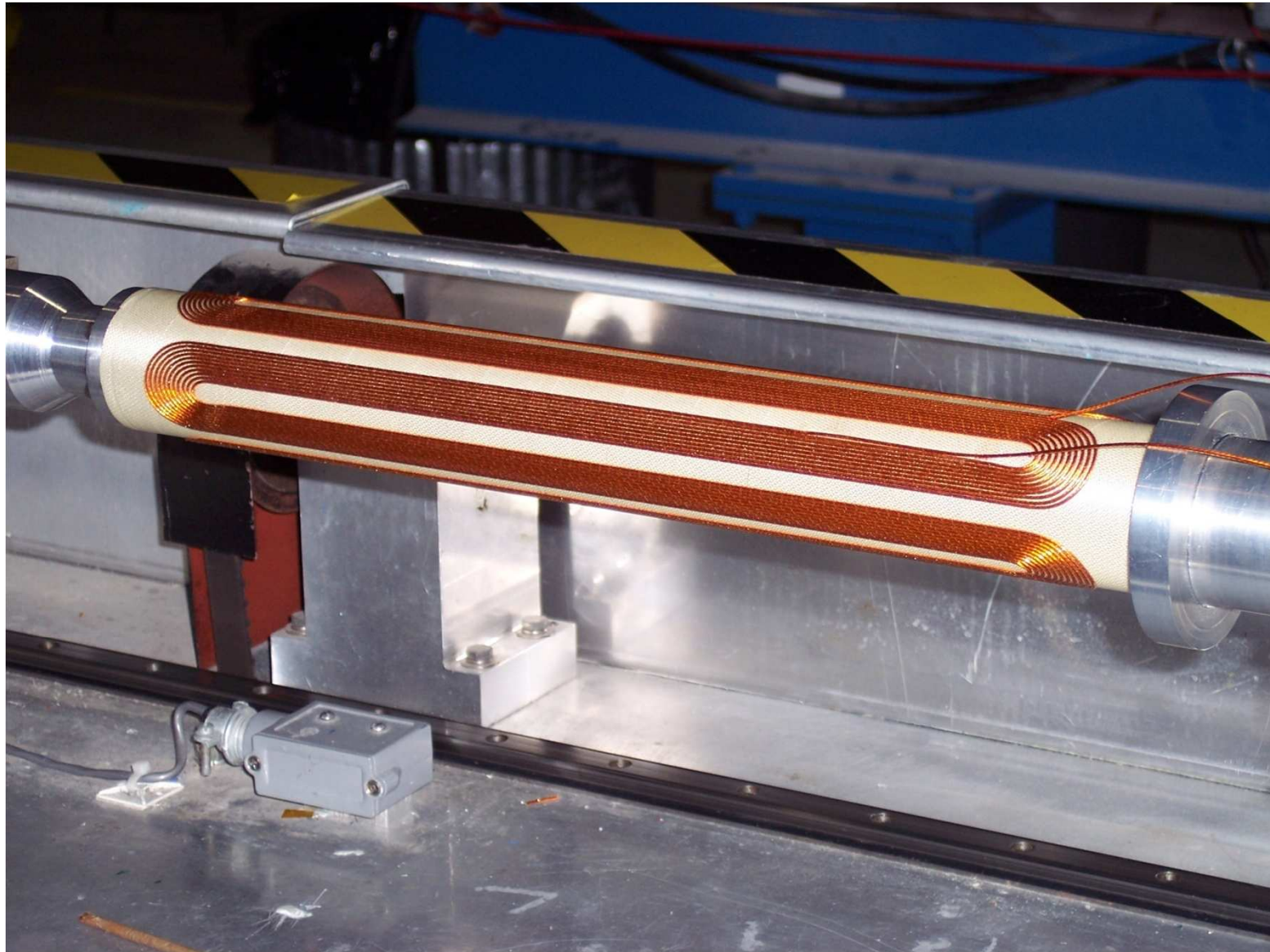
# Antiparticle Trap



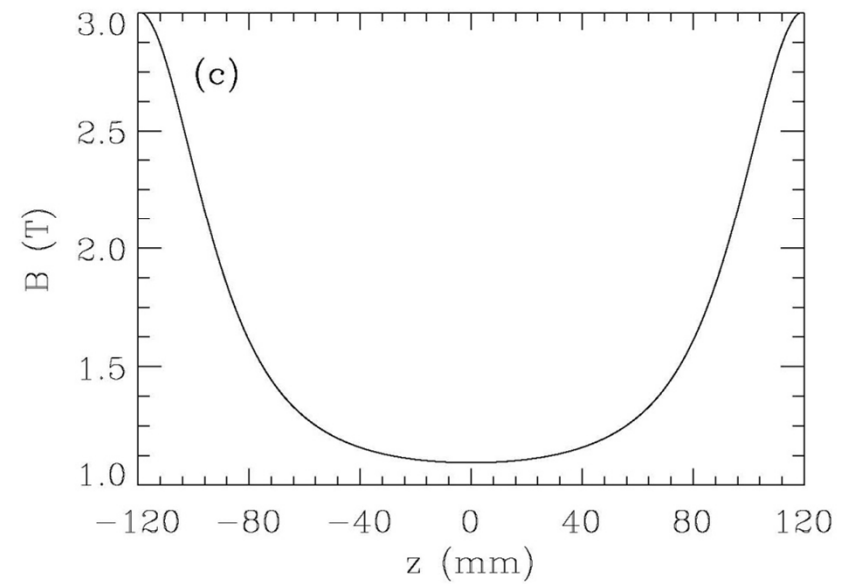
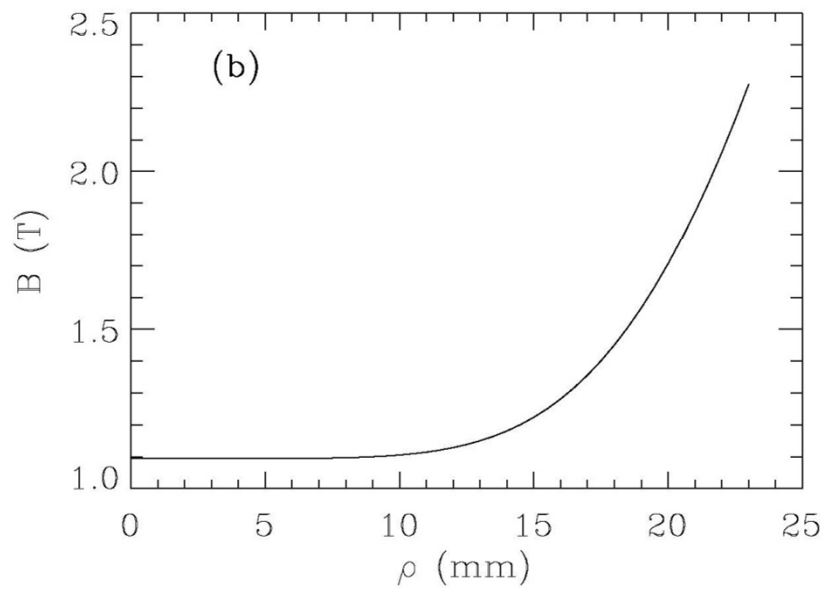
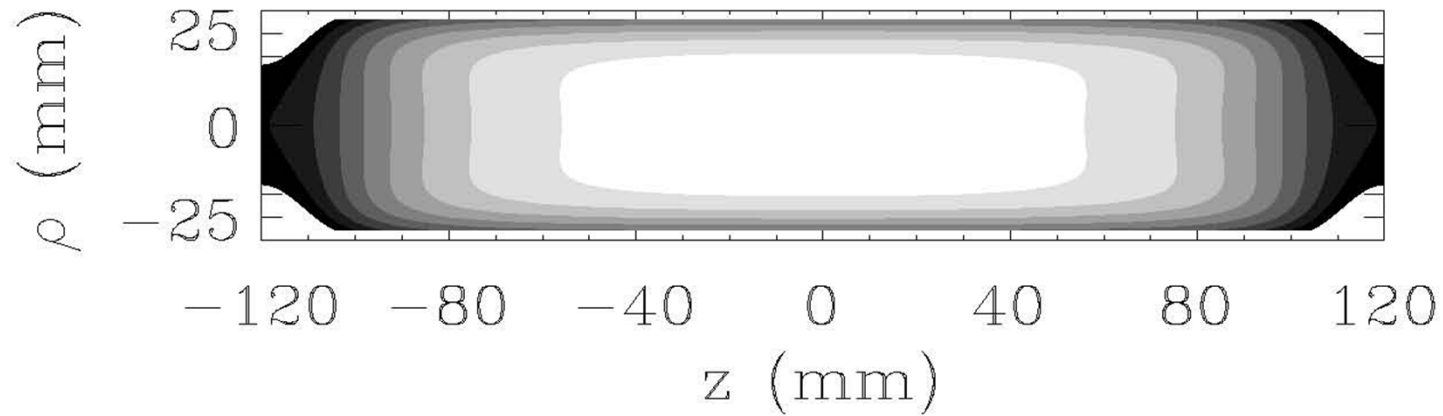
# Actual Magnets: Uniform B



# Actual Magnets: Octupole

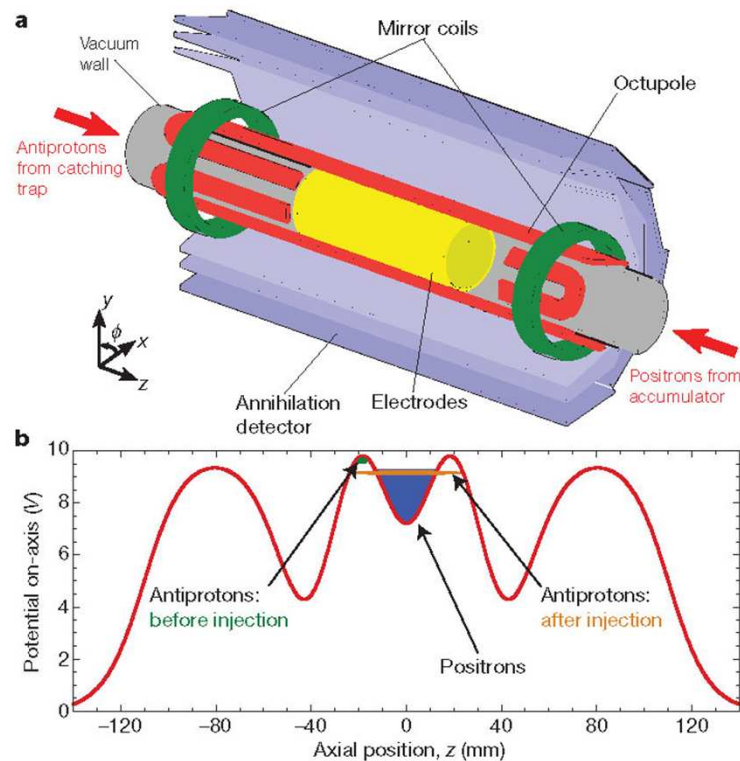


# Antiparticle Trap

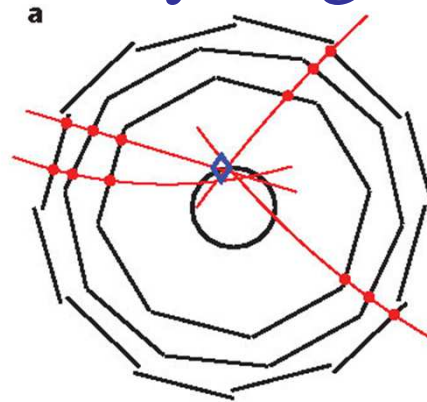


# Trapped antihydrogen

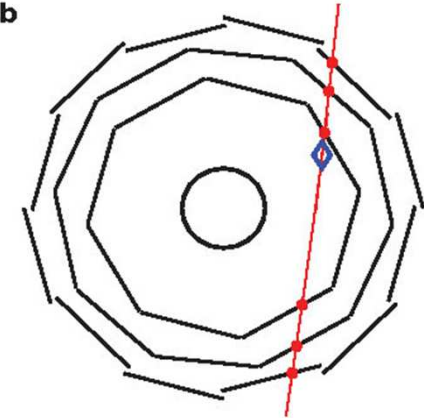
G. B. Andresen<sup>1</sup>, M. D. Ashkezari<sup>2</sup>, M. Baquero-Ruiz<sup>3</sup>, W. Bertsche<sup>4</sup>, P. D. Bowe<sup>1</sup>, E. Butler<sup>4</sup>, C. L. Cesar<sup>5</sup>, S. Chapman<sup>3</sup>, M. Charlton<sup>4</sup>, A. Deller<sup>4</sup>, S. Eriksson<sup>4</sup>, J. Fajans<sup>3,6</sup>, T. Friesen<sup>7</sup>, M. C. Fujiwara<sup>8,7</sup>, D. R. Gill<sup>8</sup>, A. Gutierrez<sup>9</sup>, J. S. Hangst<sup>1</sup>, W. N. Hardy<sup>9</sup>, M. E. Hayden<sup>2</sup>, A. J. Humphries<sup>4</sup>, R. Hydomako<sup>7</sup>, M. J. Jenkins<sup>4</sup>, S. Jonsell<sup>10</sup>, L. V. Jørgensen<sup>4</sup>, L. Kurchaninov<sup>8</sup>, N. Madsen<sup>4</sup>, S. Menary<sup>11</sup>, P. Nolan<sup>12</sup>, K. Olchanski<sup>8</sup>, A. Olin<sup>8</sup>, A. Povilus<sup>3</sup>, P. Pusa<sup>12</sup>, F. Robicheaux<sup>13</sup>, E. Sarid<sup>14</sup>, S. Seif el Nasr<sup>9</sup>, D. M. Silveira<sup>15</sup>, C. So<sup>3</sup>, J. W. Storey<sup>8†</sup>, R. I. Thompson<sup>7</sup>, D. P. van der Werf<sup>4</sup>, J. S. Wurtele<sup>3,6</sup> & Y. Yamazaki<sup>15,16</sup>



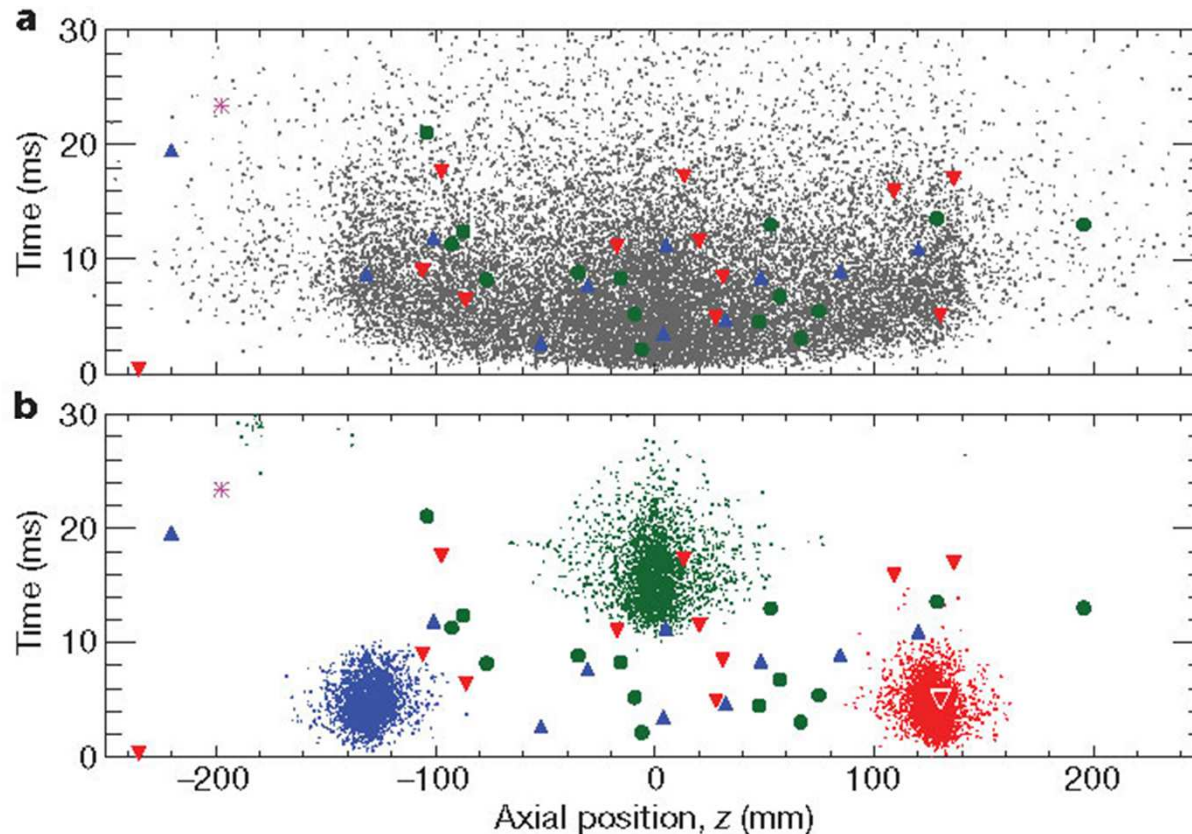
antihydrogen



cosmic



# 1<sup>st</sup> Trapped antiHydrogen

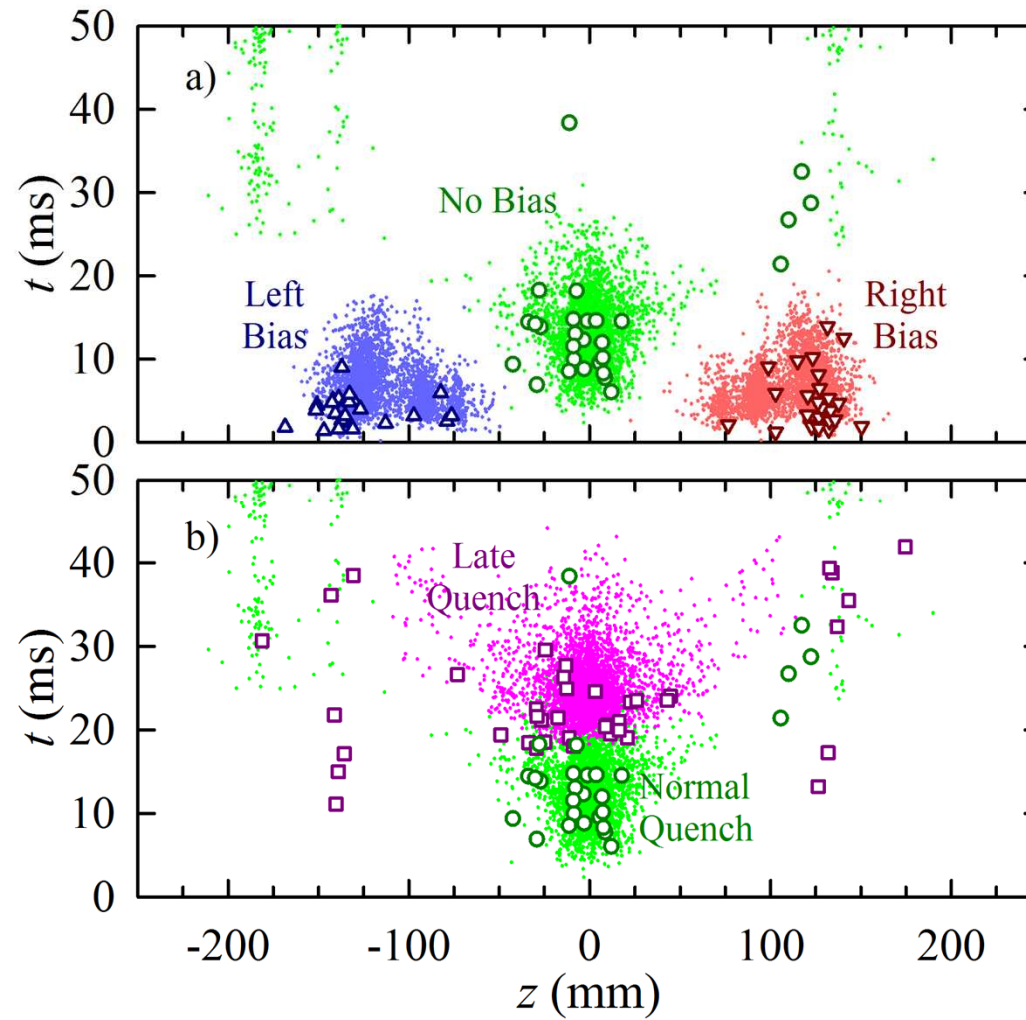


simulation of  
antihydrogen  
hits

simulation of  
antiproton  
hits

filled symbols are experimental hits in  $z$  and time after quench (blue with right E-field, green with no E-field, red with left E-field)

# Trust the pbar simulation?



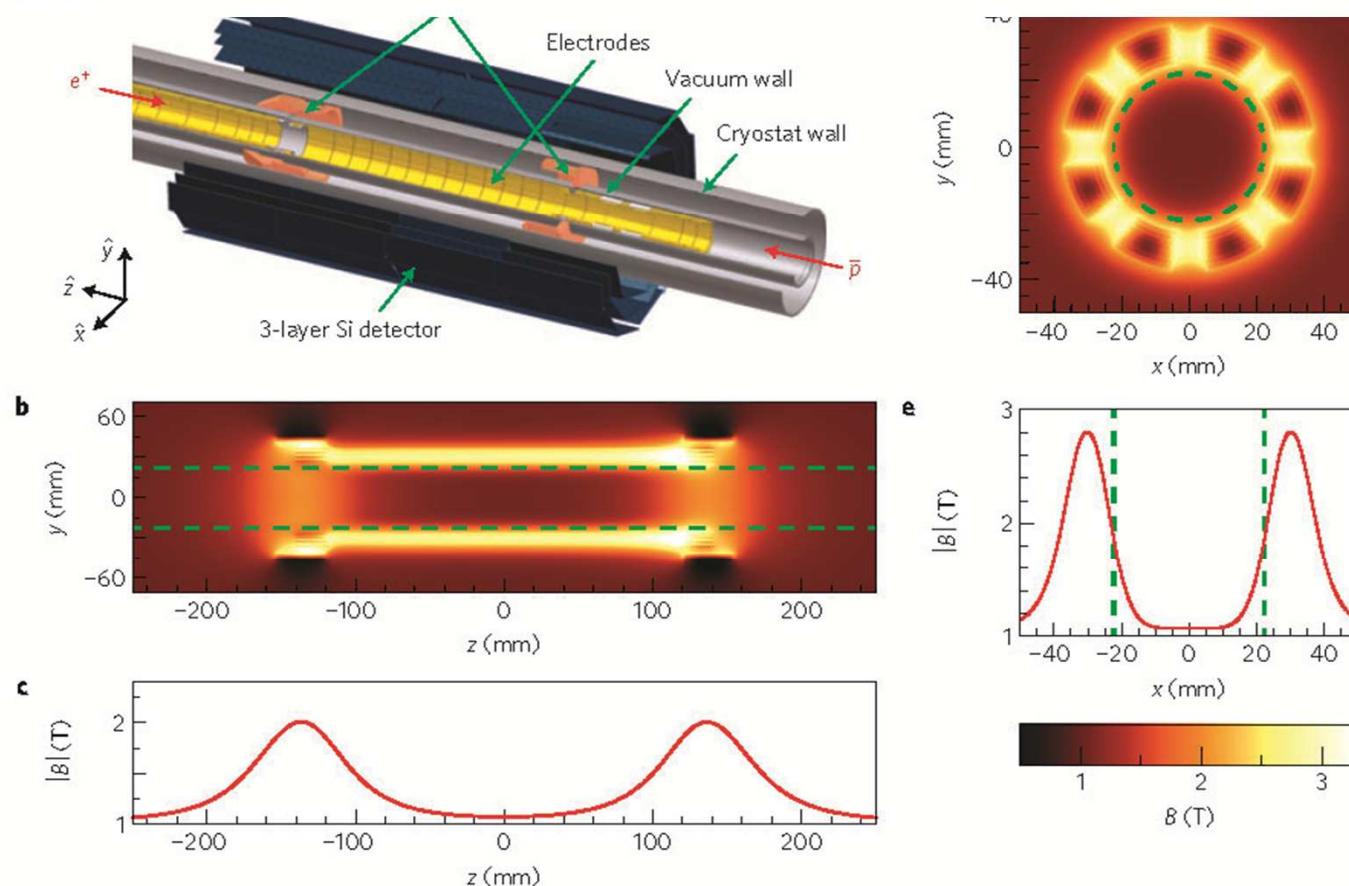
## Holiday Gift Idea No. 1



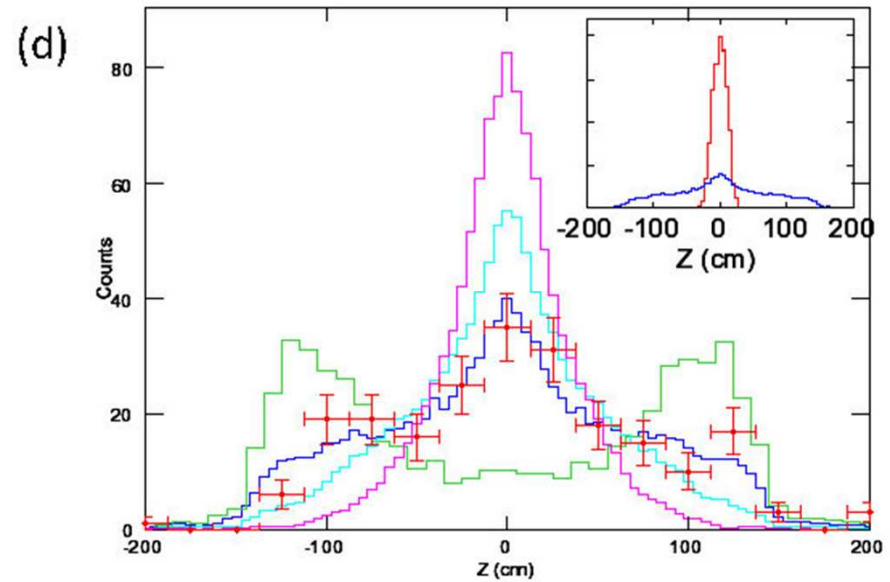
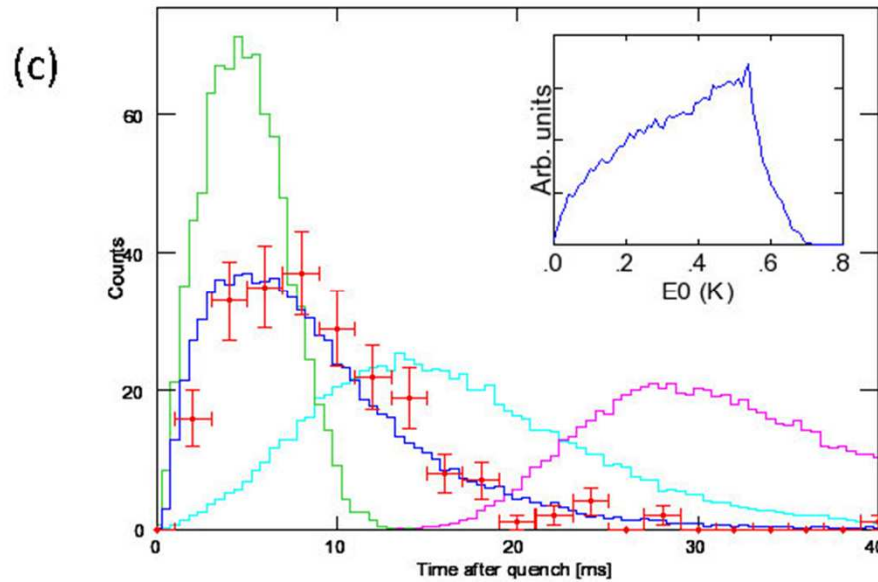
*“Are you kidding?! It’s 38 atoms of anti-hydrogen! They’re priceless!”*

# Confinement of antihydrogen for 1,000 seconds

The ALPHA Collaboration<sup>\*</sup>



# t and z hit distributions

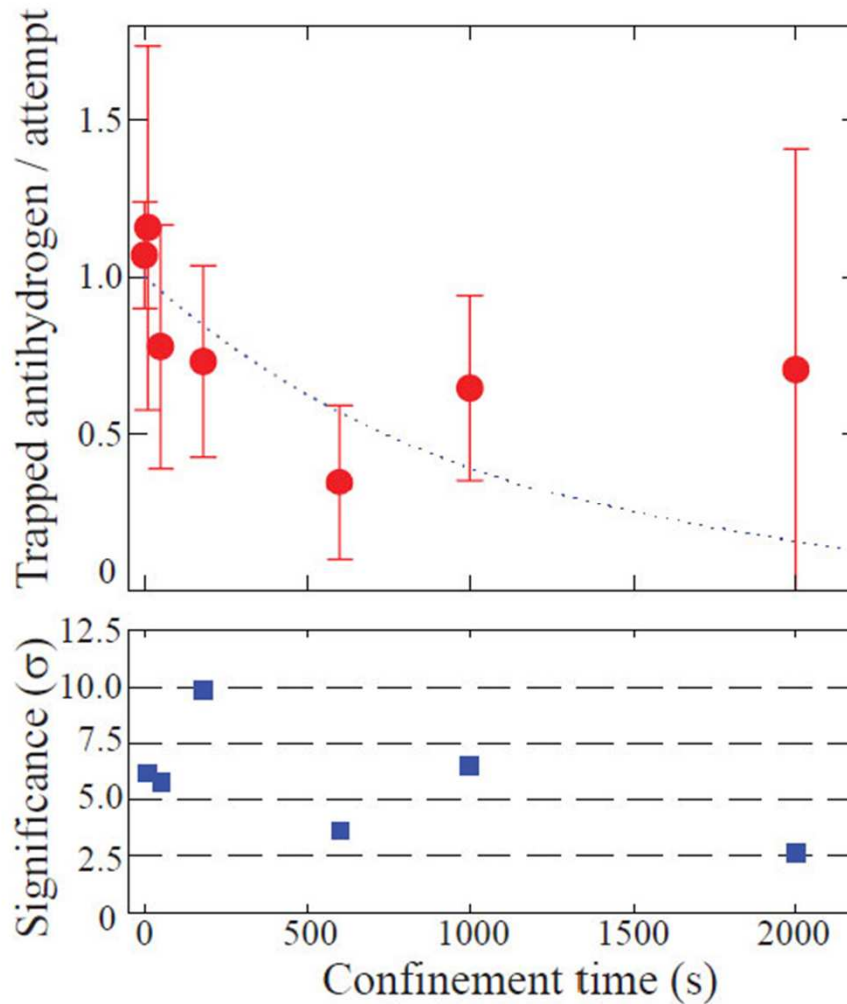


time after quench distrib.

z-hit distrib.

lines are simulations of Hbar hits  
purple 10 mK, light blue 100 mK, blue 50 K  
(any  $T \gg$  well depth), green quasi-bound

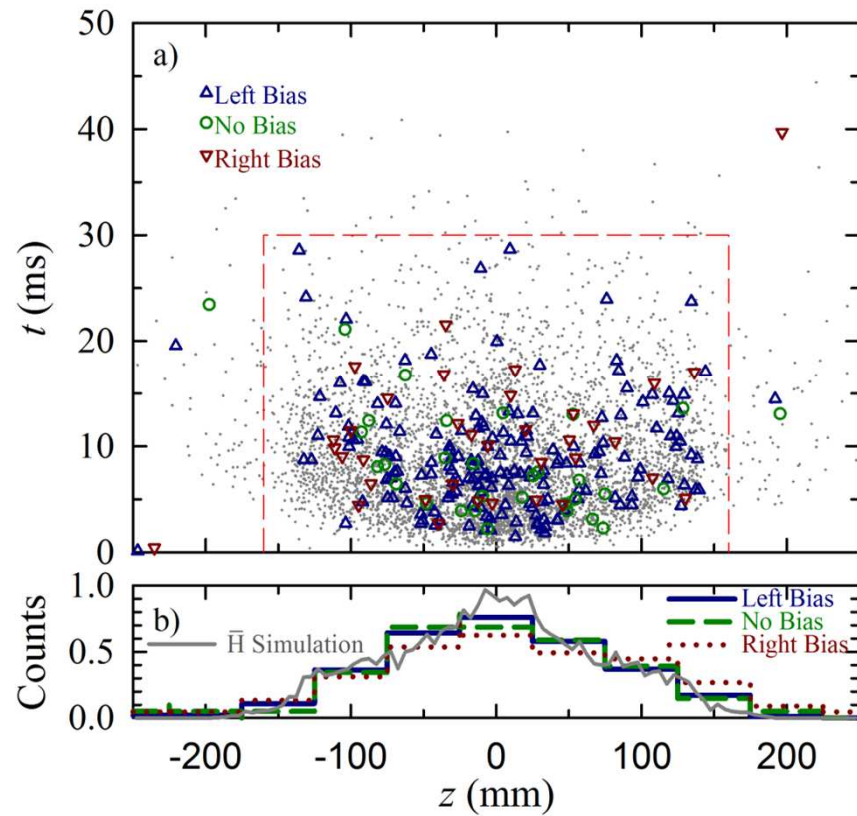
# Number trapped vs delay



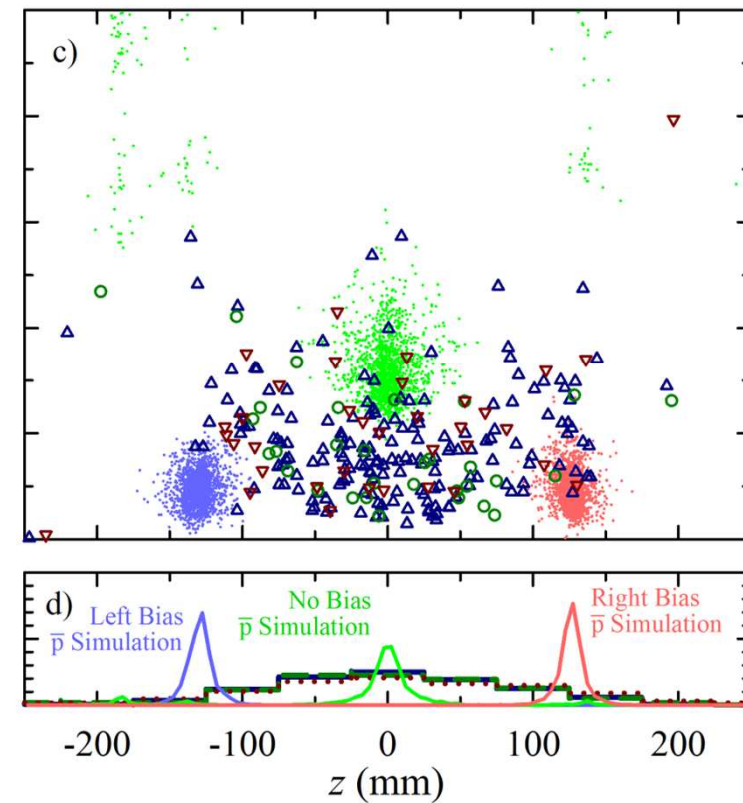
number of hits vs  
delay between  
mixing and  
magnet quench

# Full z-t data/simulation compare

Antihydrogen Simulation Comparison



Antiproton Simulation Comparison



# Do Experiments with antiH!!!

Four obvious experiments to try:

Laser spectroscopy of 1s-2s (ALPHA-II)

Microwave spectroscopy of 1s hyperfine states

Charge neutrality

Gravity

More antihydrogen useful for all experiments!!!

# Cold or Small Positron Plasma?

Would you be better off trying to make your positron plasma smaller or colder? Should it be mega-dense?

Important trends:

Very cold and dense gives antiH formation before the antiprotons come into thermal equilibrium.

Wide range of parameters, the antiproton comes into thermal equilibrium with positrons before antiH formed.

Generally, cold is better...but there is another consideration.

# Terrible News!

Positron plasma in B-field rotates.

Electric field parallel to B is  $\sim 0$  inside positron plasma

$$V \sim e \rho r^2 / (4 \epsilon_0)$$

$$E \times B \text{ drift speed} \sim e \rho r / (2 B \epsilon_0)$$

$$\omega = v/r \sim e \rho / (2 B \epsilon_0)$$

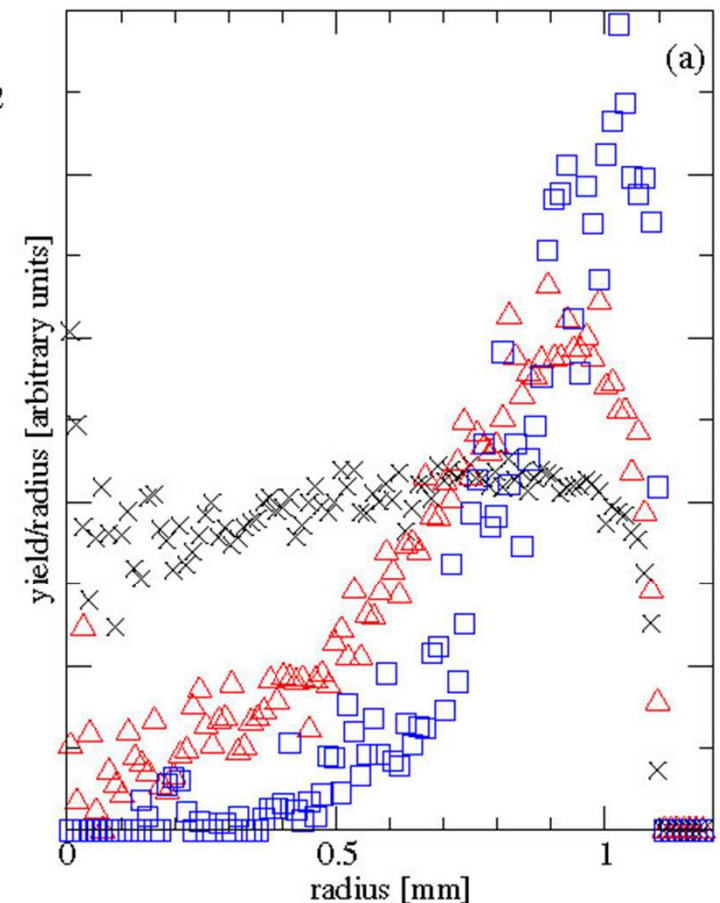
The antiproton rotates with the positron plasma (same EXB drift). **The antiH will have an initial KE from thermal distribution plus the plasma rotation.**

# Simulation of the formation of antihydrogen in a nested Penning trap: effect of positron density

S Jonsell<sup>1</sup>, D P van der Werf<sup>1</sup>, M Charlton<sup>1</sup> and F Robicheaux<sup>2</sup>

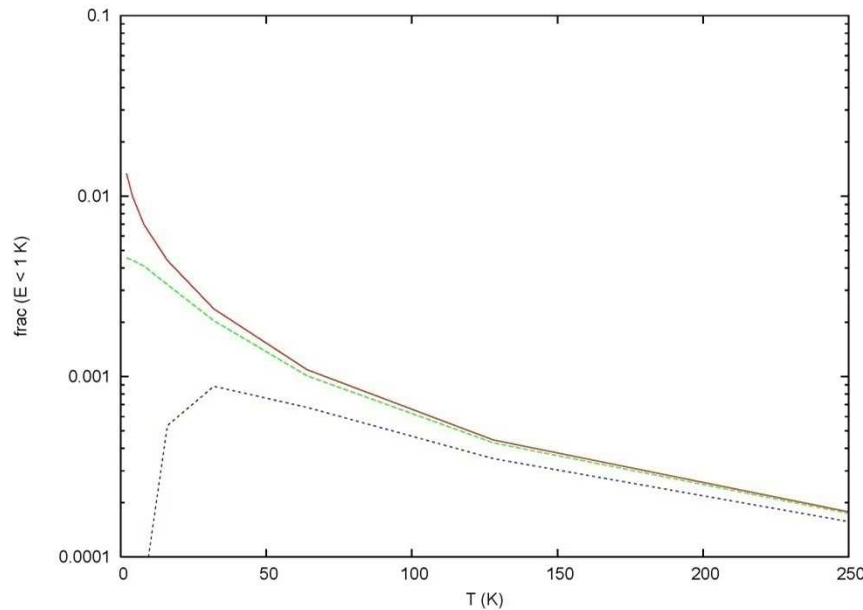
Radial position of where antiH formed during 0-0.1 ms, 0.4-0.5 ms, and > 1 ms

Radial drift to edge of positron plasma due to multiple formation and stripping of antiH

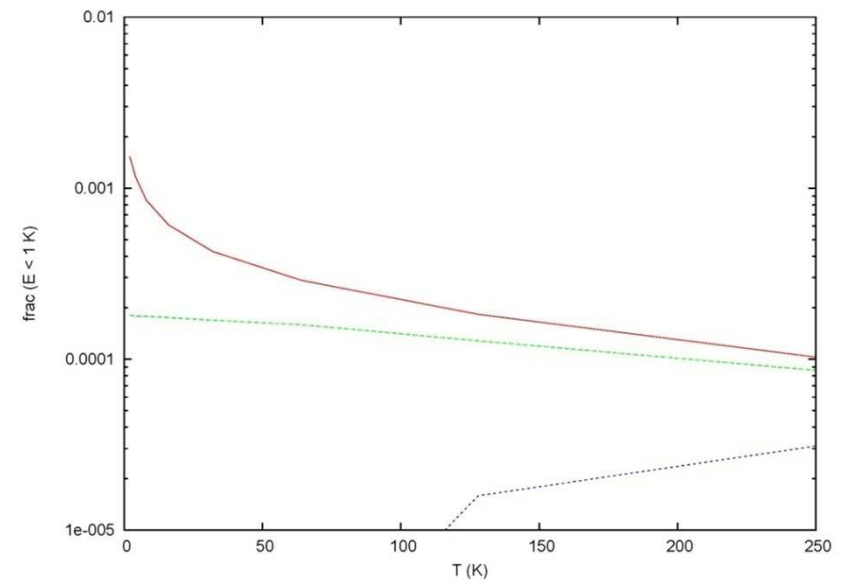


# Fraction of antiH formed $< 1$ K

Positron plasma 1 mm radius



3 mm radius



Red: evenly distributed formation of antiH

Green: formation proportional to  $r$

Blue: formation only at the edge of the plasma

**Can gain big if** make positron plasma radially small

# Plasmas Are Nasty

Expanding plasma has PE go to KE: gives substantial heating.

Charged plasma with cylindrical symmetry is stable.  
Expands slowly due to noise and slight imperfections.

Mirrors + uniform B = OK

Mirrors + octupole + uniform B = not OK

Fast expansion and heating of positrons unless radius very small. Even small plasmas are 2-4X hotter than w/out octupole.

Typical temperatures are  $\sim 40+$  K.

# An “Obvious” Problem

Anti-protons are in a plasma with electrons

Electrons cool the anti-protons & reduce  $r$  (rotating wall)

Recombination with electrons around?

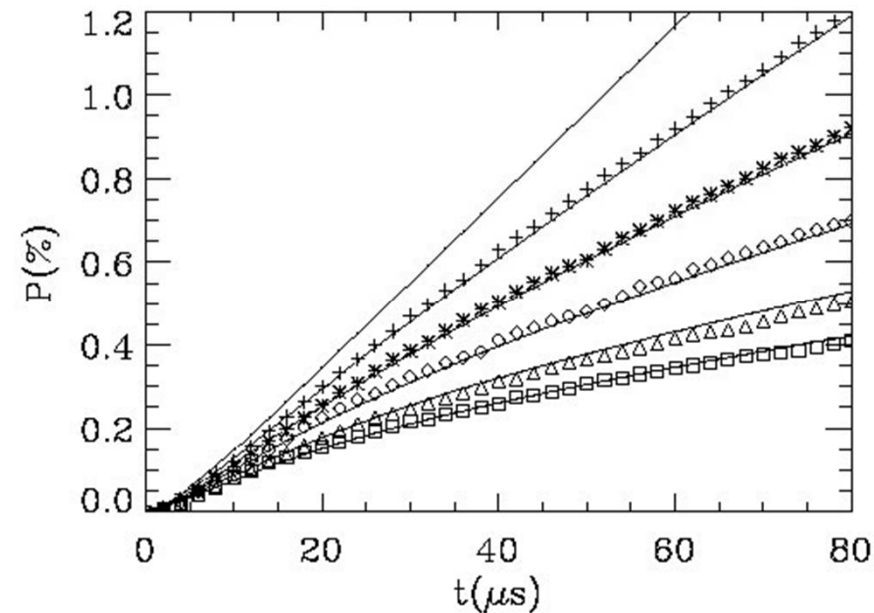
Positrons need to strongly outnumber the electrons near the anti-proton or TBR is strongly suppressed.

FR, J. Phys. B **40**, 271 (2007): electron/positron  $< 0.2$

Also, if you try a dynamic mixing, average charge/mass will fluctuate.

## Three-body recombination with mixed sign light particles

F Robicheaux



Recombination in 4 K plasma to states bound by  $> 16$  K,  
for fraction of electrons: 0.0, 0.04, 0.08, 0.12, 0.16, 0.20

Lines from simple model calculation

# An “Obvious” Solution

“Kick” the electrons out of the plasma.

Small electron mass means anti-protons hardly move during kick.

Heating of antiprotons on the 100-1000 K scale.

# An “Obvious” Problem

Anti-protons are in a plasma, with no electrons, at a temperature more than few 100X too big for trapping.



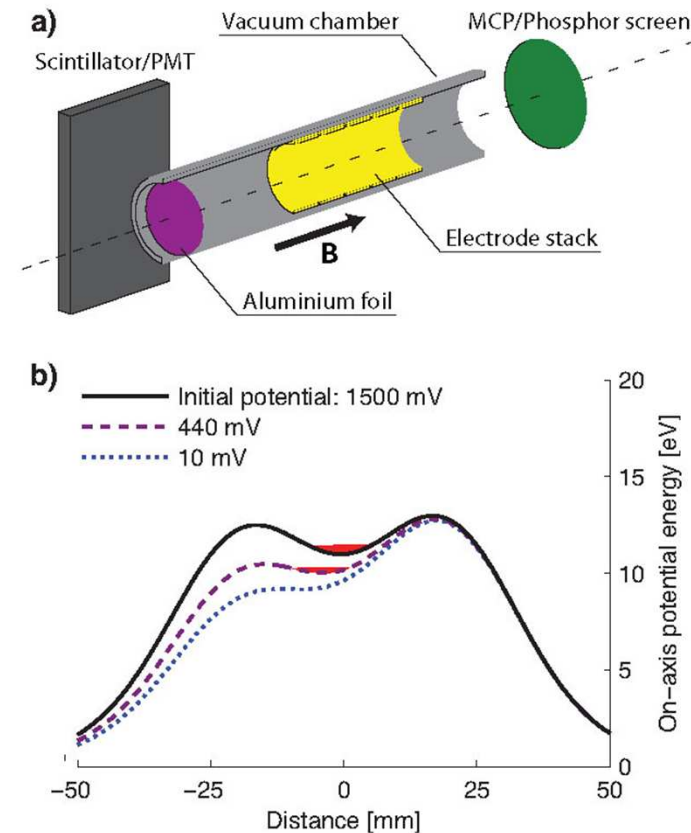
## Evaporative Cooling of Antiprotons to Cryogenic Temperatures

G. B. Andresen,<sup>1</sup> M. D. Ashkezari,<sup>2</sup> M. Baquero-Ruiz,<sup>3</sup> W. Bertsche,<sup>4</sup> P. D. Bowe,<sup>1</sup> E. Butler,<sup>4</sup> C. L. Cesar,<sup>5</sup> S. Chapman,<sup>3</sup> M. Charlton,<sup>4</sup> J. Fajans,<sup>3</sup> T. Friesen,<sup>6</sup> M. C. Fujiwara,<sup>7</sup> D. R. Gill,<sup>7</sup> J. S. Hangst,<sup>1</sup> W. N. Hardy,<sup>8</sup> R. S. Hayano,<sup>9</sup> M. E. Hayden,<sup>2</sup> A. Humphries,<sup>4</sup> R. Hydomako,<sup>6</sup> S. Jonsell,<sup>4,10</sup> L. Kurchaninov,<sup>7</sup> R. Lambo,<sup>5</sup> N. Madsen,<sup>4</sup> S. Menary,<sup>11</sup> P. Nolan,<sup>12</sup> K. Olchanski,<sup>7</sup> A. Olin,<sup>7</sup> A. Povilus,<sup>3</sup> P. Pusa,<sup>12</sup> F. Robicheaux,<sup>13</sup> E. Sarid,<sup>14</sup> D. M. Silveira,<sup>15,16</sup> C. So,<sup>3</sup> J. W. Storey,<sup>7</sup> R. I. Thompson,<sup>6</sup> D. P. van der Werf,<sup>4</sup> D. Wilding,<sup>4</sup> J. S. Wurtele,<sup>3</sup> and Y. Yamazaki<sup>15,16</sup>

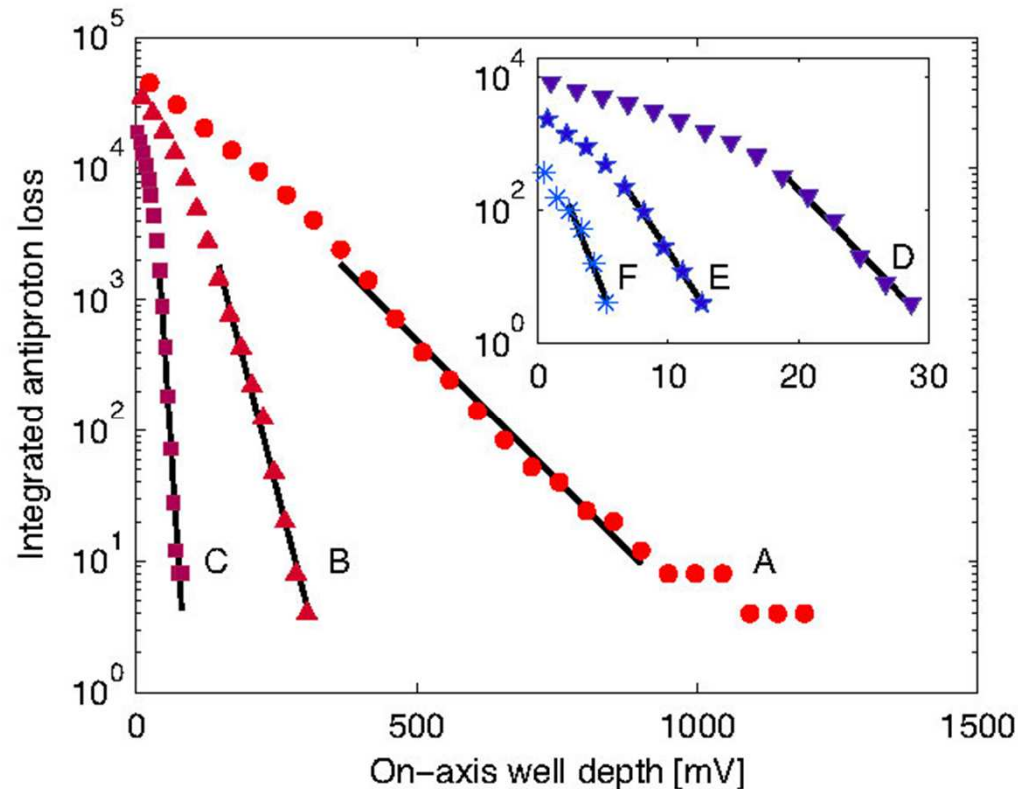
(ALPHA Collaboration)

A somewhat different case from usual evaporative cooling because of strong  $B$  & coulomb scattering.

Slowly dump antiprotons out of their trapping well.



# Dump antiprotons to measure T



Maxwell-Boltzmann tail; steeper means lower T

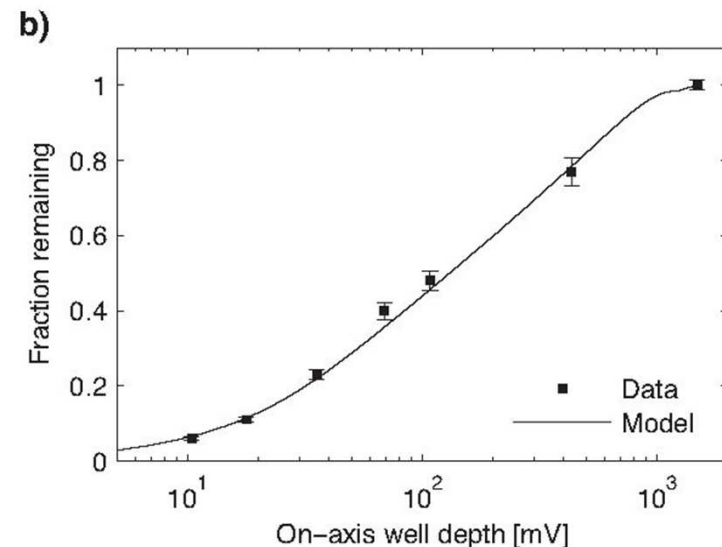
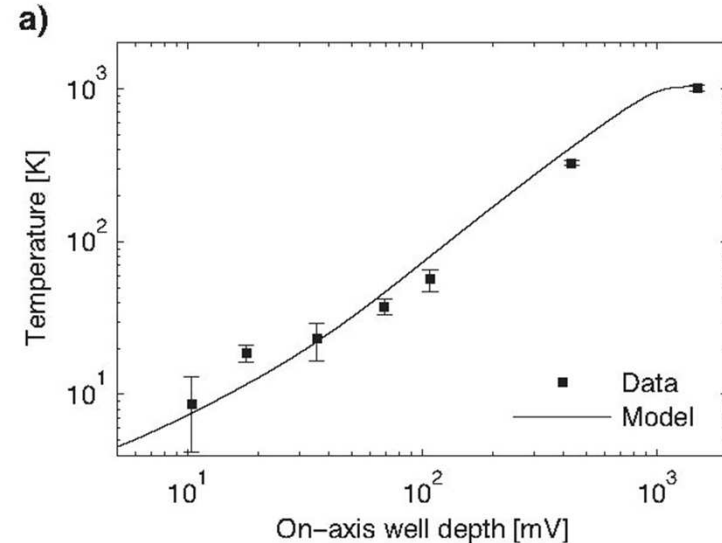
A=1040 K, B=325 K, C=57 K, D=23 K, E=19 K, F=9 K

# More low energy antiprotons

Evaporative cooling means  
you lose antiprotons

Factor of  $\sim 100$  decrease in  
temperature corresponds to  
factor of  $\sim 20$  decrease in  
number

But!! Average  $r$  increases  
because evaporate antiprotons  
from  $r \sim 0$ .



# Counterintuitive Results: Atoms

Large B-field combined with highly excited atoms can give strange result.

Microwave photons can cool center of mass motion...

Lower temperature not always better...

# Cooling AntiH: Microwaves?

# Cooling of Rydberg $\bar{\text{H}}$ during radiative cascade

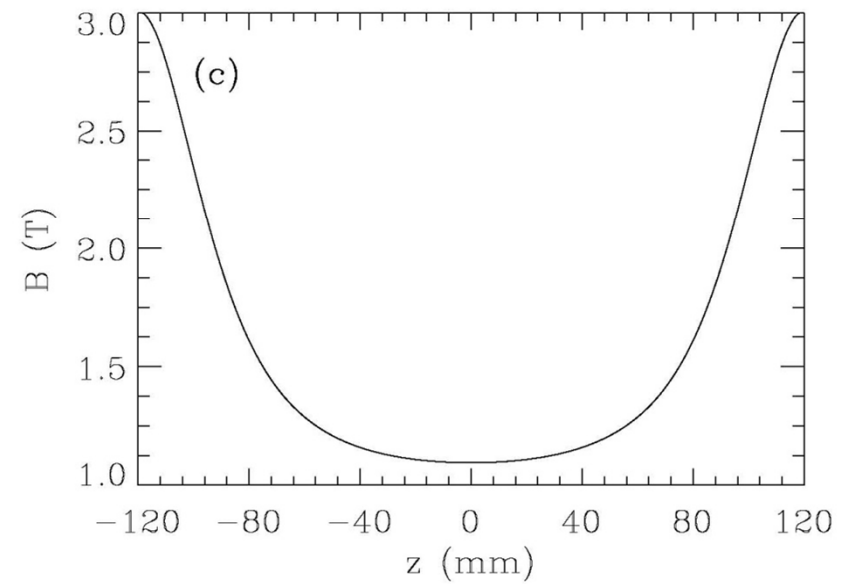
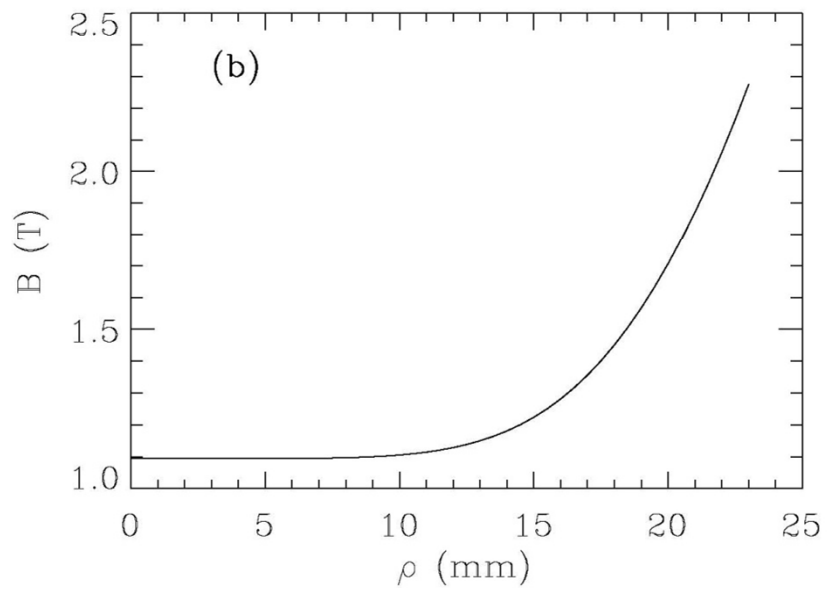
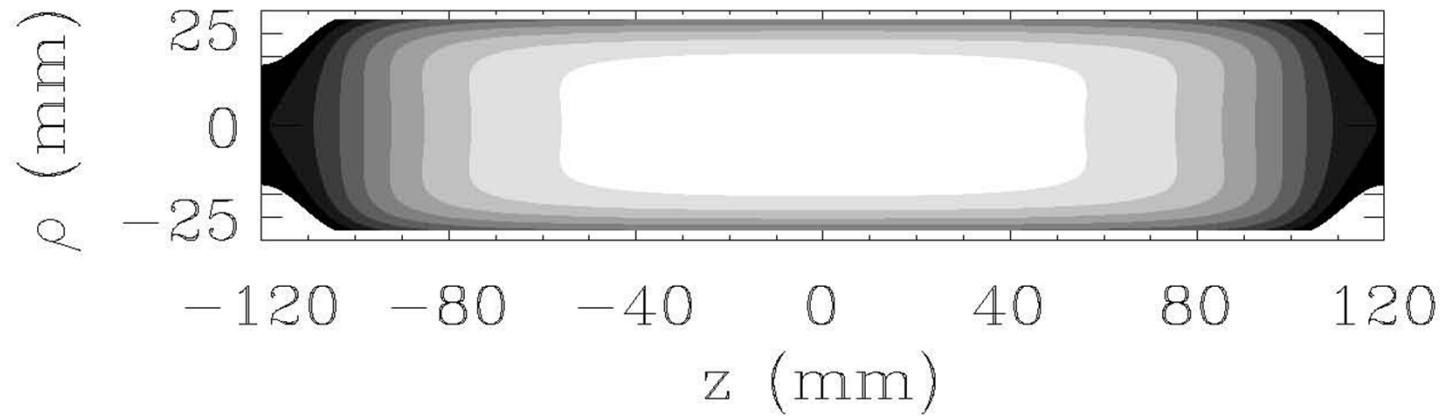
C L Taylor, Jingjing Zhang and F Robicheaux

Simulations performed by Michele Zhang & Christine Taylor (**limited geometry tests FR**)

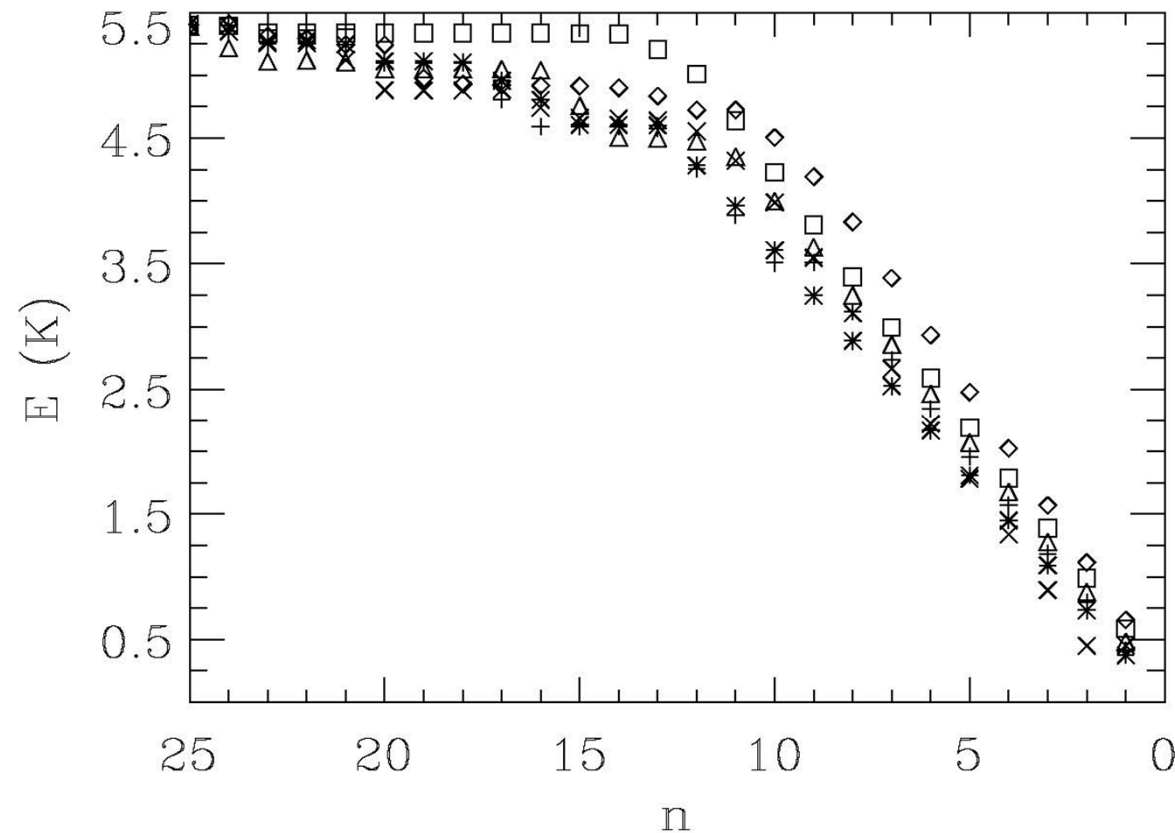
Atoms start out in perfect circular state:  $n, n-1, n-1$   
Approximate non-circular states (quasi-classical rates)

Solve Newton's equation for C.O.M. motion with random emission of photon

# Antiparticle Trap



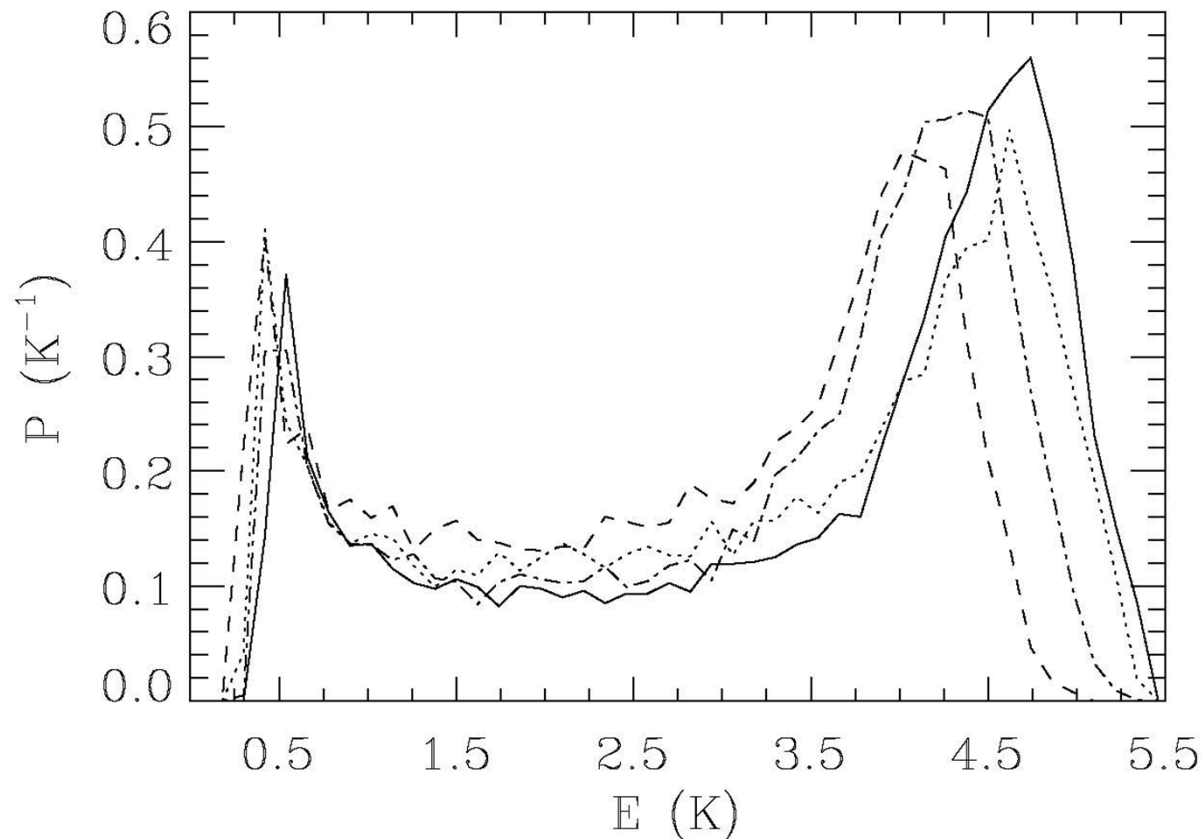
# Adiabatic vs. Sudden Cooling



Initial KE = 5.4 K ( $\sim 300$  m/s),  $n = 25$ , perp to axis, center

Only look at atoms with final  $E_{\text{com}} < 0.79$  K

# Final $E_{\text{com}}$ Distribution



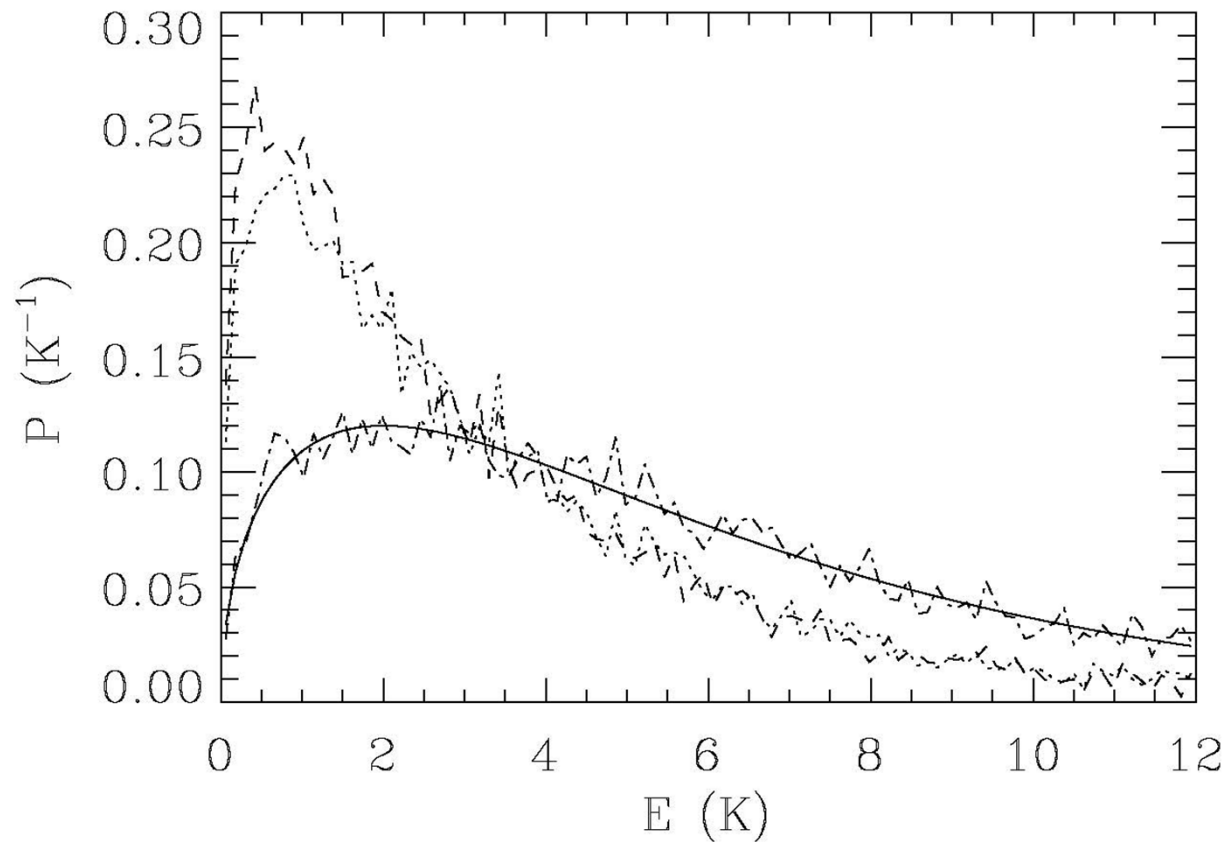
Shape and #  
predicted from  
simple theory

Initial KE = 5.4 K ( $\sim 300$  m/s),  $n = 25$ , center

solid = perp, random  $m > 15$  ; dot-dash = perp, circular

dotted = parallel, random  $m > 15$ ; dashed = parallel, circular

# Final $E_{\text{com}}$ Distribution



Fraction  
Trapped:  
18%  
16%  
6%  
Nice  
increase  
Cushion

Thermal KE = 4 K,  $n = 25$ , spheroid (2mm, 15mm radii)

dash = circ; dotted = random  $m > 15$

dash-dot = initial distribution, solid = Maxwell

## Possible mechanism for enhancing the trapping and cooling of antihydrogen

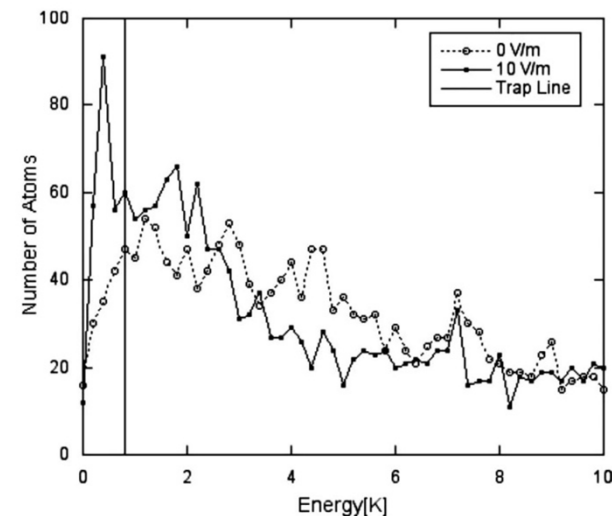
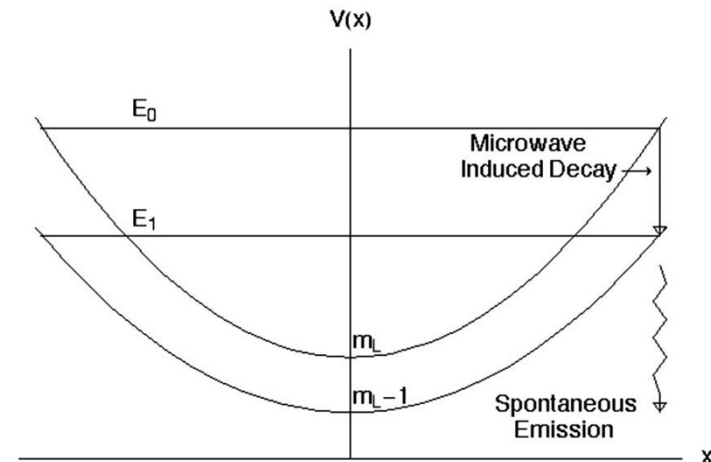
C. L. Cesar,<sup>1</sup> F. Robicheaux,<sup>2</sup> and N. Zagury<sup>1</sup>

Microwave cause  $1,m$  transition at large  $B$

Weak microwave because electric dip.

Lower  $1,m$  radiates faster

Leads to extra cooling (factor of 1.5 started at 16 K)



# TBR Temperature Surprise?

# Three-body recombination for electrons in a strong magnetic field: Magnetic moment

F. Robicheaux

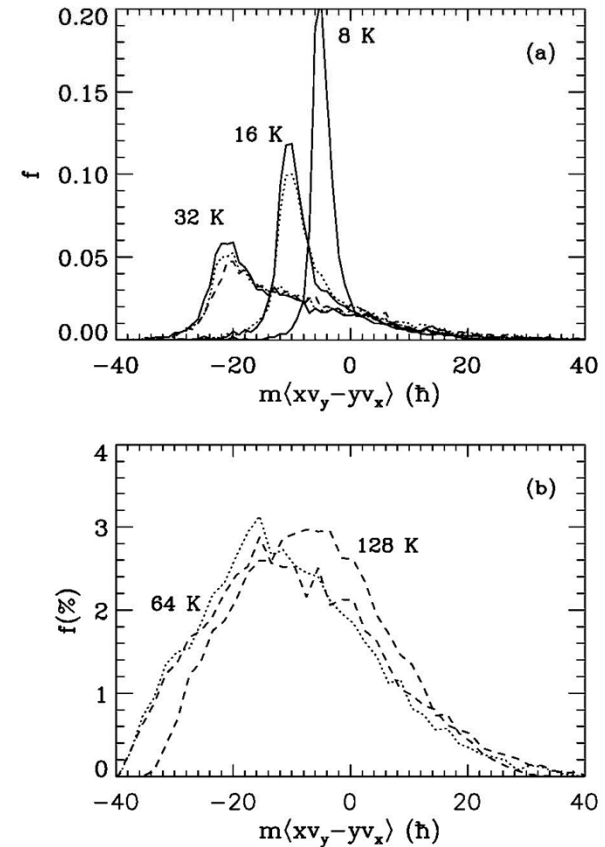
Atoms bound by more than 32 K

Fraction with  $dE/dB > 0$  and  $dE/dB > 10$  K/T **increases** with increasing positron temperature!

4 K positrons: 0.17, 0.035

8 K positrons: 0.21, 0.058

16K positrons: 0.25, 0.063



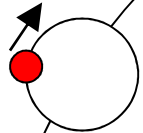
Magnetic moment vs T and BE

$\mu$  proportional to  $m\langle xv_y - yv_x \rangle$

# Guiding Center Atom: $\mu$

low field  
seeker

$e^+$



$$\mu = -m_e v^2/2 B$$

$$= -KE_{\text{cyc}}/B$$

$$\mu = k e^2/2 r B$$

$$= |PE|/2 B$$

high field  
seeker

E-field

$\bar{p}$

$$v_{\text{drift}} = E/B \ll v_{\text{cyc}}$$

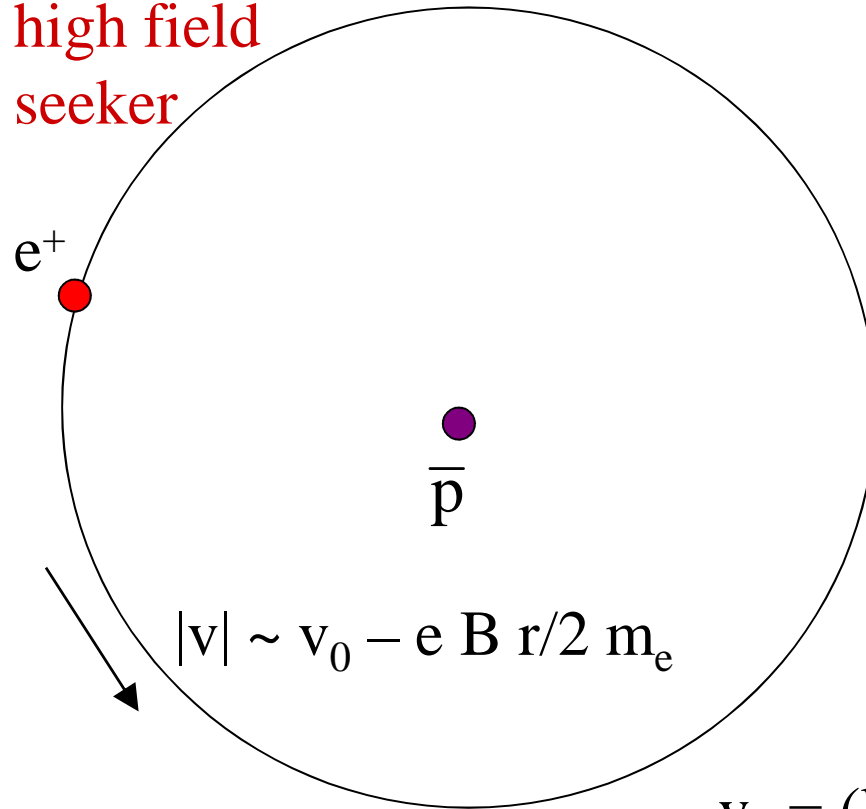
$$= k e/r^2 B$$

$|PE| < KE_{\text{cyc}}$   
means GCA low  
field seeker

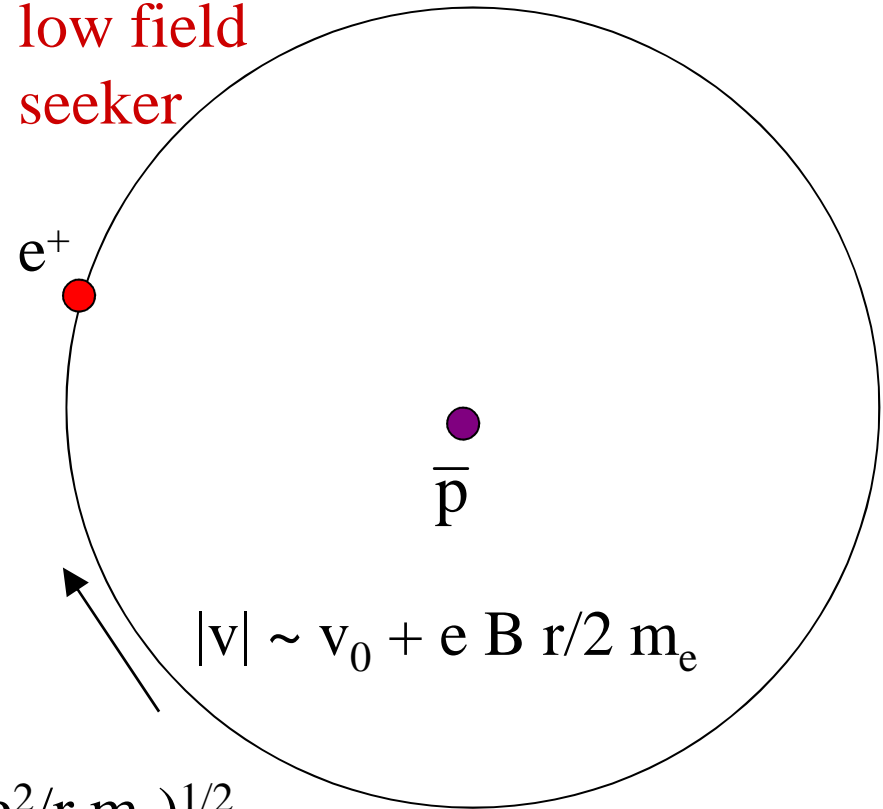
# High/Low Field Seekers

B does not strongly affect  $\mu$  for more deeply bound states ( $n < 30$ ).

high field  
seeker



low field  
seeker



$$v_0 = (k e^2 / r m_e)^{1/2}$$

$$v_0 = 73 \text{ km/s}$$

$$e B r / 2 m_e = 4 \text{ km/s}$$

for  $n = 30$ ,  $B = 1 \text{ T}$

$$|\mu| \sim e^2 (k r / m_e)^{1/2} / 2$$

$$= e |L| / 2 m_e$$

# How many antiprotons?

Should you want to have as many antiprotons as possible?

Argument for: More antiprotons means you will have more antiH and, thus, more trappable antiH.

Argument against: More antiprotons leads to more scattering between antiprotons (**heating**) and you could get charge exchange (**trapped antiH converts to untrapped!**)

# Charge Transfer Always to Untrapped

The potential wells for antiprotons mean that the antiprotons are mostly high  $KE \gg$  well depth.



# Estimate

Energy  $\sim$  couple eV gives  $v \sim 2 \times 10^4$  m/s

Geometric size atom  $\sim 2 n^2 a_0 \sim 3 \times 10^{-7}$  m for  $n \sim 50$

Time to decay  $\sim 0.1$  s for  $n \sim 50$

Transfer probability  $\sim$  density  $\times 10 (3 \times 10^{-7} \text{ m})^2 2 \times 10^4 \text{ m/s} \times 0.1 \text{ s} \sim$  density  $\times 2 \times 10^{-9} \text{ m}^3$

ALPHA density  $\sim 10^4 / (\pi [0.02 \text{ m}]^2 0.2 \text{ m}) \sim 4 \times 10^7 \text{ m}^{-3}$

ATRAP density  $\sim 0.5 \times 10^6 / (\pi [0.02 \text{ m}]^2 0.1 \text{ m}) \sim 4 \times 10^9 \text{ m}^{-3}$  (from recent PRL, I probably overestimate)