

1. a) From notes,  $I_0 = \frac{2P}{\pi \omega_0^2} = \frac{0.2 \text{ W}}{\pi \times 10^{-6} \text{ m}^2} = \boxed{6.3 \times 10^4 \frac{\text{W}}{\text{m}^2}}$   
(2 pts)

b) Generally have

$$I = \frac{E^2}{2\epsilon_0}$$

$$E = \sqrt{2\epsilon_0 I} = \sqrt{2 \times 377 \Omega \times 6.3 \times 10^4 \frac{\text{W}}{\text{m}^2}} = \boxed{6.93 \times 10^3 \frac{\text{V}}{\text{m}}}$$
  
1 pt

c)  $B = \frac{E}{c} = \boxed{2.3 \times 10^{-5} \text{ T}}$   
1 pt

d)  $n = \frac{u}{\hbar \omega}$       $u = \text{energy density} = \frac{I}{c}$

$$\hbar \omega = (1.054 \times 10^{-34} \text{ J}\cdot\text{s}) \left( \frac{2\pi c}{\lambda} \right)$$

$$= 3.14 \times 10^{-19} \text{ J} \quad \text{1 pt}$$

$$\omega = 2.98 \times 10^{15} \frac{\text{rad}}{\text{s}}$$

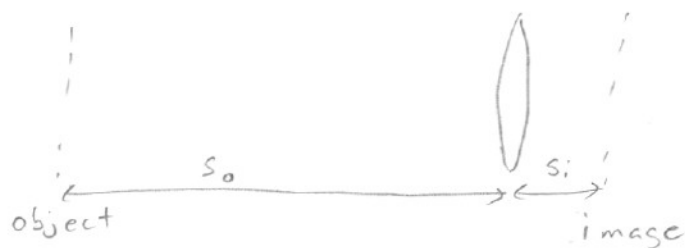
$$\nu = 4.7 \times 10^{14} \text{ Hz}$$

$$n = \frac{6.3 \times 10^4 \frac{\text{W}}{\text{m}^2}}{c} = 2.1 \times 10^{-4} \frac{\text{J}}{\text{m}^3}$$

$$n = 6.7 \times 10^{14} \frac{\text{photons}}{\text{m}^3}$$

2. Want magnification =  $\frac{7 \text{ mm}}{5 \text{ cm}} = 0.14 = \left| \frac{s_i}{s_o} \right|$  (2 pts)

Looks like



With  $s_o = \frac{s_i}{0.14} = 7.14 s_i$

Since  $s_o \gg s_i$ , have  $s_i \approx f$   
 $s_o \approx 7f$

Need  $s_o > 15 \text{ cm} \Rightarrow f > \frac{15}{7} \text{ cm} \approx 2 \text{ cm}$

So 12 mm lens won't work, but 26 mm lens should.

With that lens, want

$$\frac{1}{s_i} + \frac{1}{s_o} = \frac{1}{f}$$

$$\frac{1}{s_i} + \frac{0.14}{s_i} = \frac{1}{f}$$

$$s_i = 1.14 f$$

$$= 29.6 \text{ mm}$$

$$s_o = 7.14 s_i = 21.1 \text{ cm}$$

So lens is 21.1 cm from board,

and sensor is 2.96 cm from lens

3. From tables:

$$\delta(x-x_0) \rightarrow e^{-ik_x x_0}$$

$$e^{ixx} \rightarrow 2\pi \delta(k_x - \alpha)$$

and

$$\begin{array}{l} 1 \\ 0 \end{array} \begin{array}{l} (|y| < b) \\ \text{else} \end{array} \rightarrow 2b \operatorname{sinc} k_y b$$

So  $f(x, y) \rightarrow$

$$F(k_x, k_y) = 2b [A e^{-ik_x x_0} + 2\pi B \delta(k_x - \alpha)] \operatorname{sinc} k_y b$$

4. If  $z=0$  at front of plate, then at  $z=t_0+a$ , have

$$E = E_0 e^{inkt} e^{ik(t_0+a-t)}$$

$\uparrow$  transmission through glass       $\uparrow$  transmission through air

$$= E_0 e^{ink(t_0+a \sin \beta x)} e^{ik(t_0+a-t_0-a \sin \beta x)}$$

$$= E_0 e^{inkt_0} e^{ika} e^{i(n-1)ka \sin \beta x}$$

$$= E_0' e^{i(n-1)ka \sin \beta x}$$

$$E_0' = E_0 e^{ik(nt_0+a)}$$

But note  $ka \ll 1$

$$E \approx E_0' [1 + i(n-1)ka \sin \beta x]$$

$$= E_0' \left\{ 1 + \frac{(n-1)ka}{2} [e^{i\beta x} + e^{-i\beta x}] \right\}$$

Sum of harmonic functions

Each evolves with phase  $e^{i\alpha z}$        $\alpha = \sqrt{k^2 - \beta^2}$

So

$$E(d) = E_0' \left[ e^{ikd} + \frac{(n-1)ka}{2} (e^{i\beta x} + e^{-i\beta x}) e^{i\sqrt{k^2 - \beta^2} d} \right]$$

$$E(d) = E_0' \left[ e^{ikd} + i(n-1)ka \sin \beta x e^{i\sqrt{k^2 - \beta^2} d} \right]$$

5. Know that circular aperture of radius  $a$   
 produces field

$$E = -\frac{i}{\lambda d} e^{i\phi} \pi a^2 \frac{2 J_1\left(\frac{k\rho a}{d}\right)}{k\rho/d} E_0$$

$$\phi = kd + k \frac{\rho^2}{2d}$$

Can write field from ring as  $E_2 - E_1$

$E_2 =$  field from aperture radius  $R_2$

$E_1 =$  field from aperture radius  $R_1$

Since  $E_1$  will be "missing"

$$E_{TOT} = -\frac{i}{\lambda d} e^{i\phi} \frac{2\pi d}{k\rho} \left[ R_2 J_1\left(\frac{k\rho R_2}{d}\right) - R_1 J_1\left(\frac{k\rho R_1}{d}\right) \right] E_0$$

Interference pattern

$$\boxed{|E_{TOT}|^2 = \frac{1}{\rho^2} \left[ R_2 J_1\left(\frac{k\rho R_2}{d}\right) - R_1 J_1\left(\frac{k\rho R_1}{d}\right) \right]^2 |E_0|^2}$$

6. Matrix for half wave plate

$$M = \begin{bmatrix} \cos 2\theta & \sin 2\theta \\ \sin 2\theta & -\cos 2\theta \end{bmatrix}$$

$$\hat{J}_{in} = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 \\ i \end{bmatrix} = \text{LCP}$$

$$\text{So } \hat{J}_{out} = \frac{1}{\sqrt{2}} \begin{bmatrix} \cos 2\theta + i \sin 2\theta \\ \sin 2\theta - i \cos 2\theta \end{bmatrix}$$

$$= \frac{1}{\sqrt{2}} \begin{bmatrix} e^{i2\theta} \\ -ie^{i2\theta} \end{bmatrix}$$

$$= \frac{1}{\sqrt{2}} e^{i2\theta} \begin{bmatrix} 1 \\ -i \end{bmatrix}$$

Overall phase is irrelevant, so

$$\hat{J}_{out} = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 \\ -i \end{bmatrix} = \boxed{\text{RCP}}$$

7.a) At normal incidence, Fresnel equations

become

$$r = \frac{n_i - n_t}{n_i + n_t}$$

$$r_{1a} = \frac{n_1 - n_a}{n_1 + n_a}$$

$$r_{a2} = \frac{n_a - n_2}{n_a + n_2}$$

Supposed to be equal, so

$$\frac{n_1 - n_a}{n_1 + n_a} = \frac{n_a - n_2}{n_a + n_2} \rightarrow +2$$

$$(n_1 - n_a)(n_2 + n_a) = (n_a - n_2)(n_a + n_1)$$

$$n_1 n_2 - n_a n_2 + n_1 n_a - n_a^2 = n_a^2 - n_2 n_a + n_a n_1 - n_1 n_2$$

$$n_1 n_2 - n_a^2 = n_a^2 - n_1 n_2$$

$$n_a^2 = n_1 n_2$$

$$n_a = \sqrt{n_1 n_2}$$

b) Reflection from  $a \rightarrow 2$  interface travels extra distance  $2h$

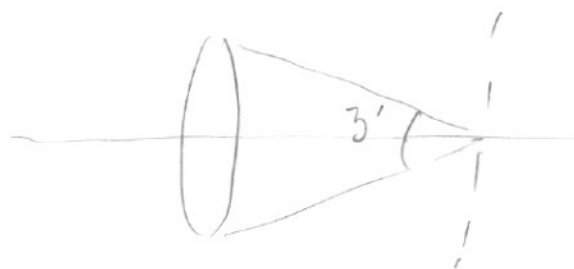
To get  $\pi$  phase shift, need  $2h = \lambda \left( \frac{1}{2} + m \right)$   
 $m = \text{integer}$

$$h = \frac{\lambda}{2} \left( \frac{1}{2} + m \right)$$

But need  $\lambda$  in medium =  $\frac{\lambda_0}{n_a}$

$$h = \frac{\lambda_0}{2n_a} \left( \frac{1}{2} + m \right) \quad \text{integer } m$$

8. From image plane, looks like light coming from source with angle  $\theta' = \frac{D}{f}$ :



So, have  $\rho_c \approx \frac{\lambda}{\theta'} = \boxed{\frac{\lambda f}{D}}$

Or: each point on source  $\rightarrow$  diffraction pattern with  $\Delta x \approx \frac{\lambda f}{D}$

Within this  $\Delta x$ , can't distinguish which source point light is from  $\Rightarrow$  field is coherent.

Or from van Cittert-Zernike,  $\rho_c \approx$  width of diffraction pattern =  $\frac{\lambda f}{D}$