Phys 531 Lecture 7 23 September 2004

Total Internal Reflection & Metal Mirrors

Last time, derived Fresnel relations

Give amplitude of reflected, transmitted waves at boundary

Focused on simple boundaries:  $air \rightarrow glass$ 

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Today, consider more complicated situations

- Total internal reflection
  - Evanescant waves
- Materials with complex index

This will wrap up unit on fundamental theory

Next time: ray optics

# Total Internal Reflection (Hecht 4.7)

So far, considered  $n_i < n_t$ 

If  $n_i > n_t$ , problem with Snell's Law:

$$\sin \theta_t = \frac{n_i}{n_t} \sin \theta_i$$

What if  $(n_i/n_t) \sin \theta_i > 1$ ?

For instance, glass  $(n_i = 1.5) \rightarrow \text{air } (n_t = 1)$ :

if 
$$\theta_i > 41.8^\circ$$
, then  $\sin \theta_i > \frac{1}{1.5}$ 

Demo!

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For  $\sin \theta_i > \sin \theta_c \equiv n_t/n_i$ , all light is reflected

Called total internal reflection = TIR only occurs when light exits medium (high  $n \rightarrow low n$ )

Would like to understand from Fresnel relations

$$r_{\perp} = \frac{n_i \cos \theta_i - n_t \cos \theta_t}{n_i \cos \theta_i + n_t \cos \theta_t}$$

$$r_{\parallel} = \frac{n_t \cos \theta_i - n_i \cos \theta_t}{n_i \cos \theta_t + n_t \cos \theta_i}$$

How can we use? Don't have a  $\theta_t$ !

Trick: use complex  $\theta_t$ 

Say  $\theta_t = a + ib$ 

Considered  $\cos \theta_t$  in homework 1 Look at  $\sin \theta_t$  now

Use

$$\sin(a+ib) = \sin(a)\cos(ib) + \cos(a)\sin(ib)$$

Just need sin(ib), cos(ib)

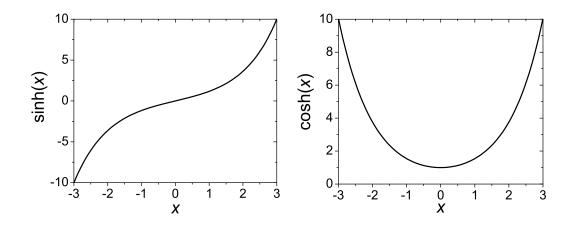
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Euler identity gives: 
$$\sin(x) = \frac{1}{2i} \left( e^{ix} - e^{-ix} \right)$$
  
 $\cos(x) = \frac{1}{2} \left( e^{ix} - e^{-ix} \right)$ 

So 
$$\sin(ib) = \frac{1}{2i} \left( e^{i(ib)} - e^{-i(ib)} \right)$$
$$= -\frac{i}{2} \left( e^{-b} - e^{b} \right)$$

Hyperbolic sine: 
$$\sinh(x) \equiv \frac{1}{2} \left( e^x - e^{-x} \right)$$
  
So  $\sin(ib) = i \sinh(b)$ 

Hyperbolic sine and cosine:



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Also

$$\cos(ib) = \frac{1}{2} \left( e^{i(ib)} + e^{-i(ib)} \right)$$
$$= \frac{1}{2} \left( e^{-b} + e^{b} \right)$$
$$\equiv \cosh(b)$$

Hyperbolic cosine

So: sin(a + ib) = sin(a) cosh(b) + i cos(a) sinh(b)

For large  $b, \, \sinh b, \cosh b \to e^b$  no problem satisfying  $n_i \sin \theta_i = n_t \sin \theta_t$ 

Want  $n_i \sin \theta_i = n_t \sin \theta_t$  with complex  $\theta_t$ 

Assume  $n_i \sin \theta_i$  and  $n_t$  real:

Then Im  $[n_t \sin \theta_t] = n_t \cos(a) \sinh(b) = 0$ 

Don't want b=0, so take  $a=\frac{\pi}{2}$ 

Gives  $\sin \theta_t = \cosh b$ 

Snell's law becomes

$$n_i \sin \theta_i = n_t \cosh b$$

Note  $\cosh b > 1$ , require  $\sin \theta_i > n_t/n_i - \sin \theta_c$ 

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**Example:** What is the complex transmission angle for light propagating from glass to air with an angle of incidence of 60°?

If  $\theta_i = 60^{\circ}$  then  $b = \cosh^{-1}[1.5\sin(60^{\circ})] = 0.755$ 

So 
$$\theta_t = \frac{\pi}{2} + 0.755i$$

Just need a good calculator!

**Question:** What are the units of b?

What happens to Fresnel coefficients?

Need 
$$\cos \theta_t = \cos \left( \frac{\pi}{2} + ib \right)$$
  
=  $-i \sinh(b)$ 

pure imaginary

So 
$$r_{\perp} = \frac{n_i \cos \theta_i - n_t \cos \theta_t}{n_i \cos \theta_i + n_t \cos \theta_t}$$

$$\rightarrow \frac{n_i \cos \theta_i - i n_t \sinh b}{n_i \cos \theta_i + i n_t \sinh b} \quad \text{complex!}$$

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Still have  $E_{r0}=rE_{i0}$  complex r: phase shift between  ${f E}_{\rm inc}$  and  ${f E}_{\rm refl}$  Both  $r_{\perp}$  and  $r_{\parallel}$  have form

$$\begin{split} r &= \frac{u+iv}{u-iv} \\ r_{\perp} : u &= n_i \cos \theta_i \text{ and } v = n_t \sinh b \\ r_{\parallel} : u &= n_t \cos \theta_i \text{ and } v = n_i \sinh b \end{split}$$

If z = u + iv, then  $r = z/z^*$ 

So 
$$|r| = \frac{|z|}{|z^*|} = 1$$
: all light reflected

Write  $z=|z|e^{i\phi}$ , then  $r=e^{2i\phi}$ 

Reflection phase shift:  $2\phi = 2 \tan^{-1} \left(\frac{v}{u}\right)$ 

To calculate numerically, just use good calculator (or computer program):

Evaluate

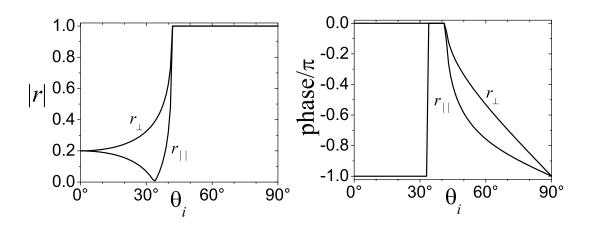
$$r_{\perp} = \frac{n_i \cos \theta_i - n_t \cos \theta_t}{n_i \cos \theta_i + n_t \cos \theta_t}$$

for 
$$\theta_t = \sin^{-1}\left(\frac{n_i}{n_t}\sin\theta_i\right)$$

Let computer deal with complex math

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Plot for glass  $\rightarrow$  air:



Use TIR to make good mirror be aware of phase shifts

For TIR, 
$$R = |r|^2 = 1$$
 so all power reflected

But also have

$$t_{\perp} = \frac{2n_i \cos \theta_i}{n_i \cos \theta_i + n_t \cos \theta_t}$$

$$t_{\parallel} = \frac{2n_i \cos \theta_t}{n_i \cos \theta_t + n_t \cos \theta_i}$$

do not equal zero!?

Have a transmitted field  $(t \neq 0)$ but doesn't carry any energy (R = 1)

Question: How might this contradiction be resolved?

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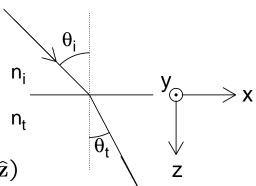
To see:

Transmitted wave:  $\mathbf{E}_t = \mathbf{E}_{t0}e^{i(\mathbf{k}_t \cdot \mathbf{r} - \omega t)}$ 

with  $\mathbf{k}_t = |k_t|(\sin\theta_t \hat{\mathbf{x}} + \cos\theta_t \hat{\mathbf{z}})$ 

For TIR, 
$$\theta_t = \frac{\pi}{2} + b$$
  $\sin \theta_t \rightarrow \cosh b$   $\cos \theta_t \rightarrow i \sinh b$ 

So  $\mathbf{k}_t \to k_t (\cosh b \hat{\mathbf{x}} + i \sinh b \hat{\mathbf{z}})$ 



Transmitted field

$$\mathbf{E}_t \to \mathbf{E}_{t0} e^{i[k_t(x\cosh b + iz\sinh b) - \omega t]}$$
$$= \mathbf{E}_{t0} e^{-k_t z \sinh b} e^{i(k_t x \cosh b - \omega t)}$$

Wave propagates in x direction Decays exponentially in z direction

• Carries no energy away from surface

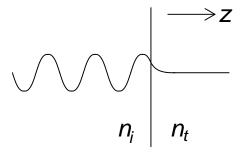
Called evanenscent wave

(Not same as exponential decay from absorption!)

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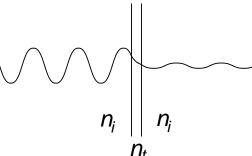
Evanescent wave can be observed

Single surface:



Introduce

second surface:



Transmitted wave appears!

Get transmission from tail of evanescent wave

For gap d, amplitude of transmitted wave  $\approx e^{-k_t d \sinh b}$ 

ullet reflection o 0 smoothly as d o 0

Called frustrated total internal reflection

Completely analogous to tunneling in QM

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Still hope T = 0 for plain TIR

Should have defined 
$$T = \operatorname{Re} \left[ \frac{n_t \cos \theta_t}{n_i \cos \theta_i} \right] |t|^2$$

In TIR, 
$$\cos \theta_t = -i \sinh b$$
  
real part = 0

So T = 0, as needed

Raises question: what if  $n_i$  or  $n_t$  is complex?

## Reflection from metals (Hecht 4.8)

Saw previously that in absorbing medium

$$n \to n + i \frac{\alpha}{2k_0}$$

 $\alpha$  = absorption coefficient

Get 
$$\mathbf{E} = \mathbf{E}_0 e^{i(n\mathbf{k}_0 \cdot \mathbf{r} - \omega t)}$$
  
 $\rightarrow \mathbf{E}_0 e^{-\alpha \hat{\mathbf{k}} \cdot \mathbf{r}/2} e^{i(n\mathbf{k} \cdot \mathbf{r} - \omega t)}$ 

and 
$$I \propto |E_0|^2 \propto e^{-\alpha \hat{\mathbf{k}} \cdot \mathbf{r}}$$

Wave attenuates as it propagates

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**Question:** How could you distinguish an evanescent wave from TIR and a plane wave exponentially decaying due to absorption? (Supposing no knowledge about the media.)

Normally don't want optical material to absorb

Important exception: mirrors light doesn't penetrate medium not much loss

How to apply Fresnel relations? Typically  $n_i$ ,  $\theta_i$  real  $n_t$  complex

Same equations apply

- Snell's law:  $\sin \theta_t = \frac{n_i}{n_t} \sin \theta_i$  (so  $\theta_t$  complex)
- Use  $\cos \theta_t = \sqrt{1 \sin^2 \theta_t}$  as before
- $\bullet$  Plug into equations for  $r_{\perp}$  ,  $r_{\parallel}$  Get complex result

Hard to do by hand; easy on computer

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Consider a very good absorber  $\alpha \to \infty$ 

Then 
$$n_i \sin \theta_i = \left(n_t + i \frac{\alpha}{2k_0}\right) \sin \theta_t$$
 means  $\theta_t \to 0$  and  $r_\perp = \frac{n_i \cos \theta_i - n_t \cos \theta_t}{n_i \cos \theta_i + n_t \cos \theta_t}$  
$$\to \frac{n_i \cos \theta_i - n_t - i \alpha/2k_0}{n_i \cos \theta_i + n_t + i \alpha/2k_0}$$
 
$$\to -1$$

Similarly  $r_{\parallel} 
ightarrow 1$ 

So perfect absorber = perfect reflector

To make a true absorber: need moderate  $\alpha$  + porous surface reflected waves bounce many times Example: soot ( = carbon dust)

Use high- $\alpha$  material for mirror

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Highest  $\alpha$ : metals

Homework problem 5 from assignment 1: In conductive medium get current  ${\bf J}=\sigma {\bf E}$   $\sigma=$  conductivity

$$Got k = \sqrt{\epsilon \mu_0 \omega^2 + i\omega \mu_0 \sigma}$$

Rewrite

$$k = k_0 \sqrt{\frac{\epsilon}{\epsilon_0} + i \frac{\sigma}{\epsilon_0 \omega}}$$

Good conductor: silver  $\sigma \approx 6 \times 10^7 \; (\Omega \; \text{m})^{-1} \; \; (\text{at dc})$ 

If 
$$\lambda = 500$$
 nm and  $\epsilon \approx \epsilon_0$  
$$\frac{\sigma}{\epsilon_0 \omega} \approx 2000$$

So 
$$k \approx k_0 \sqrt{2000i}$$
  
  $\approx 45k_0 \frac{1+i}{\sqrt{2}} = 30k_0(1+i)$ 

Then expect  $n \approx \alpha/2k_0 \approx 30$ 

Actually, not that good at optical freqs find n=0.3 and  $\alpha/2k_0=4$  Still get  $R\approx 0.95$  across visible

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### Practical notes

• Typical metals:

Silver:  $R \approx 0.95$  in visible, NIR

- oxidizes quickly in air

Gold:  $R \approx 0.95$  in NIR

- doesn't oxidize

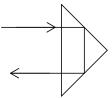
Aluminum:  $R \approx 0.85$  in visible, NIR

- oxidizes but easy to protect (SiO)
- ullet Metals don't have Brewster angle typically dip in  $R_{||}$ , but not to zero
- ullet Wave penetrates fraction of  $\lambda$  typical 50-100 nm

Get better mirrors using dielectric layers

- discuss theory later
- can get R = 0.99 easily, 0.99999 with effort
- more expensive than metal

#### Could use TIR:



#### Drawbacks:

- Reflection losses from first surface
- Beam displacement inconvenient
- Limited range of  $\theta_i$

Usually use when displacement desired

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### Summary

- ullet Get TIR with internal incidence,  $heta_i > heta_c$
- Perfect reflection, with phase shift
- Evanescant wave at surface
- For TIR or absorbing media,
   Fresnel equations are complex
- Highly absorbing medium → good mirror