1. Find the Fourier transform of the function $f(x)$ where:

$$
f(x)= \begin{cases}0 & (\text { if } x<-a / 2) \\ e^{i \beta x} & \text { (if }-a / 2<x<a / 2) \\ 0 & \text { (if } a / 2<x<b-a / 2) \\ e^{i \beta x} & (\text { if } b-a / 2<x<b+a / 2) \\ 0 & (\text { if } b+a / 2<x)\end{cases}
$$

Here $a$ and $b$ are positive real constants with $b>a$.
2. Calculate the convolution

$$
g(t)=\int_{-\infty}^{\infty} f_{1}(T) f_{2}(t-T) d T
$$

for

$$
\begin{aligned}
& f_{1}(t)= \begin{cases}1 & \text { (if } \left.-\tau_{1}<t<\tau_{1}\right) \\
0 & \text { (otherwise) }\end{cases} \\
& f_{2}(t)= \begin{cases}1 & \text { (if } \left.-\tau_{2}<t<\tau_{2}\right) \\
0 & \text { (otherwise) }\end{cases}
\end{aligned}
$$

where $\tau_{1}>\tau_{2}$. Sketch a plot of $g(t)$, and note its full-width at half-maximum.
Hint: you will need to separately consider the cases (a) $t<-\left(\tau_{1}+\tau_{2}\right)$; (b) $-\left(\tau_{1}+\tau_{2}\right)<t<\tau_{2}-\tau_{1}$; (c) $\tau_{2}-\tau_{1}<t<\tau_{1}-\tau_{2}$; (d) $\tau_{1}-\tau_{2}<t<\tau_{1}+\tau_{2}$; (e) $\tau_{1}+\tau_{2}<t$.
3. For each of the following aperture functions, calculate the propagating wave $E(x, y, z)$, assuming light with wave number $k$ :
(a) $A(x, y)=e^{i k y / 4}$
(b) $A(x, y)=e^{-i k(x+y) / 2}$
(c) $A(x, y)=e^{i k(x+y)}$
(d) $A(x, y)=\sin ^{2}(k y / 4)$
4. Suppose that light of wavelength $\lambda=633 \mathrm{~nm}$ can be described in the $z=0$ plane by an aperture function $A(x, y)$. It is known that $A(x, y)$ has a Fourier spectrum which is zero for $\sqrt{k_{x}^{2}+k_{y}^{2}}>10^{6} \mathrm{~m}^{-1}$. Show that the Fraunhofer diffraction pattern will be contained within a cone and calculate the cone angle $\theta$. (See picture on reverse.)

## Problem 4:


5. The Fourier method we have developed for diffraction can be applied to other problems as well. For instance, consider a pulse of light $E(z, t)$ incident on a dispersive medium with index $n(\omega)$ and length $d$. The front face of the medium is at $z=0$, and the incident pulse is

$$
E(z=0, t) \equiv A(t)=E_{0} e^{-t^{2} / \tau^{2}} e^{-i \omega_{0} t}
$$

for constant pulse duration $\tau$ and carrier frequency $\omega_{0}$. We can ask: what is the form of the pulse when it exits the medium?

Using the Fourier transform, the incident pulse can be written as a sum of harmonic functions $e^{-i \omega t}$. Each of these components propagates through the medium as $e^{i(k z-\omega t)}$ an thus exits after after acquiring a phase shift $\mathcal{H}(\omega)=e^{i k d}$ for $k=k(\omega)=$ $n \omega / c$. Symbolically,

$$
\text { (input) } e^{-i \omega t} \rightarrow e^{-i \omega t} e^{i k d} \text { (output) }
$$

By resumming the transmitted harmonic functions, the complete transmitted pulse can be calculated.

Suppose the wavenumber $k$ can be approximated as

$$
k\left(\omega_{0}+\Delta\right)=k_{0}+k_{1} \Delta+\frac{1}{2} k_{2} \Delta^{2}
$$

where $\Delta=\omega-\omega_{0}, k_{0}=k\left(\omega_{0}\right), k_{1}=d k / d \omega$, and $k_{2}=d^{2} k / d \omega^{2}$. Use the procedure outlined above to calculate the electric field of the transmitted pulse. In particular, find:
(a) the time at which the peak of the pulse exits the medium, and
(b) the duration of the pulse when it exits the medium (ie, the new value of $\tau$ ).

Hint: The Gaussian beam example from lecture 15, slides 32-35 illustrates a similar calculation.

