

Photons

Last time, finished polarization

Learned about retarders,
Jones matrices

Today: introduce quantum optics
Whole new way to look at light

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Outline:

- Photon optics
- Photon detectors
- Quantum noise
- Quantum states

Photon optics not too hard

Helpful to know quantum mechanics
not totally necessary

Next time:

Quantum field theory of light

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Photon Optics (Hecht 3.3.3)

Simple version of quantum theory

Say that light is really composed of particles
particles called *photons*

Energy of photon = $\hbar\omega$

\hbar = Planck's constant = 1.054×10^{-34} J s

ω = oscillation frequency of light

(Don't worry about what is oscillating for now)

Polychromatic light:

many photons with different ω 's

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- Pulse of light with energy U (in J)

contains $N = \frac{U}{\hbar\omega}$ photons

N = *photon number*

- Beam with power P (in W = J/s)

delivers $\Phi = \frac{P}{\hbar\omega}$ photons/s

Φ = *photon flux*

- Beam with irradiance $I(\mathbf{r})$ (in W/m²)

has $\phi = \frac{I}{\hbar\omega} \frac{\text{photons/s}}{\text{m}^2}$ at \mathbf{r}

ϕ = *photon flux density*

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Also light with energy density $u(\mathbf{r})$ (in J/m^3)

has photon density $\frac{u}{\hbar\omega}$ $\left(\frac{\text{photons}}{\text{m}^3}\right)$

Example:

Sunlight has an irradiance of about $250 \text{ W}/\text{m}^2$

Assume average wavelength = 500 nm

Then average $\hbar\omega = 4 \times 10^{-19} \text{ J}$

Photon flux density = $6 \times 10^{20} \text{ photons}/(\text{s m}^2)$

Looking at sun, pupil area $\approx 10^{-6} \text{ m}^2$

collect about 6×10^{14} photons in 1 s

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Photons are not “classical” particles!

→ propagate according to wave equation
not Newton’s laws

Procedure: use wave techniques to calculate $I(\mathbf{r})$

Then I gives flux density

Imagine photons “follow” wave
like surfers in ocean

But photons and wave are inseparable

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Best to interpret probabilistically:

Average number of photons in volume $d^3r =$

$$\langle N \rangle = \frac{u(\mathbf{r}) d^3r}{\hbar\omega} = \frac{I(\mathbf{r}) d^3r}{\hbar\omega c}$$

If $\langle N \rangle \ll 1$, interpret as probability to find one photon

Important:

No definite trajectory for individual photons

Picture surfers scrambling back and forth on wave

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Avoids “two-slit” paradox:

In two slit interferometer, which slit does photon pass through?

Correct interpretation:

It doesn't matter

The wave passes through both

Interference pattern says where photons can go

Question: What if a wave passes through beam splitter, and the outputs are separated by a large distance? Does it make sense to ask where one of the photon is?

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If wave determines photon distribution,
why use photons?

Because detectors see photons, not waves

Or: light can only transfer energy in units of $\hbar\omega$

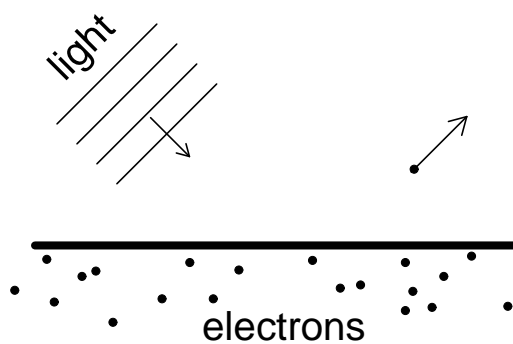
So photons important whenever light is emitted or absorbed

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Example: photoelectric effect

Shine light on metal:

Electrons absorb energy from light and escape



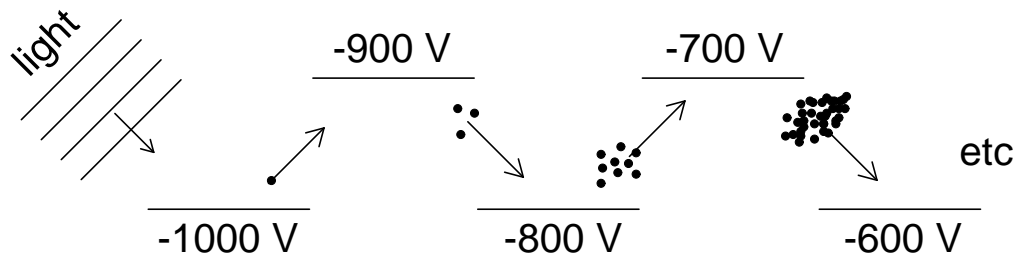
Find that maximum electron energy = $\hbar\omega$
= max energy absorbed from photon

Doesn't depend on irradiance, just ω

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Another example: photomultiplier tube

Light hits metal plate, detaches single electron



Accelerate electron to second plate:

detach more electrons

Cascade many plates, get big current pulse

See blip on output: detection of one photon

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Other photon properties:

- Momentum $p = \hbar k$ (Hecht 3.3.4)

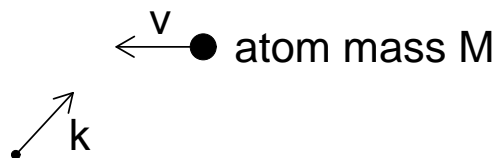
Interesting effect: suppose atom absorbs photon

makes internal transition $U_0 \rightarrow U_1$

Say atom velocity before absorption = \mathbf{v}

Velocity after absorption = $\mathbf{v} + \hbar \mathbf{k} / M$

gets kick from photon



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Initial energy: $U_{\text{tot}} = U_0 + \frac{1}{2}Mv^2 + \hbar\omega$

Final energy: $U_{\text{tot}} = U_1 + \frac{1}{2}M \left(\mathbf{v} + \frac{\hbar\mathbf{k}}{M} \right)^2$

Energy conserved, so

$$\omega = \frac{U_1 - U_0}{\hbar} + \frac{\hbar k^2}{2M} + \mathbf{v} \cdot \mathbf{k}$$

First term: standard QM

Second term: “radiative correction”

Third term: Doppler shift

Derive Doppler effect from QM!

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- Angular momentum (Hecht 8.1.5)

Circular polarization states have definite spin

Right circular: $\mathbf{L}_{\text{photon}} = -\hbar\hat{\mathbf{k}}$

Left circular: $\mathbf{L}_{\text{photon}} = +\hbar\hat{\mathbf{k}}$

Linear polarized states

= superposition of $\hat{\mathbf{e}}_{\mathcal{R}}$ and $\hat{\mathbf{e}}_{\mathcal{L}}$

Photon equally likely to be “spin-up” and “spin-down”

Important for transition selection rules

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Detecting Photons

Already mentioned one way to see single photons:
photo-multipier tube (PMT)

Are there other methods?

Can characterize light detectors by three parameters:

- quantum efficiency
- background rate
- saturation rate

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Quantum efficiency $\varepsilon =$
probability that incident photon is detected

For PMT, typically about 5%
- limited by emission of first electron

Background rate $R_0 =$
count rate observed with no incident light

For good PMT, as low as ~ 1 count/s
- due to thermal emission of electrons

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Saturation rate $R_{\text{sat}} \approx$
max measurable count rate

For PMT, about 10^6 counts/s
- takes about $1 \mu\text{s}$ to recover between 'clicks'

To see single photons, need incident flux Φ with

$$R_0 < \epsilon\Phi < R_{\text{sat}}$$

Impossible if $R_0 \geq R_{\text{sat}}$

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Some detectors:

Device	ϵ (%)	R_0 (counts/s)	R_{sat} (counts/s)
PMT	1–20	< 1 to 10^5	$10^6 - 10^9$
Photodiode	20–90	10^4+	$< R_0$
Avalanche PD	30–70	10 to 5000	$10^6 - 10^7$

(Note PMTs have much larger area than APDs)

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Quantum Noise

Important effect of photon theory:
light is noisy

Suppose laser beam with power P
measure with ideal photon counter

In time T , detect $N = PT/\hbar\omega$ photons

But photons are randomly distributed in wave

Don't expect to get exactly N every time

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Guess that photons are distributed independently:

Probability to detect a photon doesn't depend
on when previous photon was detected

Why?

- Photons emitted by independent atoms
- Don't expect correlations in emission times

(Reconsider this assumption next time)

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Then photons obey *Poisson statistics*

If on average you detect N photons in time T

then N fluctuates by $\Delta N = N^{1/2}$

Same as fluctuations in flipping coins,
most other statistical noise

Sometimes called “shot noise”:

- comes from photons being discrete “shots”

Better name is “quantum noise”

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So no such thing as light with constant P

Always detect some fluctuations

Size of fluctuations depends on time scale

Measure for time T , get relative noise

$$\frac{\Delta P}{P} = \frac{\Delta N}{N} = \frac{1}{N^{1/2}} = \left(\frac{\hbar\omega}{PT} \right)^{1/2}$$

Noise goes up as T goes down

(just like normal noise goes down with averaging)

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Question: If your detector has a quantum efficiency less than one, will the quantum noise scale as the square root of the number of **incident** photons, or the number of **measured** photons?

All real detectors also have “technical noise”
= noise sources besides quantum mechanics

For instance:

- noise in background signal
- electrical pick-up
- vibration-induced noise
- noise from temperature drifts

etc.

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Most often, technical noise dominates

Quantum noise typically important for:

- Photon counting applications
(N is low, so $\Delta N/N$ is large)
- High speed applications (T is small)
→ need high power to measure fast signals

Can distinguish quantum and technical noise:

Quantum: noise $\propto P^{1/2}$

Technical: noise $\propto P$ or indep of P

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Quantum States

So far, haven't really needed quantum mechanics
just idea that light energy is discrete

But quantum mechanics of light is interesting too

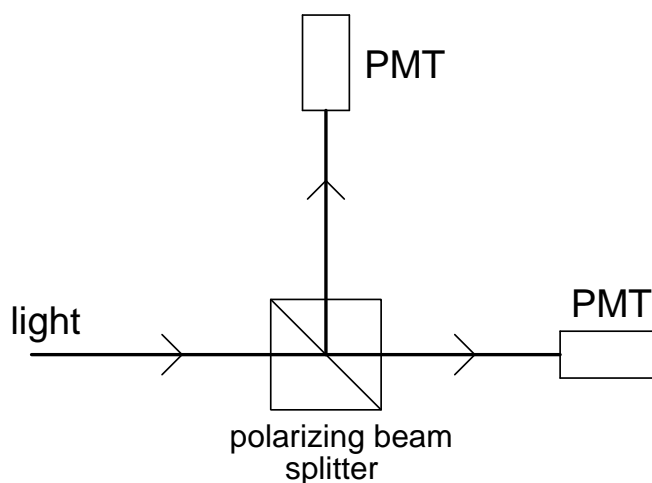
Easiest to see using polarization:

Suppose beam with $\hat{j} = \cos \alpha \hat{x} + \sin \alpha \hat{y}$

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Count number of photons polarized along x, y

For instance:



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If average photon number = N , expect

$$N_x = \cos^2 \alpha N$$

$$N_y = \sin^2 \alpha N$$

on average

If $N \ll 1$, often see no photons

but when one is detected, have prob

$$\cos^2 \alpha \text{ to be along } x$$

$$\sin^2 \alpha \text{ along } y$$

Similar to QM for particle in superposition state

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Natural to describe photon with wave function

$$\psi = \cos \alpha \hat{X} + \sin \alpha \hat{Y}$$

where

$$\hat{X} = \text{state polarized along } x$$

$$\hat{Y} = \text{state polarized along } y$$

Note $\psi \sim \hat{j}$, for right interpretation of $\hat{\mathbf{x}}$, $\hat{\mathbf{y}}$

More generally, have

$$\psi(\mathbf{r}) \sim \mathbf{E}(\mathbf{r})$$

Electric field of wave \approx wave function of photons

Optics and quantum mechanics very similar

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But quantum mechanics has more possibilities

Imagine light pulse with two photons

QM allows polarization state

$$\psi_{12} = \chi \equiv \frac{1}{\sqrt{2}} (\hat{X}_1 \hat{Y}_2 + \hat{Y}_1 \hat{X}_2)$$

Always measure one photon polarized along x
and one along y

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But light not polarized in any usual way

Compare: light polarized at 45° has

$$\begin{aligned} \psi_{12} &= \left[\frac{1}{\sqrt{2}} (\hat{X}_1 + \hat{Y}_1) \right] \left[\frac{1}{\sqrt{2}} (\hat{X}_2 + \hat{Y}_2) \right] \\ &= \frac{1}{2} (\hat{X}_1 \hat{X}_2 + \hat{X}_1 \hat{Y}_2 + \hat{Y}_1 \hat{X}_2 + \hat{Y}_1 \hat{Y}_2) \end{aligned}$$

Have 50% chance to measure both photons in
same state

Not the same as state χ

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States like χ called *nonclassical*:

Wavefunction can't be expressed as an ordinary wave

Generally, don't worry about it:

Almost all light sources produce classical light

(Hard work to make state like χ)

But non-classical states are useful

Tricks like quantum noise suppression, quantum teleportation, quantum computing

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Summary:

- Photons \approx "particles" of light
 - but photons follow wave
- Photon energy $\hbar\omega$, momentum $\hbar k$,
ang. momentum \hbar
- Detectors measure photons, not wave
good detectors count individual photons
- Photon statistics \rightarrow quantum noise
- Photon wave function \approx electric field
 - but nonclassical states possible

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