

04/06/05

Lecture 27.

CD

Electro-optic effect: can be described in terms of impermeability tensor.

$$\eta_{ij}(\vec{E}) = \eta_{ij} + \sum_k r_{ijk} E_k + \sum_{lm} s_{ijklm} E_l E_m$$

\uparrow Pockels \uparrow Kerr.

Examples KDP (symmetry class $\bar{4}2m$) has

$$r = \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \\ r_{41} & 0 & 0 \\ 0 & r_{41} & 0 \\ 0 & 0 & r_{63} \end{bmatrix}$$

Then index ellipsoid becomes

$$\frac{x^2}{n_o^2} + \frac{y^2}{n_o^2} + \frac{z^2}{n_e^2} + 2r_{41}yzE_x + 2r_{41}xzE_y + 2r_{63}xyE_z = 1$$

$$(n_o = 1.51, n_e = 1.47)$$

How can we use this to make a modulator?

want to change effective n's. x & y are symmetric, so just consider applying E_x or E_z .

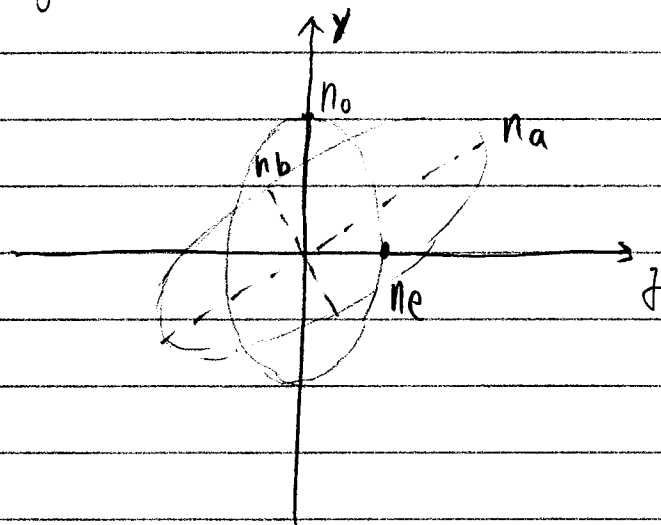
Apply E_x :

Only $y-z$ plane is affected. n_x remains n_o
So take $y-z$ = plane of polarization
→ light propagates along x

In $y-z$ plane, polarization ellipse is

$$\frac{y^2}{n_o^2} + \frac{z^2}{n_e^2} + 2r_{41} \frac{y}{z} E_x = 1$$

The electric field will rotate the ellipse from its original orientation.



Need to find new principal axes; easiest way is to diagonalize quadratic form

$$\begin{bmatrix} y & z \end{bmatrix} \begin{bmatrix} \frac{1}{n_o^2} & r_{41} E_x \\ r_{41} E_x & \frac{1}{n_e^2} \end{bmatrix} \begin{bmatrix} y \\ z \end{bmatrix}$$

Eigenvectors of matrix are principal axes, and eigenvalues are $\frac{1}{n_a^2}$, $\frac{1}{n_b^2}$.

Generally, eigenvalues of $\begin{bmatrix} A & C \\ C & B \end{bmatrix}$ is

$$\begin{vmatrix} A-\lambda & C \\ C & B-\lambda \end{vmatrix} = 0$$

$$\lambda^2 - (A+B)\lambda + AB - C^2 = 0$$

$$\text{are } \lambda_{\pm} = \frac{1}{2} \left[A+B \pm \sqrt{(A-B)^2 + 4C^2} \right]$$

$$\text{Here } A-B = \frac{1}{1.51^2} - \frac{1}{1.47^2} = -0.0242$$

$$\text{suppose } E_x = 10 \frac{\text{kV}}{\text{mm}} = 10^7 \frac{\text{V}}{\text{m}} \quad \text{pretty large}$$

$$\text{Then } C = 8.6 \frac{\text{pF}}{\text{V}} \cdot 10^7 \frac{\text{V}}{\text{m}} = 8.6 \times 10^{-5}$$

Then $4C^2 \ll (A-B)^2$ for any practical E_x

$$\lambda_{\pm} \approx A - \frac{C^2}{(B-A)}, \quad B + \frac{C^2}{B-A}$$

$$\begin{aligned} \text{SO } \frac{1}{n_a^2} &= \frac{1}{n_0^2} - \frac{C^2}{B-A} & \frac{1}{n_b^2} &= \frac{1}{n_0^2} + \frac{C^2}{B-A} \\ &= \frac{1}{n_0^2} - \Delta\left(\frac{1}{n^2}\right) & &= \frac{1}{n_0^2} + \Delta\left(\frac{1}{n^2}\right) \end{aligned}$$

$$\text{So } \Delta\left(\frac{1}{n^2}\right) = 2 \frac{\Delta n}{n^3} = \frac{c^2}{B-A}$$

$$\Delta n = \frac{1}{2} \cdot n^3 \frac{(r_{41} E_x)^2}{B-A}$$

$$\approx \frac{1}{2} (1.5)^3 \frac{(8.6 \times 10^{-5})^2}{0.0242}$$

$$\approx 5 \times 10^{-7}$$

Shift is not linear, and very small, Not good for modulation.

Now try applying E_z , light propagating along z .

Index ellipse:

$$\frac{x^2}{n_o^2} + \frac{y^2}{n_o^2} + 2xy/r_{63} E_z = 1$$

Matrix is

$$\begin{bmatrix} \frac{1}{n_o^2} & r_{63} E_z \\ r_{63} E_z & \frac{1}{n_o^2} \end{bmatrix}$$

$A = B$ so

$$n_{\pm} = \frac{1}{n_o^2} \pm r_{63} E_z \rightarrow \text{linear effect}$$

$$\text{So } \Delta n = \frac{1}{2} n^3 \Delta n = \pm \frac{1}{2} n_o^3 r_{63} E_z$$

$$= \pm 1.8 \times 10^{-9} \text{ at } 10 \text{ kV/mm}$$

much better.

* This is just like quantum mechanics:
off-diagonal perturbation only has significant effect if original states are degenerate.

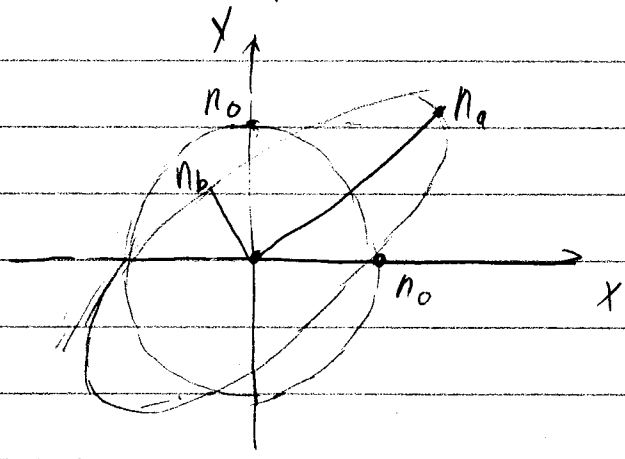
Also need eigenvectors.

easy to see new axes are:

$$x' = \frac{1}{\sqrt{2}} (x + y)$$

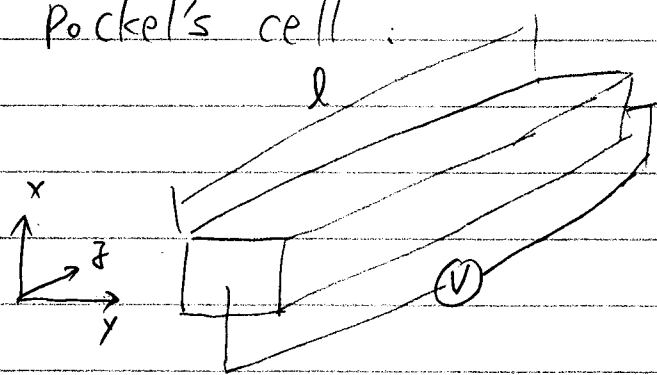
$$y' = \frac{1}{\sqrt{2}} (x - y)$$


Ellipse




$$\hat{x}': n_a = n_0 + \Delta n, \hat{y}': n_b = n_0 - \Delta n.$$

To make Pockel's cell:



 polarizer horizontal

Transparent electrodes

 polarizer, vertical

⑥

Incident light polarized along $\hat{x} = \frac{1}{\sqrt{2}}(\hat{x}' + \hat{y}')$

So after crystal, have

$$\begin{aligned}\vec{E}_{\text{out}} &\propto \frac{1}{\sqrt{2}} \left[\hat{x}' e^{+ik_0 l \Delta n_0} + \hat{y}' e^{-ik_0 l \Delta n_0} \right] \\ &= \frac{1}{\sqrt{2}} \left[\hat{x}' + \hat{y}' e^{-2ik_0 l \Delta n_0} \right]\end{aligned}$$

Phase shift $\phi = (2k_0 l) \left(\frac{1}{2} n_0^3 r_{63} E_z \right)$

$$\begin{aligned}\text{Here } E_z &= \frac{V}{l}, \text{ so } \phi = k_0 n_0^3 r_{63} V \\ &= \frac{2\pi}{\lambda_0} n_0^3 r_{63} V.\end{aligned}$$

Recall:

$$I_{\text{out}} = I_{\text{in}} \sin^2 \frac{\phi}{2}$$

So to get max output need $\phi = \pi$

$$\pi = \frac{2\pi}{\lambda} n_0^3 r_{63} V$$

$$\text{Write } V_{\pi} = \frac{\lambda}{2 n_0^3 r_{63}}$$

$$\text{at } \lambda = 633 \text{ nm}, \quad V_{\pi} = 8.7 \text{ kV}$$

Actually, KDP is a hard example.

Simpler if you can use a diagonal component of r .

For instance, r_{xxz} gives $\Delta n_x = \frac{1}{2} n_x^3 r_{13} E_z$.

If take light propagating along \hat{z}

ellipse is

$$\frac{x^2}{n_x^2} + \frac{y^2}{n_y^2} + r_{13} E_z x^2 = 1$$

axes don't change

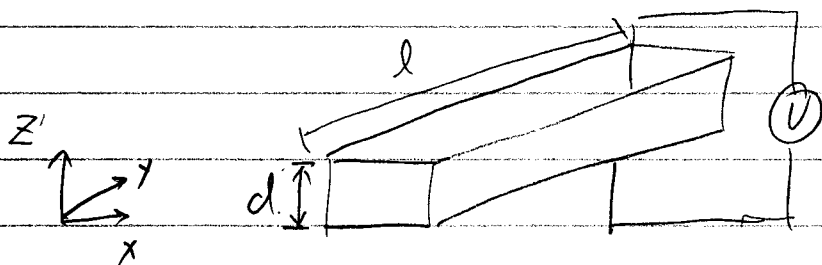
$$\hat{x} \rightarrow \hat{x}' \quad , \quad \hat{y} \rightarrow \hat{y}'$$

Also can take light propagating along \hat{y} , then

ellipse is

$$\frac{x^2}{n_x^2} + \frac{z^2}{n_z^2} + r_{13} E_z x^2 = 1$$

Then \vec{z} is transverse to light



Don't need transparent electrodes
and can get better modulation using long skinny
crystal

$$\phi = k l \frac{1}{2} n_x^3 Z_2 r_{31}$$

$$= \frac{1}{2} k l n_x^3 r_{31} \frac{V}{d}$$

$$= \pi \frac{l}{\pi_0 d} n_x^3 r_{31} V$$

$$V_{\pi} = \frac{\pi_0 d}{l} \cdot \frac{1}{n_x^3 r_{31}}$$

much smaller if $d \ll l$