

02/09/05

Lecture 10

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Next big topic: Laser theory.

Starting with how light interacts with matter.

Matter (atoms): quantum mechanics, (always)

Light: three possible approaches

\* Classical E & M theory: treat light as waves, obeying Maxwell's equations.

Almost sufficient to describe laser, need little input from quantum mechanics.

\* photon theory: semi-classical approach  
simplified quantum approach.

treat light as quantum particles, that are absorbed and emitted in discrete events.

\* Quantum Optics: treat light as quantum wave, coupled to quantum atoms.

This approach is the most accurate, but also hardest.

We'll take photon theory approach. This approach use the concept of photon from quantum mechanical treatment of light, but simplifies things by glossing over details of emission & absorption events.

Recall basic laser idea;  
put fluorescent medium in a cavity,  
every round trip light passes through medium  
→ gets amplified.  
→ study amplification / absorption process.

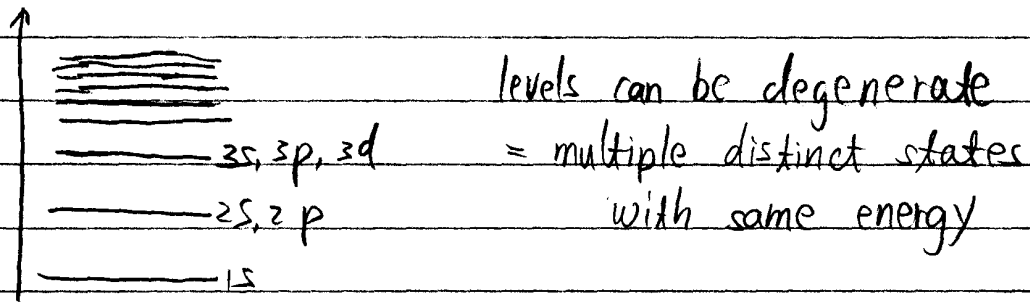
First briefly review quantum levels of matter, and light quanta - photon, then focus on their interactions.

Start with medium.

QM says that a given type of matter has certain allowed values for internal energy called energy levels

Examples

\* Hydrogen atom



# of states = degeneracy =  $g$

Note: only relative degeneracy usually matters

In H, ground state is  $1s$ ,  $g_1 = 1$

But electron has spin,  $g_1 = 1 \times 2 = 2$

But nucleus has spin,  $g_1 = 1 \times 2 \times 2 = 4$

Other levels have same contributions.

for  $n=2$   $g_2 = 2s + 2p_1 + 2p_0 + 2p_{-1} = 4$

consider electron spin  $g_2 = 4 \times 2 = 8$

further more consider nucleus spin  $g_2 = 4 \times 2 \times 2 = 16$

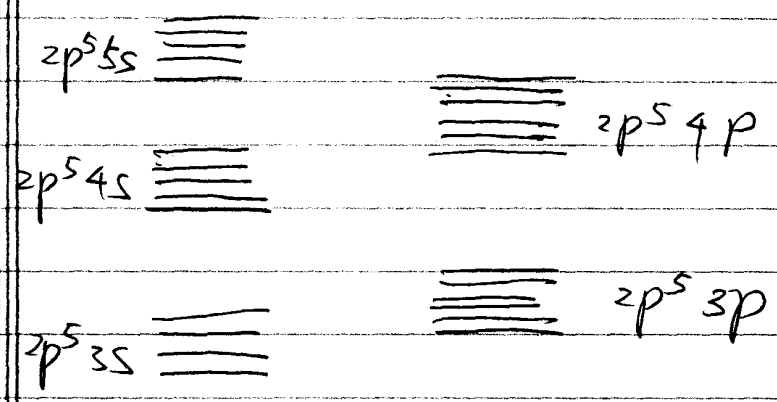
only  $\frac{g_2}{g_1}$  usually matters, so ignore spins

since they make same contribution for both  $n=1$

and  $n=2$  states, and take  $g_1 = 1$

$g_2 = 4$

\* He: more complicated atom, degeneracy is broken

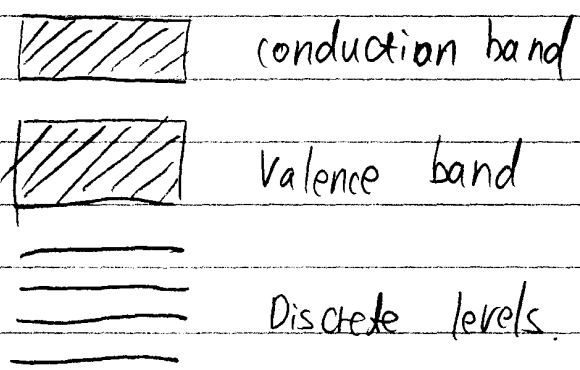


Each level has many nearby "sublevels"

Lump sublevels together or treat separately, depending on energy resolution of process considered.

Narrow linewidth laser: treat separately.  
 Broad linewidth: lump together.

\* Si: solid, have so many nearby states, they can't be distinguished,  $\rightarrow$  energy band



Typically need to integrate over energies in band ...  
 We'll more stick to discrete-level systems.

## Light:

Consists of photons: each photon frequency  $\nu$  has energy  $E = h\nu$

Photons are associated with modes

Given mode contains integer # of photons

$$\text{Mode: } E_{\vec{n}} \sim \sin \frac{n_x \pi x}{d} \sin \frac{n_y \pi y}{d} \sin \frac{n_z \pi z}{d}$$

for cube cavity

$$E_{n,l,m} \sim \text{Hermite-Gaussian mode}$$

for laser cavity

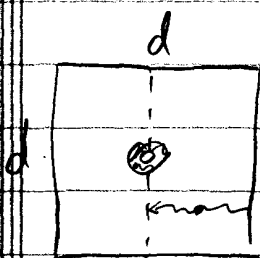
Although we have been talking about Gaussian mode for laser cavity, but to start the discussion of light & matter think about cube cavity mode for now, which is easier since field is uniform. The results should be easy to generalize to Gaussian mode.

Note: free space - limit of cube as  $d \rightarrow \infty$

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## Interactions,

Say a cavity contains both a photon and an atom, what happens?



photon bounces between cavity walls  
 passes through plane of atom  
 at rate  $\frac{c}{d}$  ( $\sim \frac{1}{\lambda} = \frac{1}{d/c}$ )

2D picture of  
 cube cavity

Every time photon passes plane of atom, has probability  $P'$  to be absorbed.

' $P'$ ' depends on size of cavity, smaller cavity  $P'$  larger.

write.  $P' = \frac{\sigma}{A}$

$A$ : cross sectional area  
 of cavity  $= d^2$

$\sigma$ : "cross section" of atom  
 having units of area.

Treat atom as little ball with cross area  $\sigma$ .  
 photon is absorbed if it hits ball.

$$\text{So } p' = \frac{\sigma}{d^2}$$

Absorb at rate  $\frac{c}{d} \cdot \frac{\sigma}{d^2}$  (prob. per unit time)

$$= \frac{c}{V} \cdot \sigma$$

$V$ : cavity volume

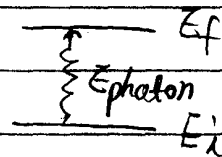
If  $N$  photons, absorption at rate

$$P_{\text{abs}} = N \frac{c}{V} \sigma$$

However  $\sigma = \sigma(\nu)$  is frequency-dependent.

Atom can absorb light if  $E_{\text{photon}} \approx E_f - E_i$   
for two atomic levels  $i$  &  $f$

So  $\sigma(\nu)$  has peaks at atomic resonances.



To absorb a photon at  $\approx E_f - E_i$ , atom need to initially be in state  $i$ , the lower energy state.

Depends on state of atom, there are another two possible, but less obvious process.

## Stimulated emission:

QM is invariant under time reversal, so if atom is initially in higher state  $f$ , can emit photons at rate.

$$P_{\text{stim}} = N \frac{c}{V} \sigma(\nu)$$

$\sigma(\nu)$  same as for absorption, since it's a characteristic of atom.

Note, stimulated emission depends on  $N$ :

presence of photons in a mode "encourages" atom to emit photon into that mode.

## Spontaneous emission:

even if cavity is empty, atom in higher state  $f$  can emit photons at rate

$$P_{\text{spont}} = \frac{c}{V} \sigma(\nu)$$

again, same  $\sigma(\nu)$

can only be clearly explained by QED: zero-point energy.