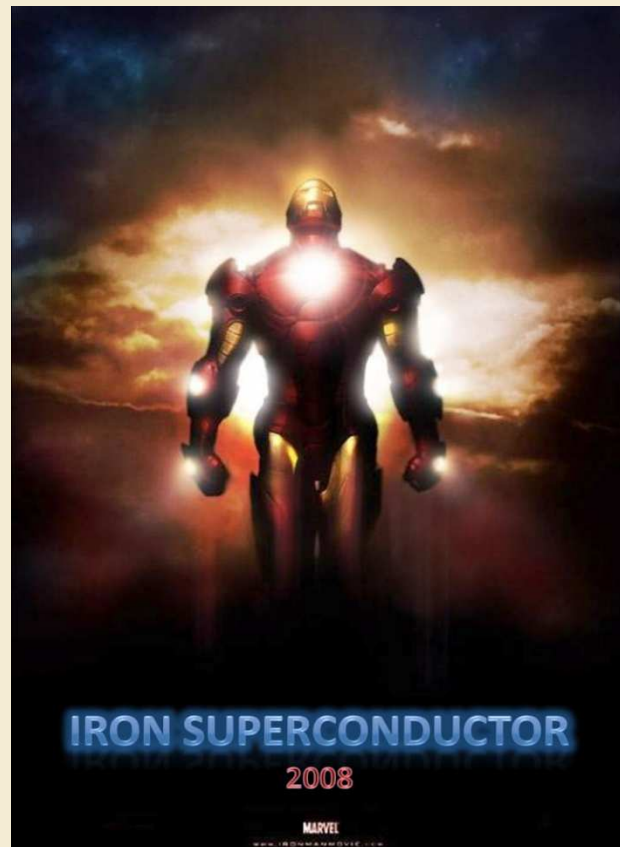
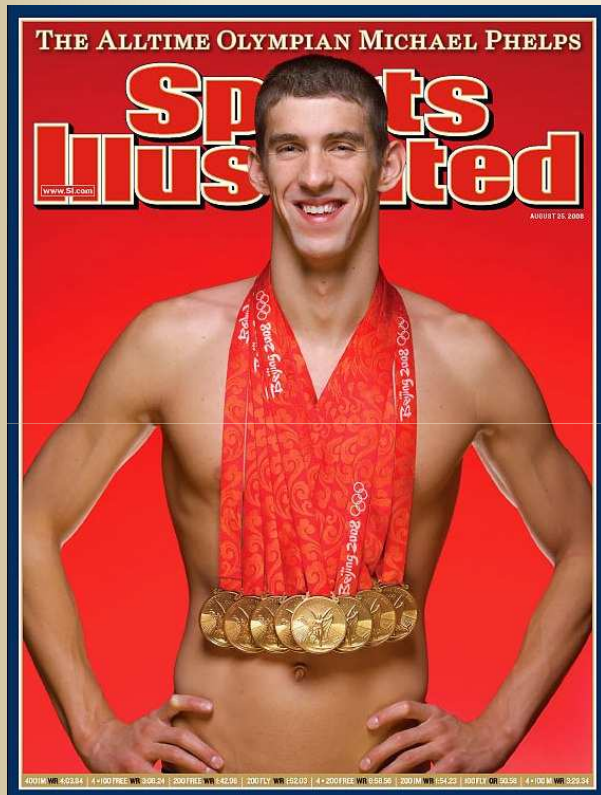


Feb. 04, University of Virginia

2008





of
IRON Superconductors

Outline

- **Introduction**

- **Point Contact (PC) spectroscopy**

 - Spin Transfer Torque (**STT**) Effects

- **Point Contact Andreev Reflection (**PCAR**) spectroscopy**

 - Spin polarization

 - Superconducting **gap**

 - Pseudogaps ?

- **Summary**

Discovery of Superconductivity

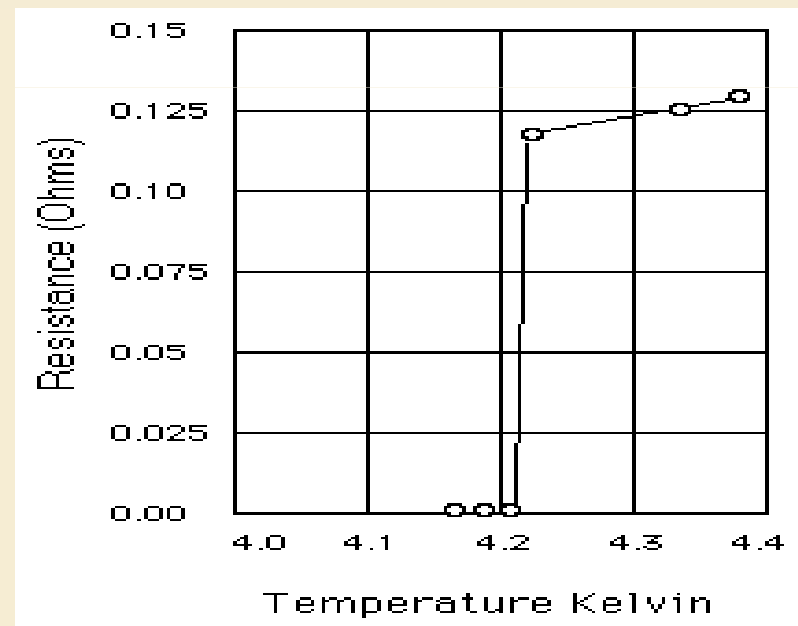
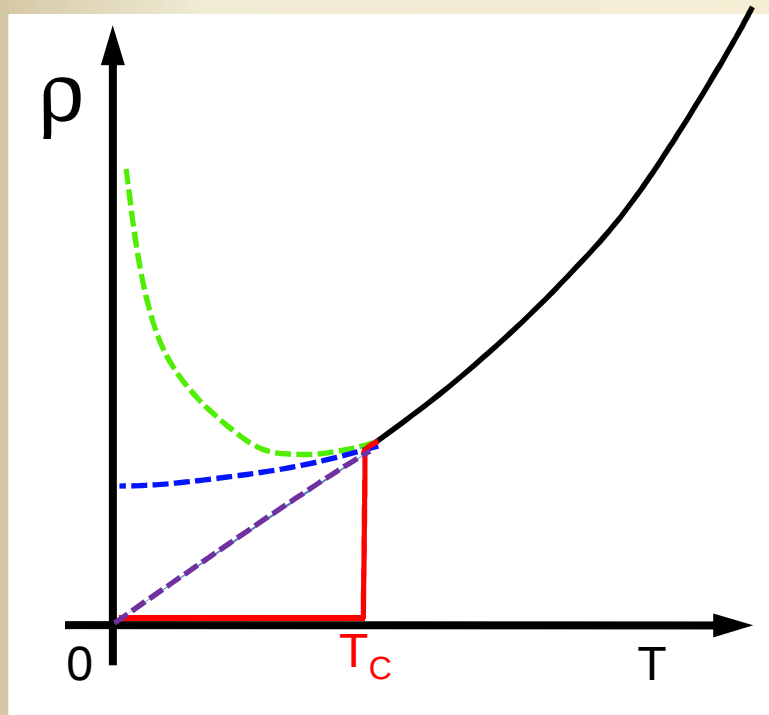


Heike Kamerlingh Onnes

Liquefied Helium (1908), b.p. = 4.17 K

Discovered superconductivity (1911) in Hg $T_c = 4.21$ K

Nobel prize (1913)



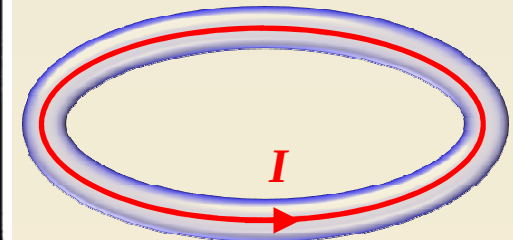
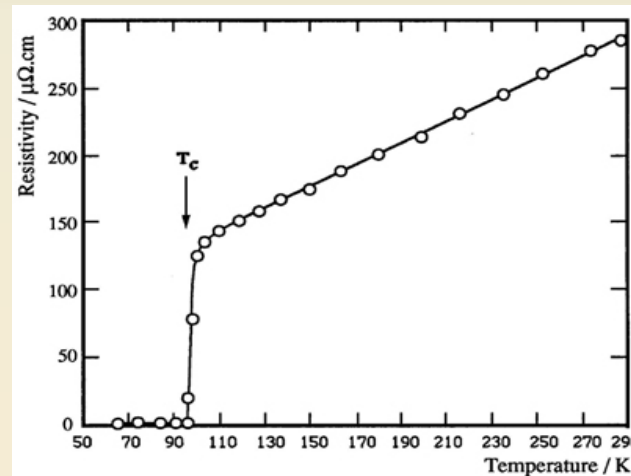
The saga of superconductivity began.

Perfect Conductivity and Perfect Diamagnetism

Zero resistance

Perfect conductivity

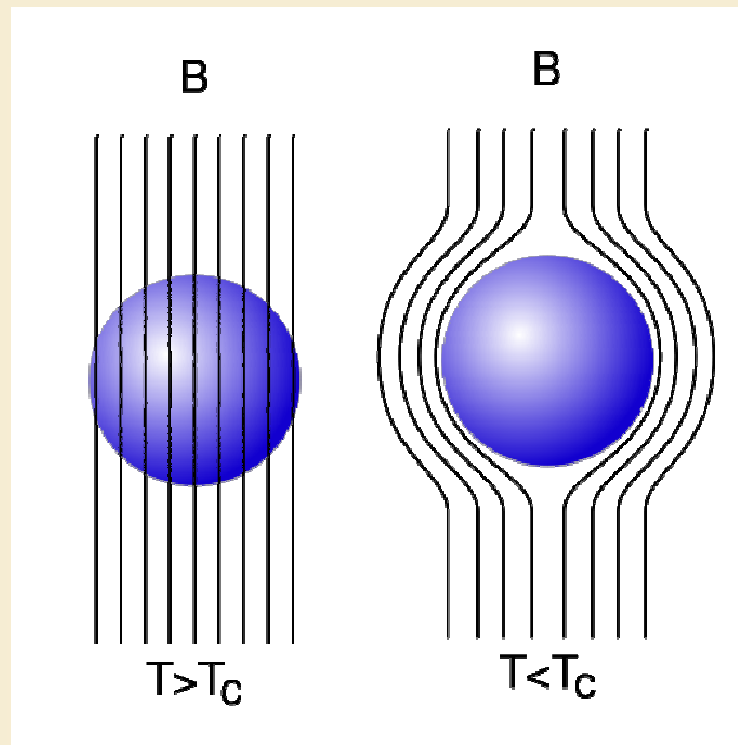
Persistent current ***I***
over 10^5 years



Meissner effect

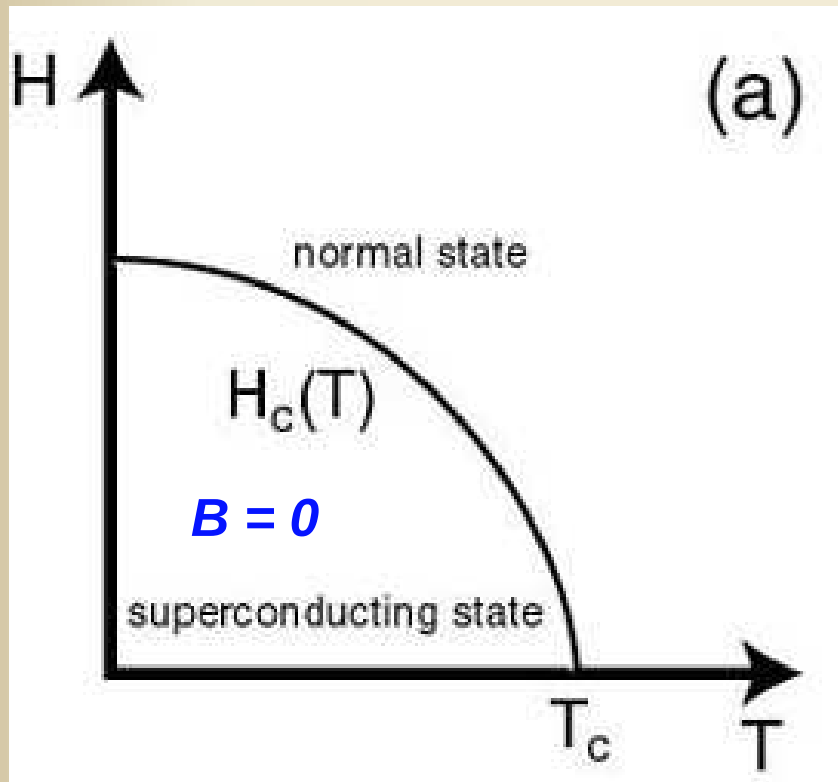
Perfect diamagnetism

No field inside SC



Type I and Type II Superconductors

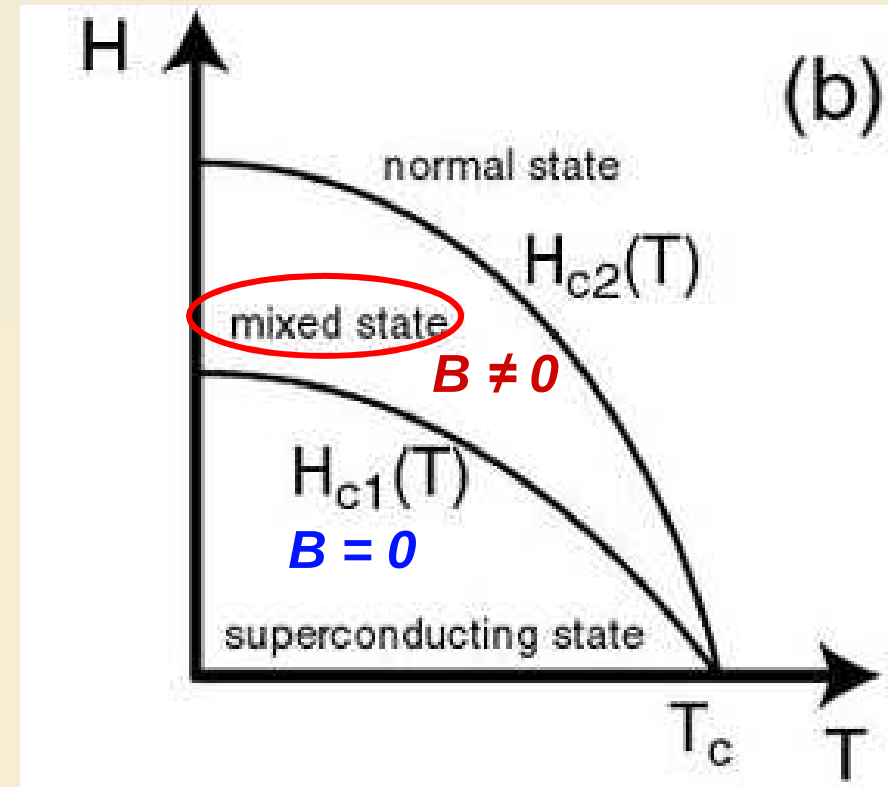
Type I (Pb, Al, etc)



Meissner ($H < H_c$)

$H_c(0) \approx 0.1$ Tesla

Type II (NbTi, Nb₃Sn, etc)

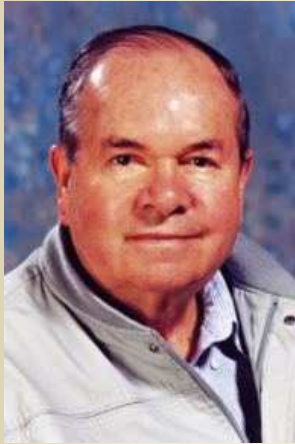


Meissner ($H < H_{c1}$)

Mixed state ($H_{c1} < H < H_{c2}$)

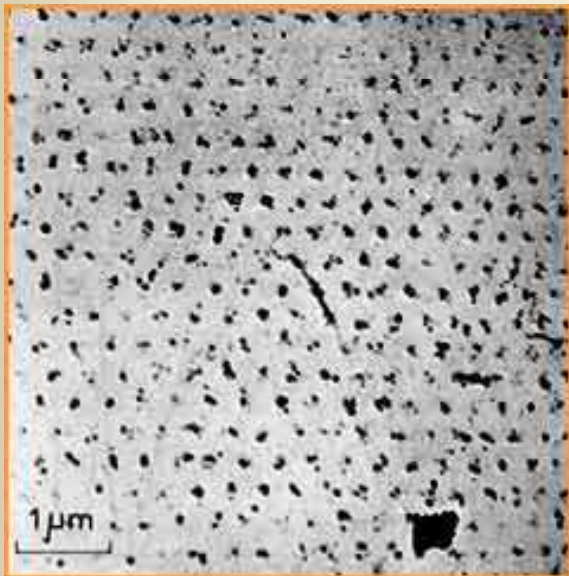
$H_{c2}(0)$ very large

Vortices in Mixed State

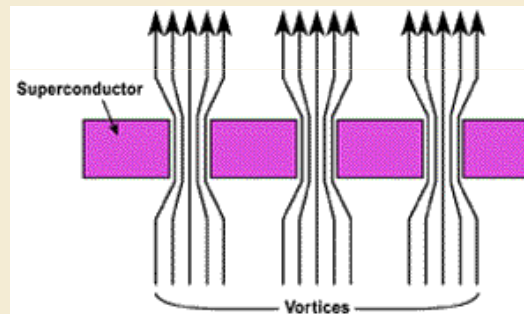


Abrikosov Vortex Lattices in Superconductors (1957)

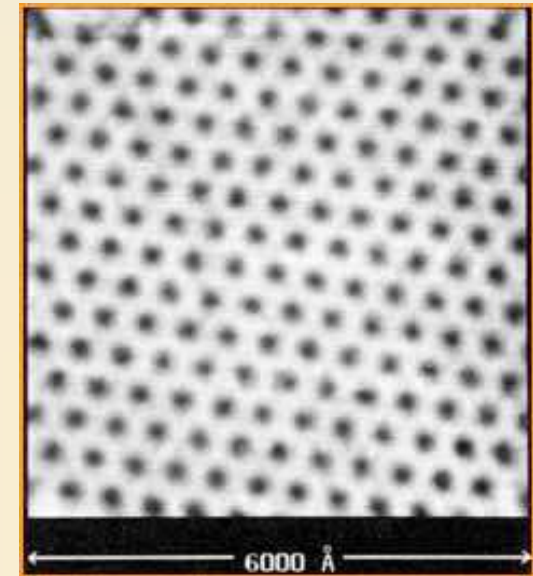
Nobel Prize (2003)



Essmann (1967) (Bitter, Pb)



$$\Phi_0 = hc/2e \approx 20 \text{ G}\cdot\mu\text{m}^2$$



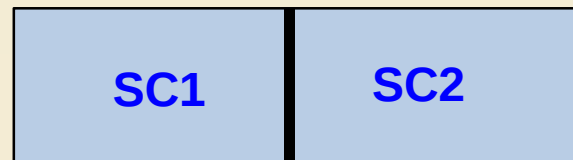
Hess (1989) (STM, NbSe₂)

Josephson Effect



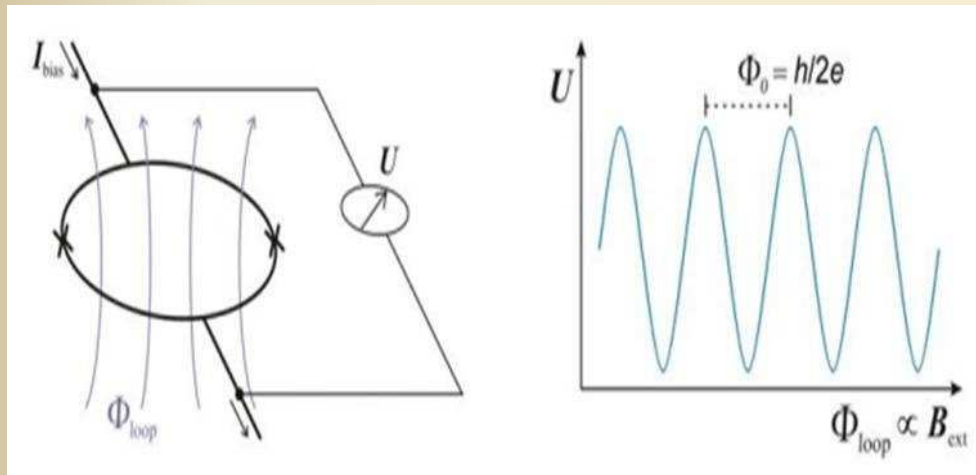
Brian D. Josephson

1962 Theoretical prediction of Josephson effect
1973 Nobel prize



$$I(t) = I_C \sin(\phi(t))$$

ϕ : phase difference



$$\Phi_0 = \frac{h}{2e} = 2.067833636 \times 10^{-15} \text{ Wb}$$

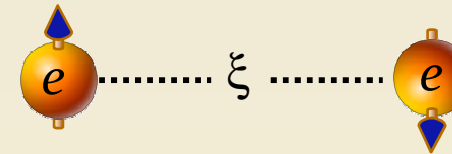
Superconducting Quantum Interference Device (SQUID)

Most sensitive detector

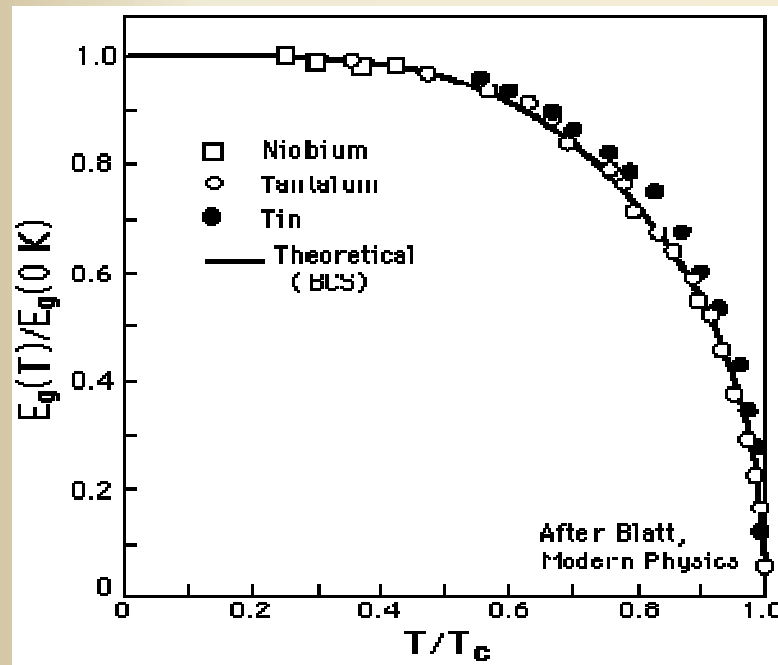
Bardeen-Cooper-Schrieffer (BCS) Theory (1957)



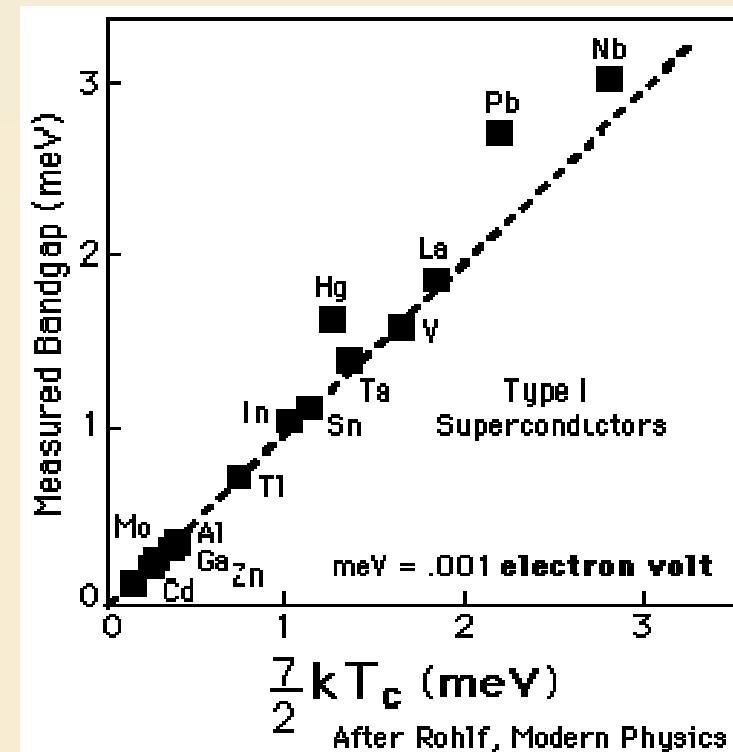
Bardeen-Cooper-Schrieffer
Nobel Prize, 1972



Cooper pairs

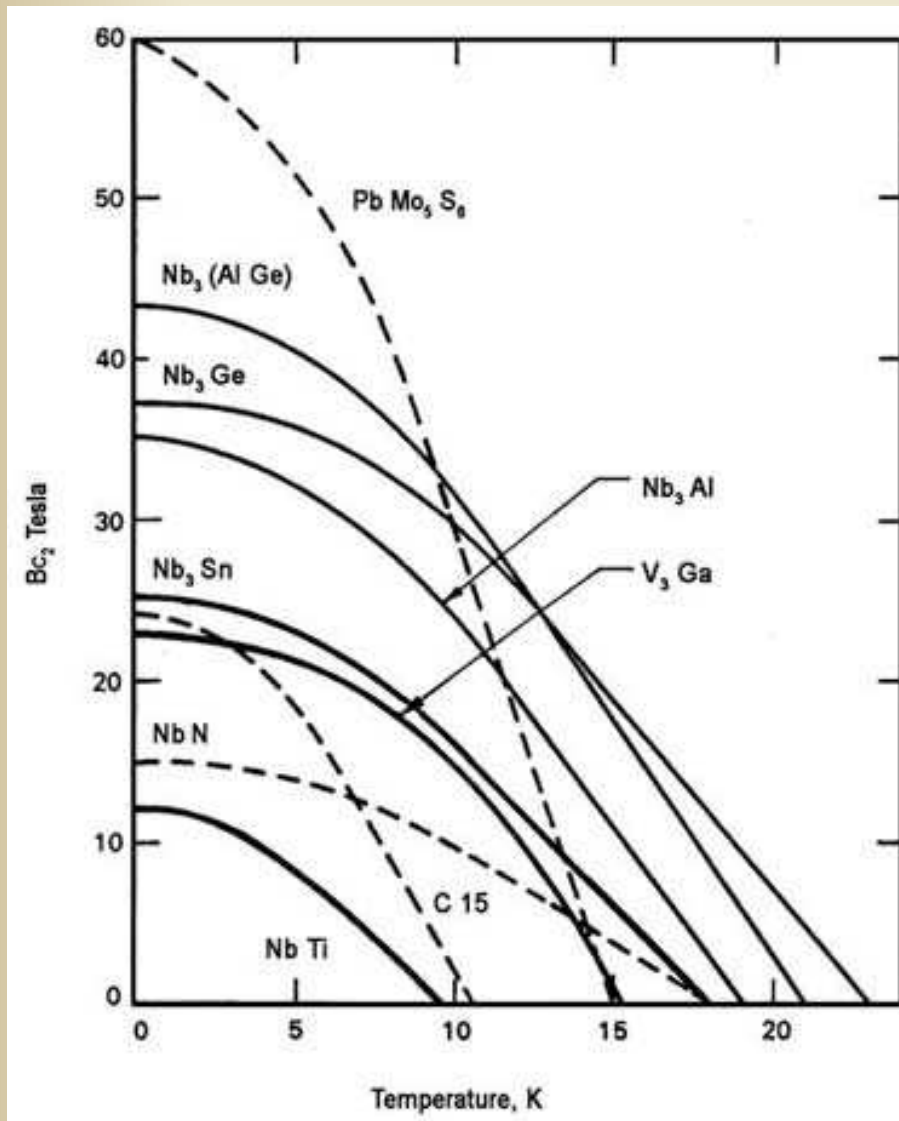


Superconducting Gap
 $E_g = 2 \Delta$



$$2\Delta(0)/k_B T_C = 3.53$$

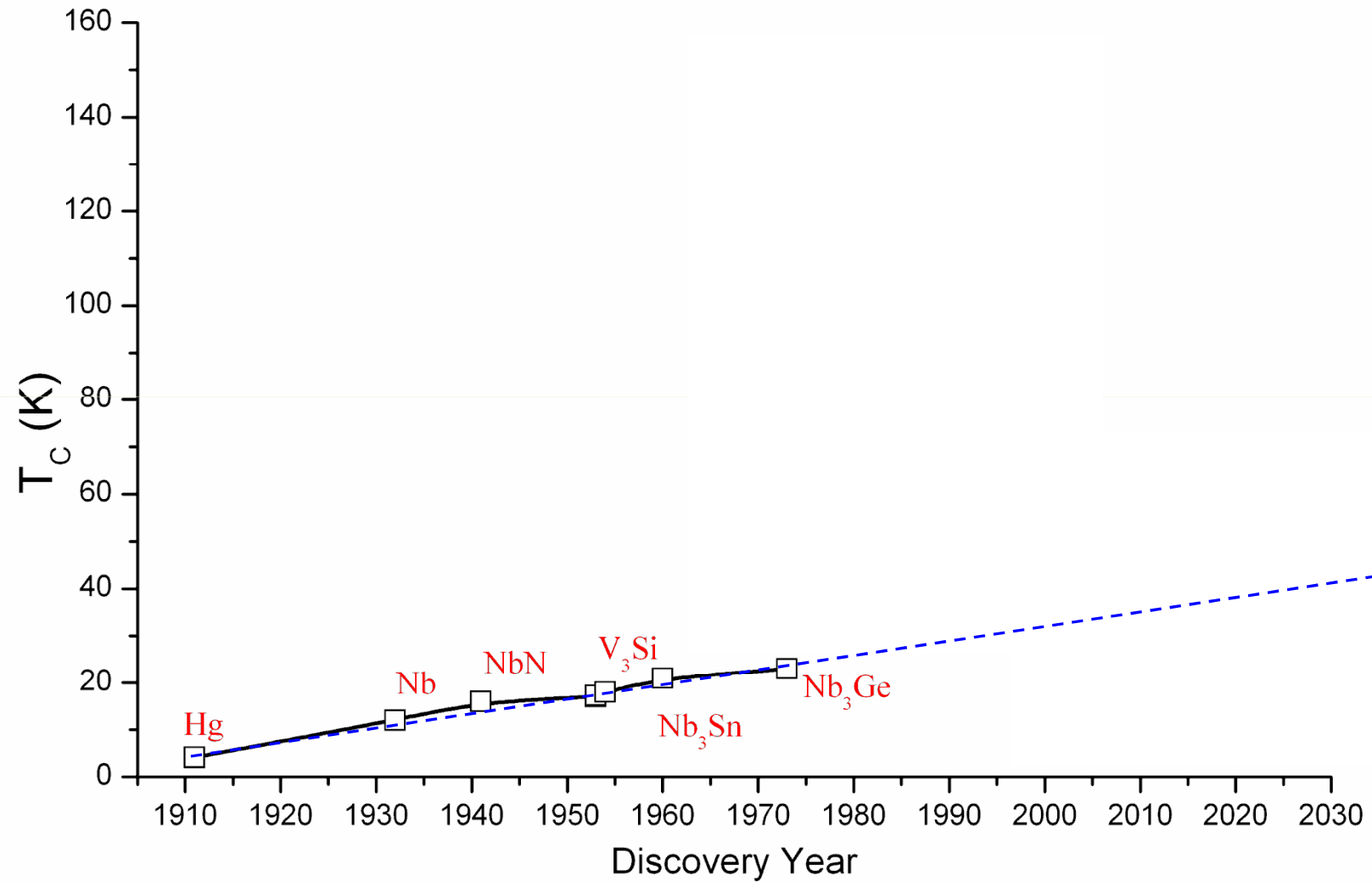
Very large H_{c2} for Type II Superconductors



High H_{c2} is necessary
Superconducting magnets
Magnetic levitation

High T_C is always beneficial

T_C vs. Discovery Year



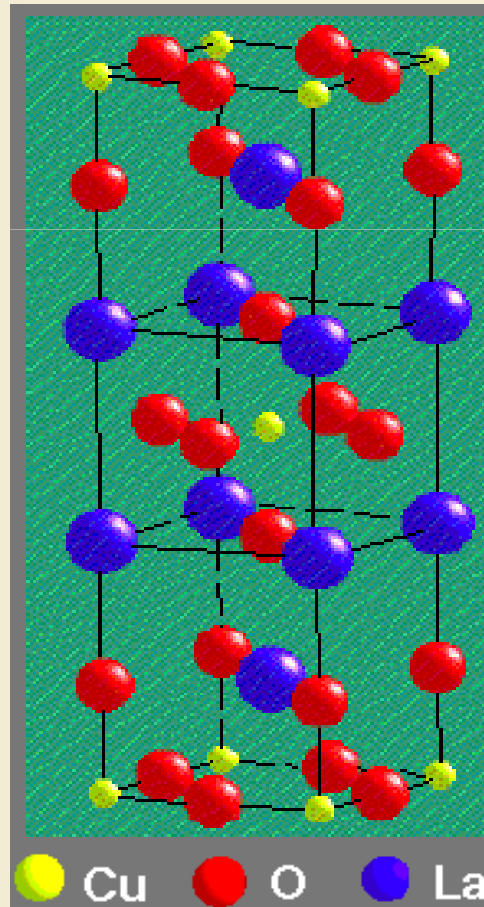
Discovery of Cuprates (1986)



Alex Müller and Georg Bednorz

Discovered $(\text{La-Ba})_2\text{CuO}_4$ in 1986, $T_c = 30 \text{ K}$

Nobel prize, 1987



CuO_2 Plane

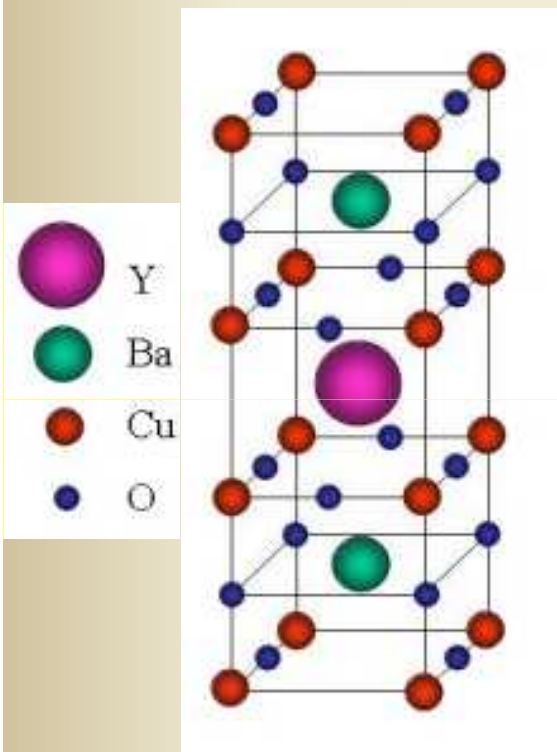
La_2CuO_4

$(\text{La-Sr})_2\text{CuO}_4$

$(\text{La-Ba})_2\text{CuO}_4$

$T_c \sim 30 - 40 \text{ K}$

Cuprates



Cu chains

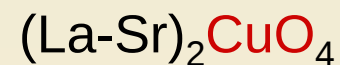
Cu planes

Cu planes

Cu chains

Cuprates

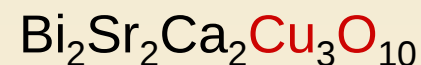
T_c



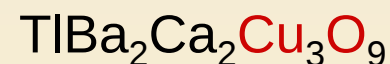
35 K



90 K



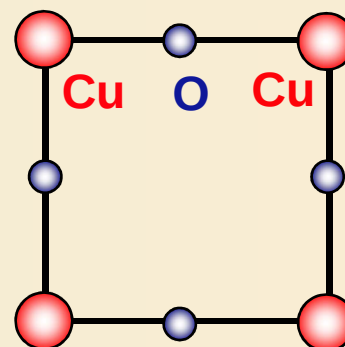
105 K



125 K



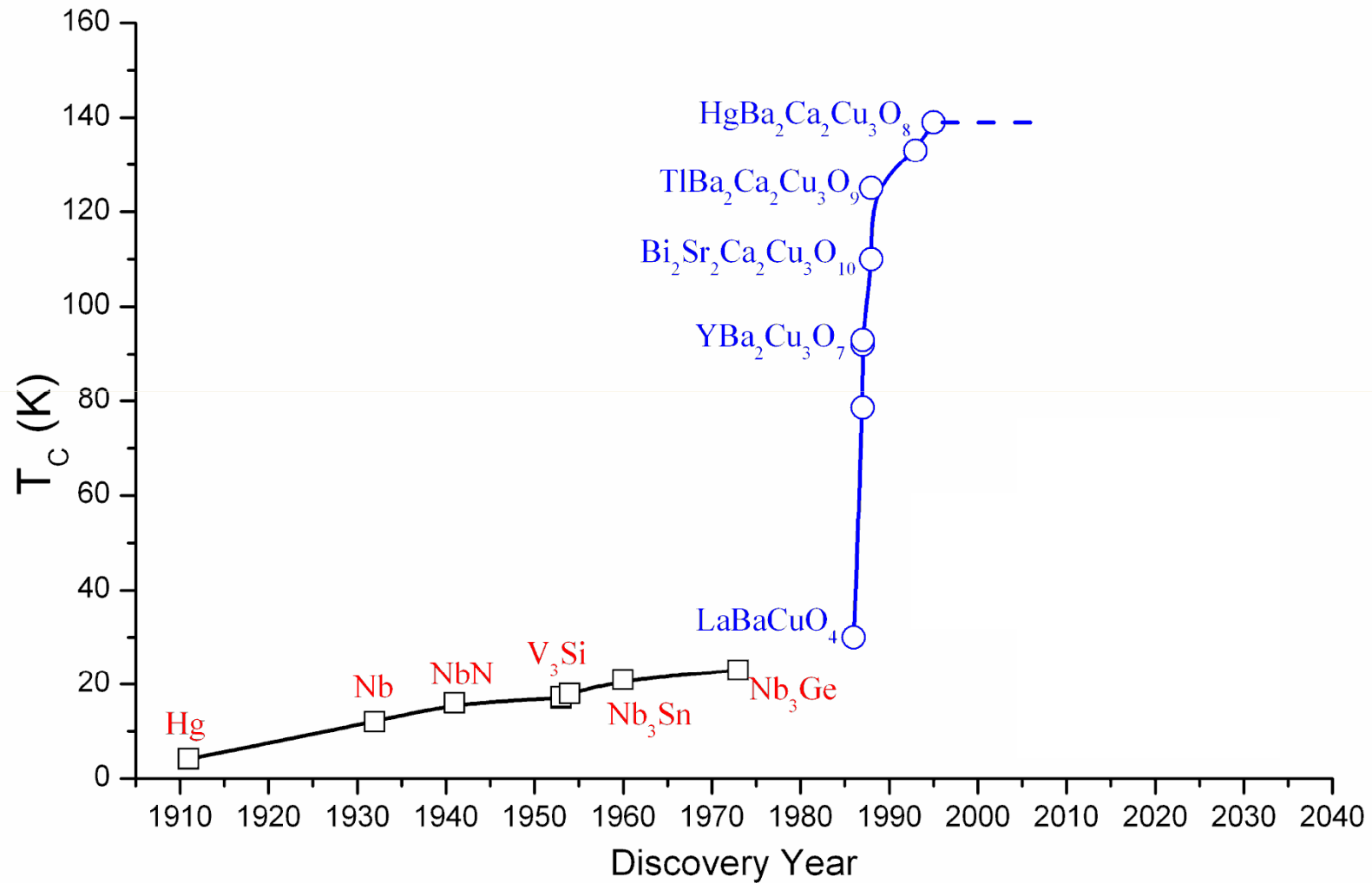
140 K



CuO_2 plane

Responsible for high T_c and a host of unusual properties

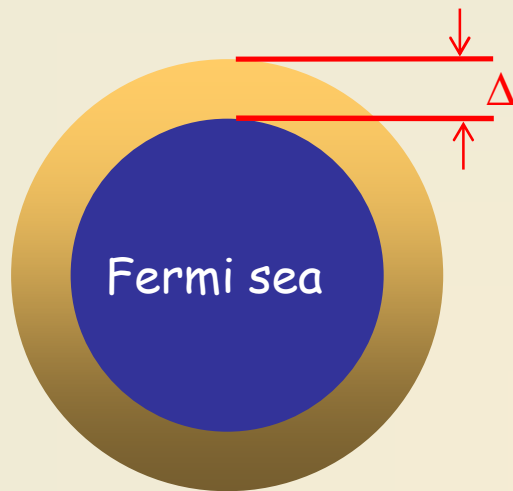
T_C vs. Discovery Year



Mind the Gap



s-wave and *p*-wave Superconductor



s-wave (**isotropic**)

Conventional SC (Nb, Pb, Al,...)

$T_c < 25$ K



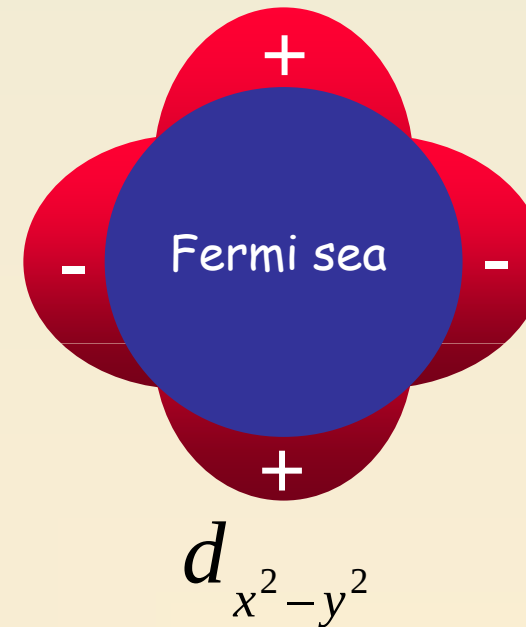
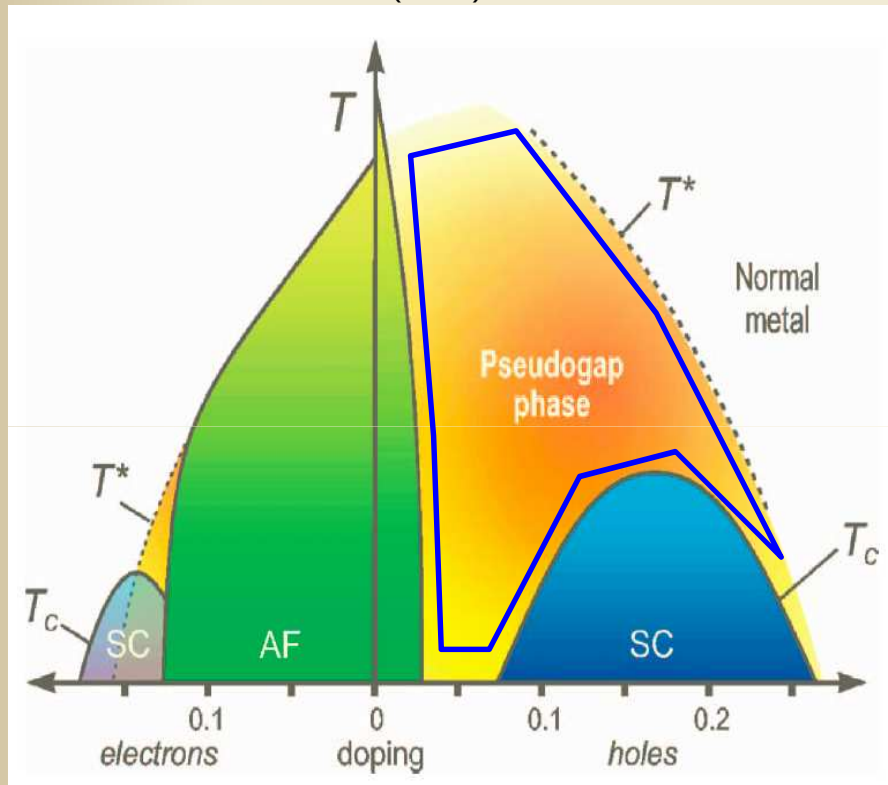
p-wave

Superfluid ^3He , $\text{Sr}_2\text{RuO}_4(?)$

$T_c \sim 1$ mK - 1 K

d-wave Superconductor

Fisher *et al.*, RMP **79**, 353 (2007)



d-wave symmetry

Parent compound: **Mott insulator**

Cuprates: **CuO₂ plane**

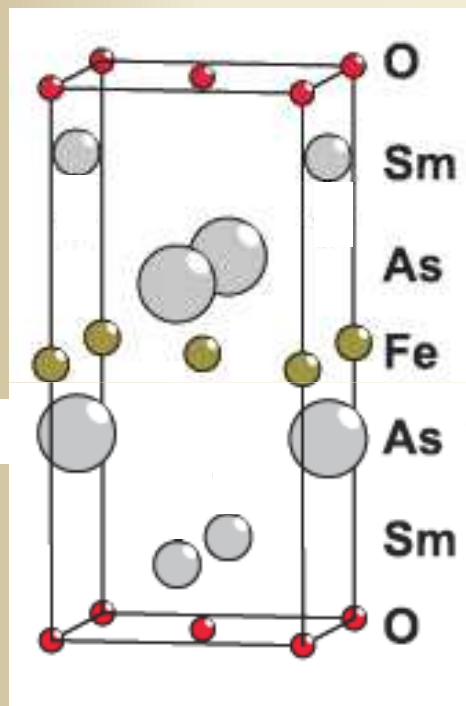
$T_c < 160$ K

New Oxypnictides Fe Superconductors (2008)

“Single layer”

SmFeAsO

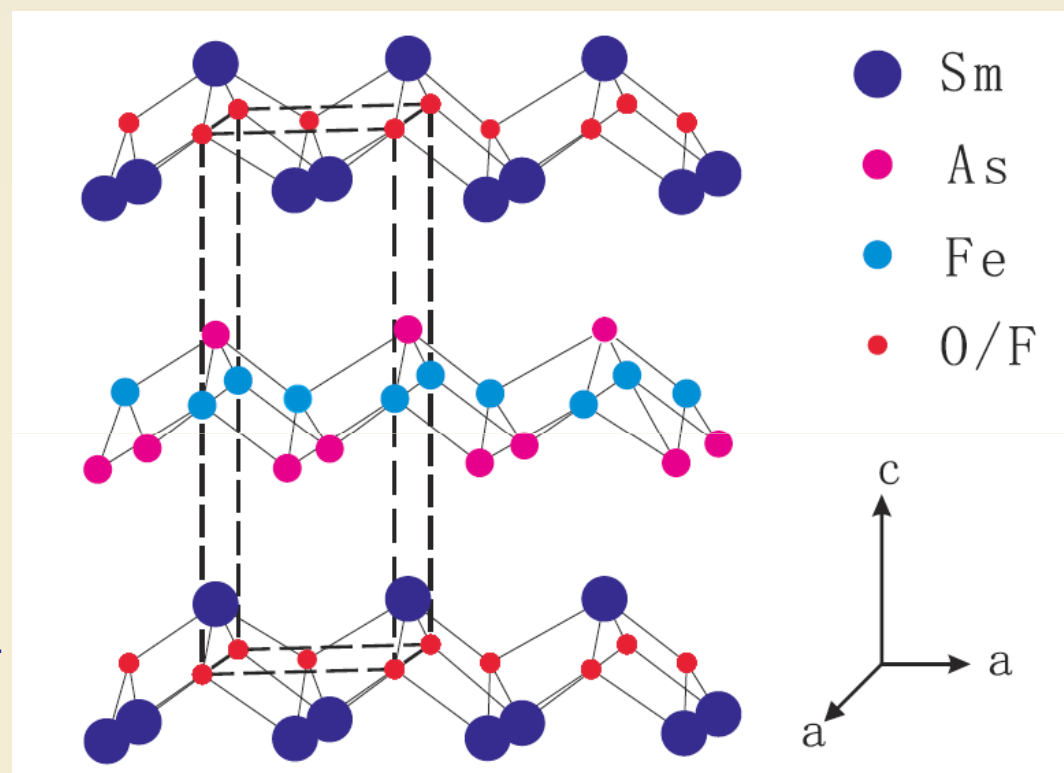
Pnictogen (P, As, Sb)



(SmO)⁺

(FeAs)⁻

(SmO)⁺



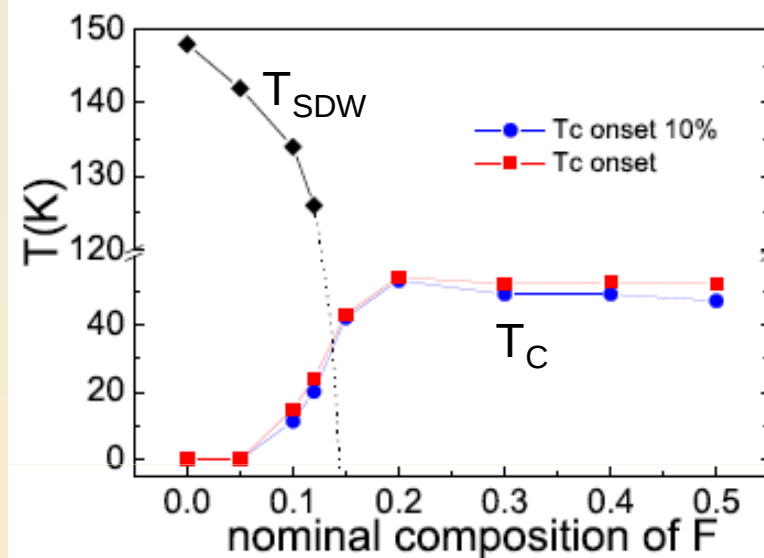
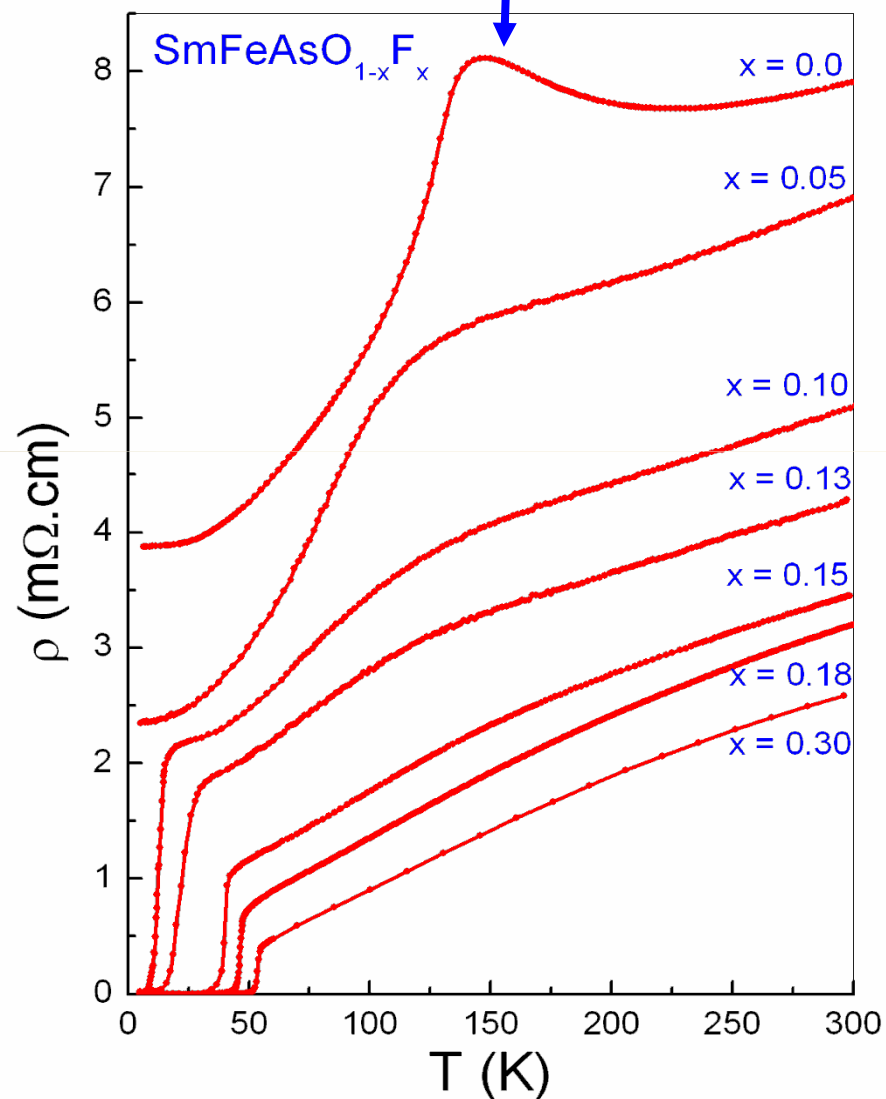
No CuO₂ plane

Contains Fe

Contains As

Discovery of Fe-Superconductors (2008)

(SDW)



New Fe-SC:

$\text{SmFeAsO}_{1-x}\text{F}_x$ (1111) $T_{\text{C}} \sim 55\text{K}$

Parent compound:

SmFeAsO metallic

Spin density wave (SDW) $\sim 160\text{K}$

Fe-Superconductors

New Fe-SC:

$\text{SmFeAsO}_{1-x}\text{F}_x$ (1111), $T_c \sim 55\text{K}$

Parent compound:

SmFeAsO metallic

Spin density wave (**SDW**) at 160K

$$\Delta = ?$$

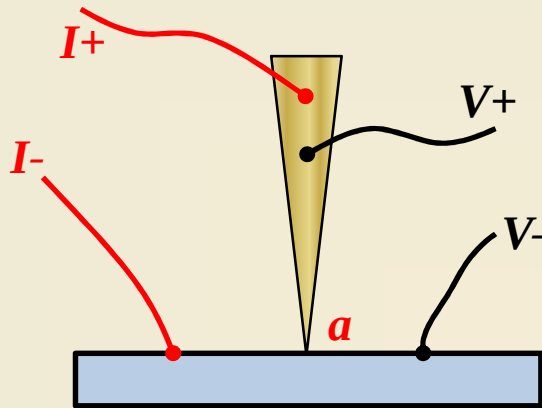
What **value**?

What **structure**?

What **temperature dependence**?

Are there **pseudogaps**?

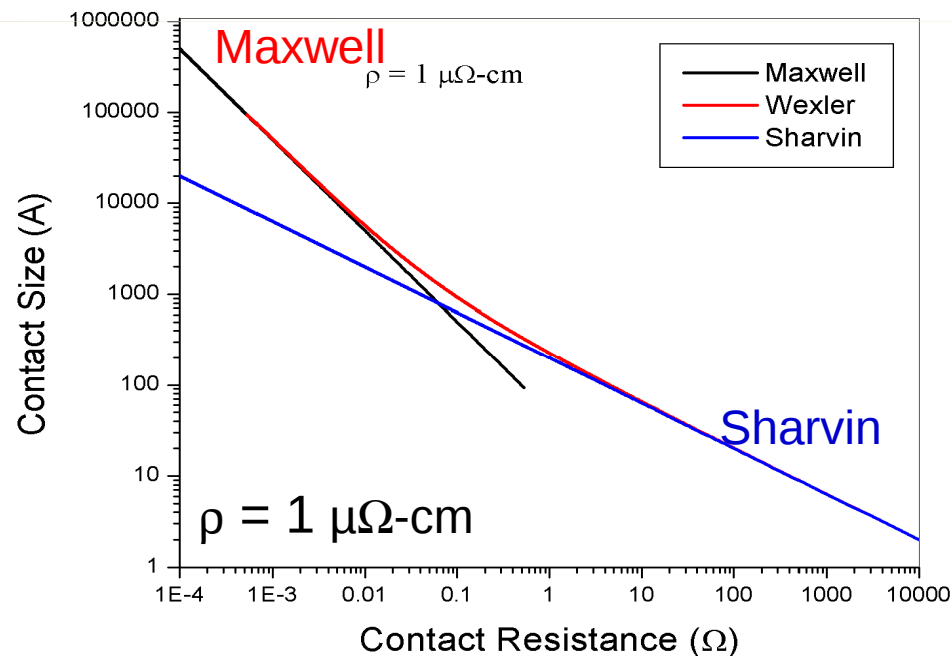
Point Contact



$$R_{Maxwell} = \frac{\rho}{2a} \quad a \gg l, \text{ diffusive}$$

$$R_{Wexler} = \frac{4\rho l}{3\pi a^2} \left(1 + \frac{3\pi}{8} \frac{a}{l}\right)$$

$$R_{Sharvin} = \frac{4\rho l}{3\pi a^2} \quad a \ll l, \text{ ballistic}$$



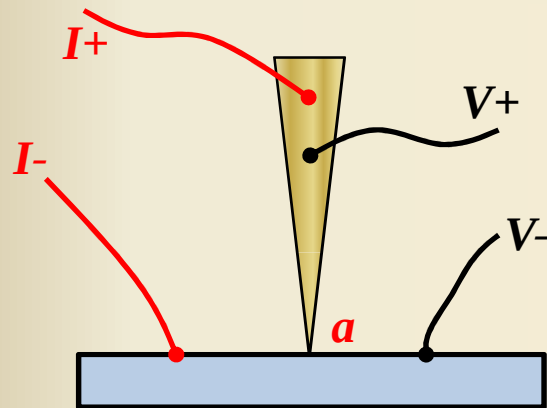
ρ : resistivity

a : contact size

l : mean free path

Different regime with different physics

Point Contact Spectroscopy



□ Phonon spectra

Phonon

□ Spin transfer torque effect

Contact size: $a \sim \text{nm}$

Current: $I = 1 \text{ mA}$

Current density: $j = 10^{10} \text{ A/cm}^2$

□ Point Contact Andreev Reflection

Spin polarization

Superconducting gap

□ Thermal transport

ballistic, diffusive

□ ...

Spin Transfer Torque Effect

Contact size: $a \sim \text{nm}$

Current: $I = 1 \text{ mA}$

Current density: $j = 10^{10} \text{ A/cm}^2$

T. Y. Chen, Y. Ji and C. L. Chien, *Appl. Phys. Lett.* **84**, 380 (2004).

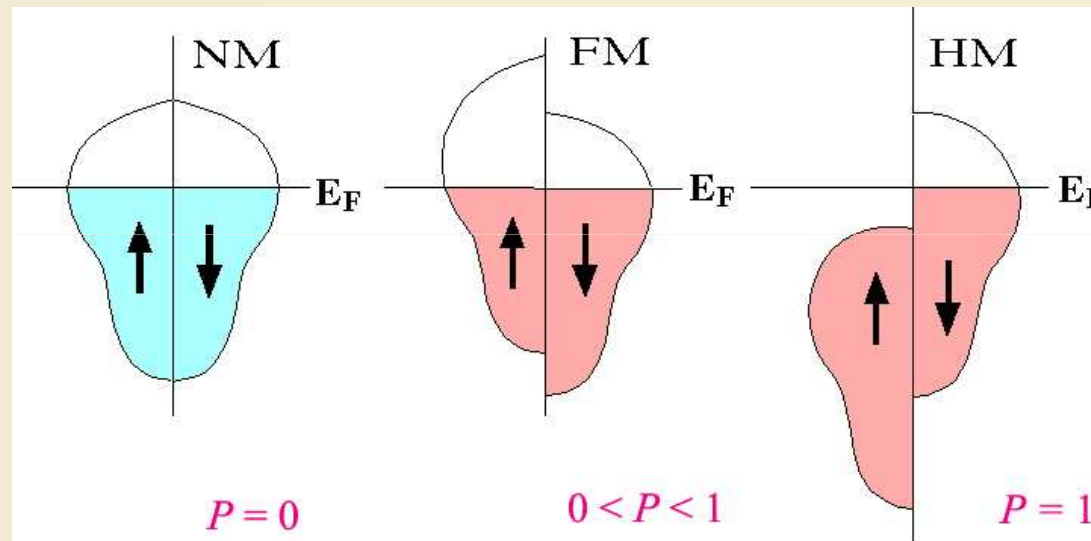
T. Y. Chen, Y. Ji, C. L. Chien and M. D. Stiles, *Phys. Rev. Lett.* **93**, 026601 (2004).

T. Y. Chen, S. X. Huang, C. L. Chien and M. D. Stiles, *Phys. Rev. Lett.* **96**, 207203 (2006).

Spin Polarization (P)

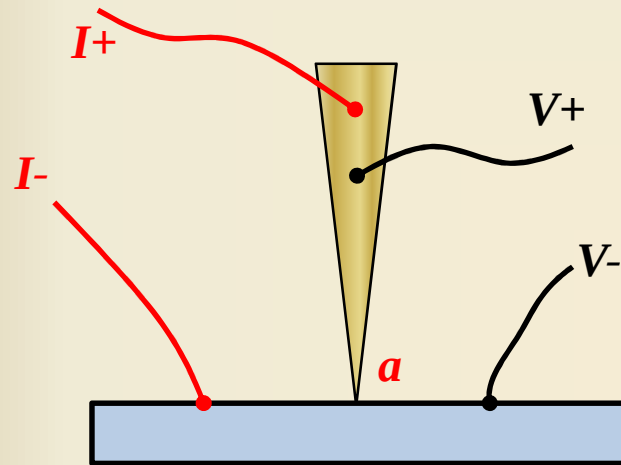
$$P_0 = \frac{N_{\uparrow}(E_F) - N_{\downarrow}(E_F)}{N_{\uparrow}(E_F) + N_{\downarrow}(E_F)}$$

$N(E_F)$: density of states at Fermi energy



P : crucial for spintronics

Point Contact Andreev Reflection (PCAR)



Tip:

Conducting, semiconducting,
superconducting

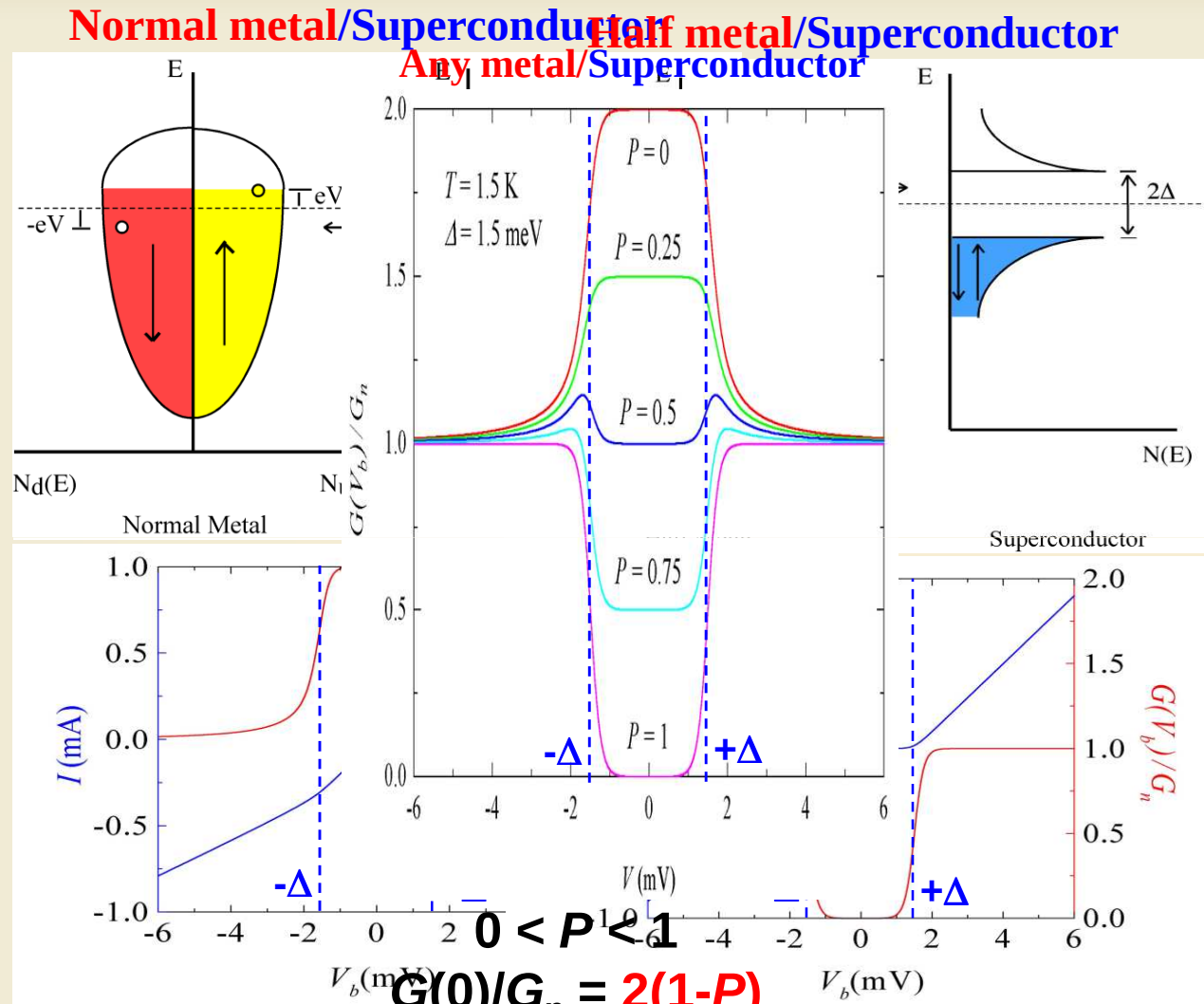
Sample:

Conducting, semiconducting,
superconducting

Normal metal/superconductor interface: Andreev Reflection

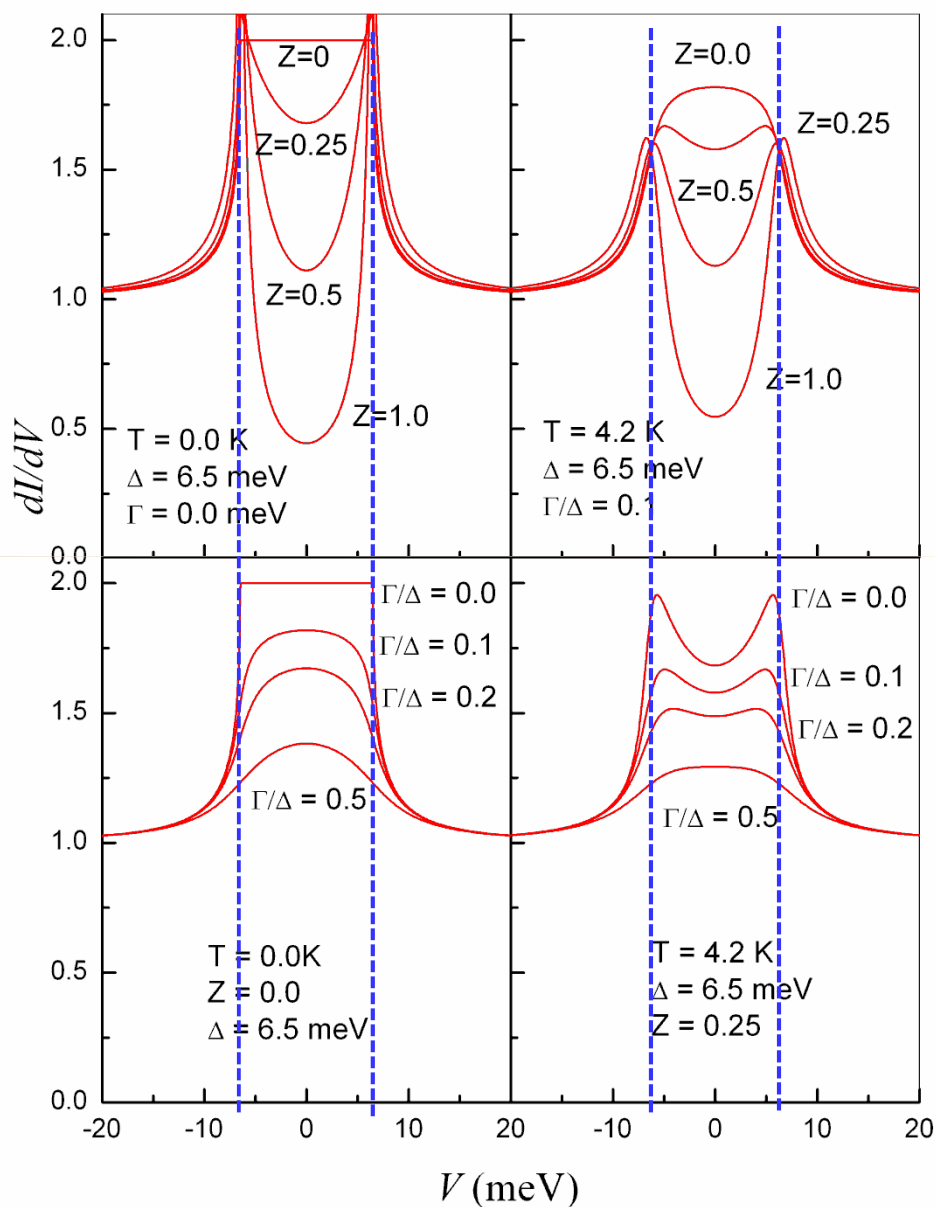
Measures **spin polarization (P)** and **superconducting gap (Δ)**

Determination of Spin Polarization by PCAR



$G(0)/G_n = 2(1-P)$
 $G(0)/G_n = 2$ $P = 0$ $G(0)/G_n = 0$ $P = 1$
 P of metal, and Δ of SC can be determined.
Ideal contact, clean limit

Interfacial Scattering (Z) and inelastic scattering (Γ)



□ Interfacial: Z

suppresses in-gap conductance

□ Inelastic: Γ

affects conductance at gap

□ Thermal: T

smears out conductance

Analyzing entire conductance

Blonder-Tinkham-Klapwijk (BTK) theory

Point Contact Andreev Reflection Spectroscopy

Using known SC (Nb, Pb) to determine *P* of metal

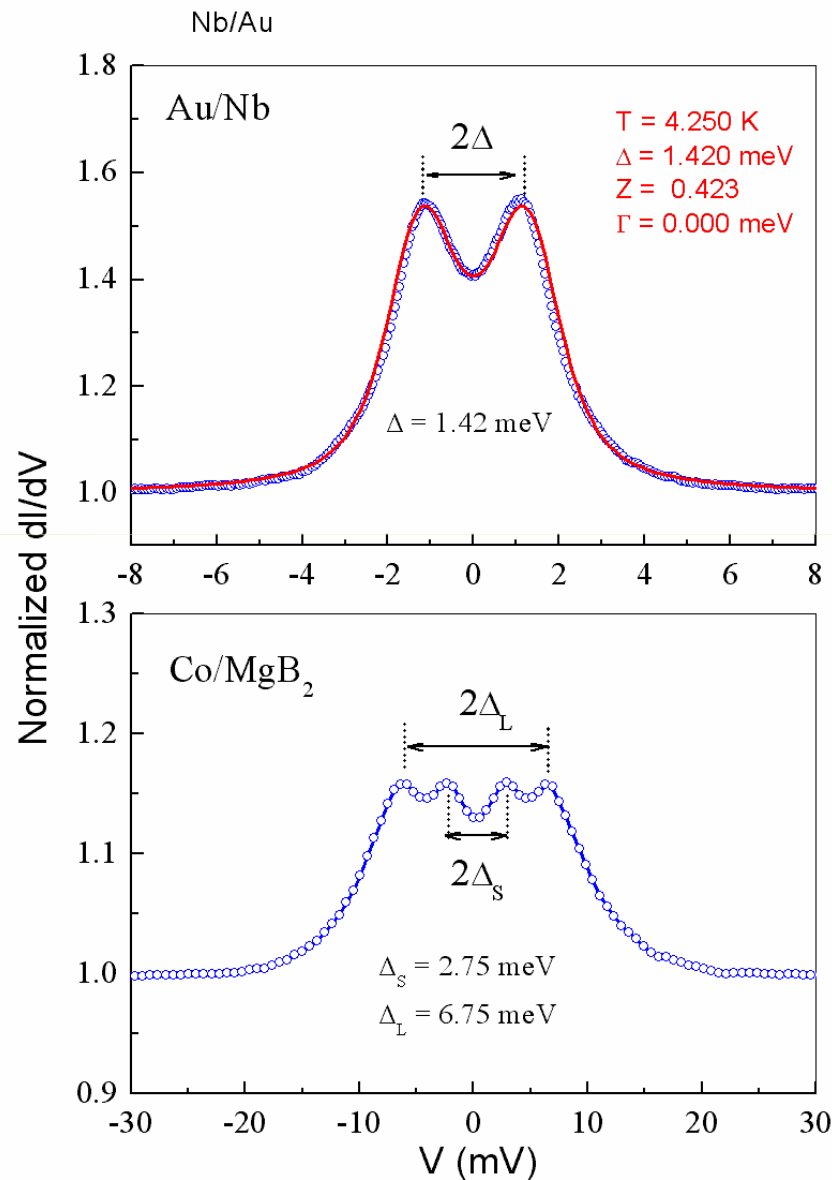
- S. X. Huang, T. Y. Chen, and C. L. Chien, *Appl. Phys. Lett.* **92**, 242509 (2008).
- T. Y. Chen, C. L. Chien and C. Petrovic, *Appl. Phys. Lett.* **91**, 142505 (2007).
- C. Leighton, M. Manno, A. Cady, J. W. Freeland, L. Wang, K. Umemoto, R. M. Wentzcovitch, T. Y. Chen, C. L. Chien, P. L. Kuhns, M. J. R. Hoch, A. P. Reyes, W. G. Moulton, E. D. Dahlberg, J. Checkelsky and J. Eckert, *J. Phys.: Condens. Matter* **19**, 315219(2007) (invited review).
- L. Wang, T. Y. Chen, C. L. Chien, and C. Leighton, *Appl. Phys. Lett.* **88**, 232509 (2006).
- L. Wang, T. Y. Chen, C. L. Chien, J. G. Checkelsky, J. C. Eckert, E. D. Dahlberg, K. Umemoto, R. M. Wentzcovitch, and C. Leighton, *Phys. Rev. B* **73**, 144402 (2006).
- L. Wang, K. Umemoto, R.M. Wentzcovitch, T. Y. Chen, C.L. Chien, J.G. Checkelsky, J.C. Eckert, E.D. Dahlberg and C. Leighton, *Phys. Rev. Lett.* **94**, 056602 (2005).
- L. Wang, T. Y. Chen and C. Leighton, *Phys. Rev. B* **69**, 094412 (2004).
- Lance Ritchie and Gang Xiao, Y. Ji, T. Y. Chen, C. L. Chien, Ming Zhang, Jinglan Chen, Zhuhong Liu, Guangheng Wu, and X. X. Zhang, *Phys. Rev. B* **68**, 104430 (2003).

Using known metal (Au, Co) to determine *Gap* of SC

Nb, MgB₂, SmFeAsO_{0.85}F_{0.15}

T. Y. Chen, Z. Tesanovic, R. H. Liu, X. H. Chen, and C. L. Chien, *Nature* **453**, 1224 (2008)

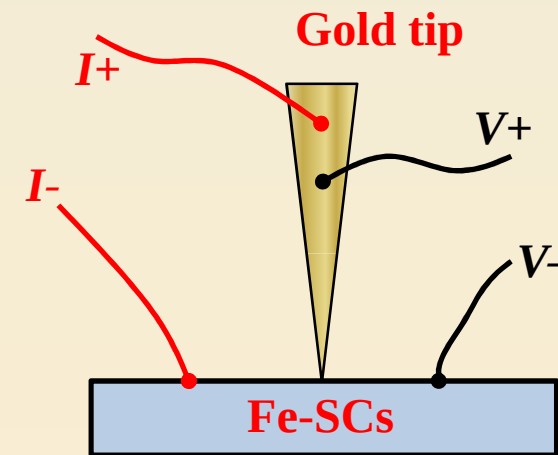
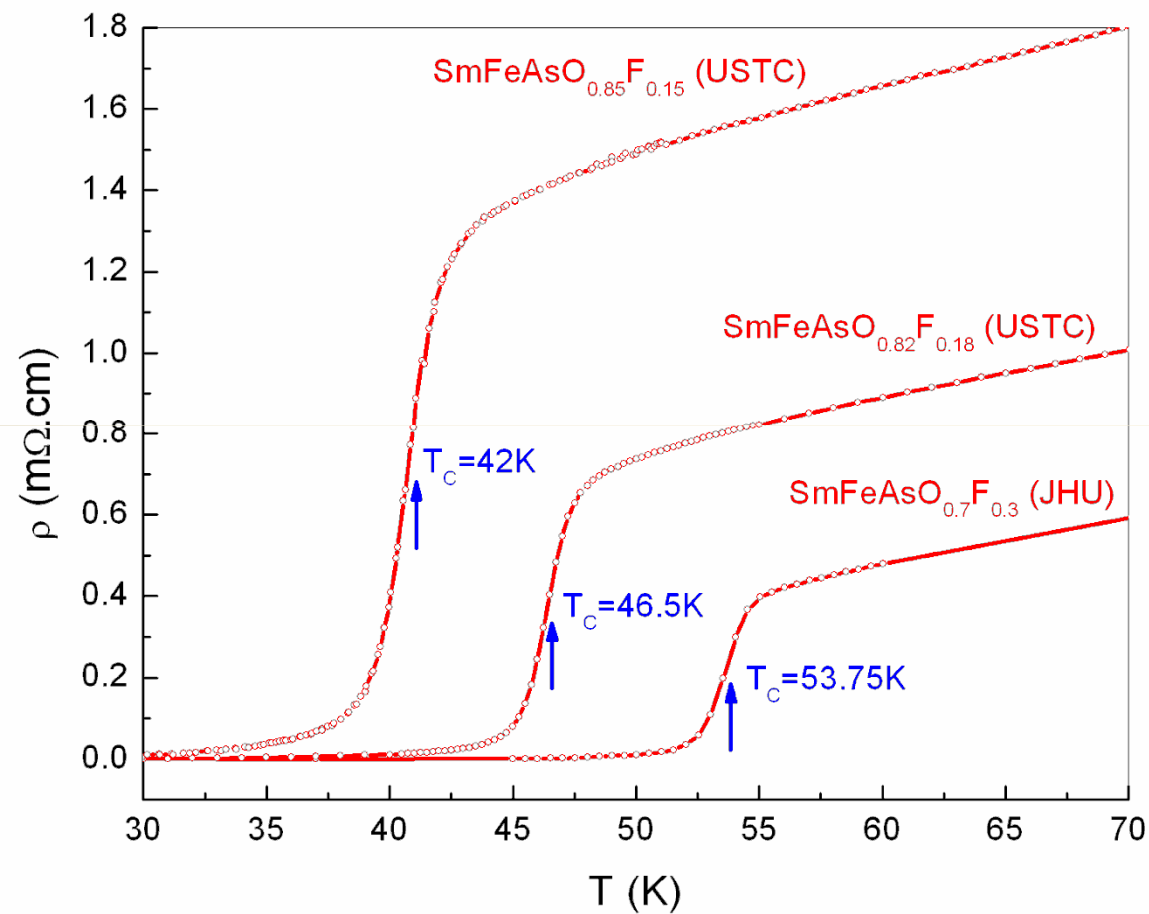
AR Spectroscopy Measurements of Gaps



One gap of Nb ($T_C = 9.27$ K)
 $\Delta_{\text{Nb}} = 1.42$ meV

Two gaps of MgB₂ ($T_C = 36.5$ K)
 $\Delta_L = 6.75$ meV
 $\Delta_S = 2.75$ meV

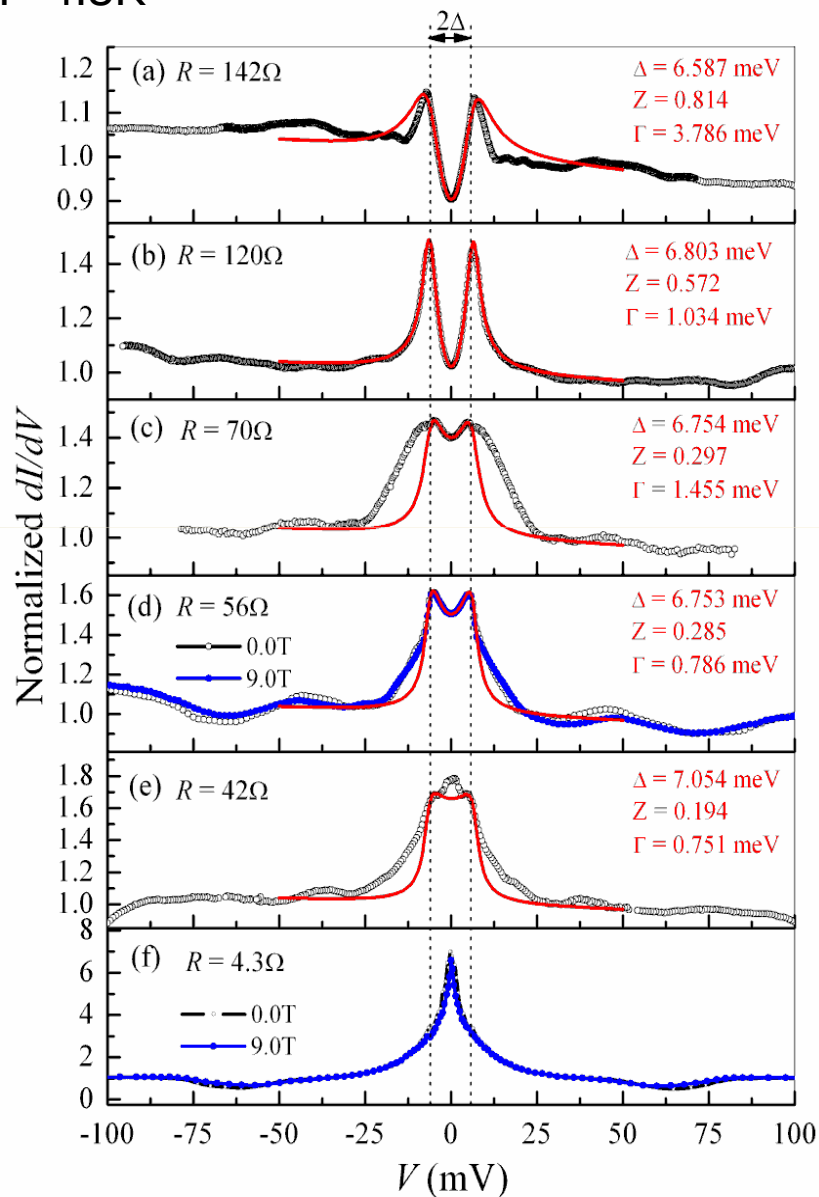
Superconducting $\text{SmFeAsO}_{1-x}\text{F}_x$ ($x=0.15, 0.18, 0.30$)



SmFeAsO_{0.85}F_{0.15} ($T_C = 42$ K)

Small
contact

$T = 4.5$ K



Large
contact

Same 2Δ for all $G(V)$

□ Structure:

~ isotropic

□ Value:

$$2\Delta = 13.34 \pm 0.71 \text{ meV}$$

$$2\Delta/k_B T_C = 3.68$$

close to 3.53

(BCS s-wave)

Complications:

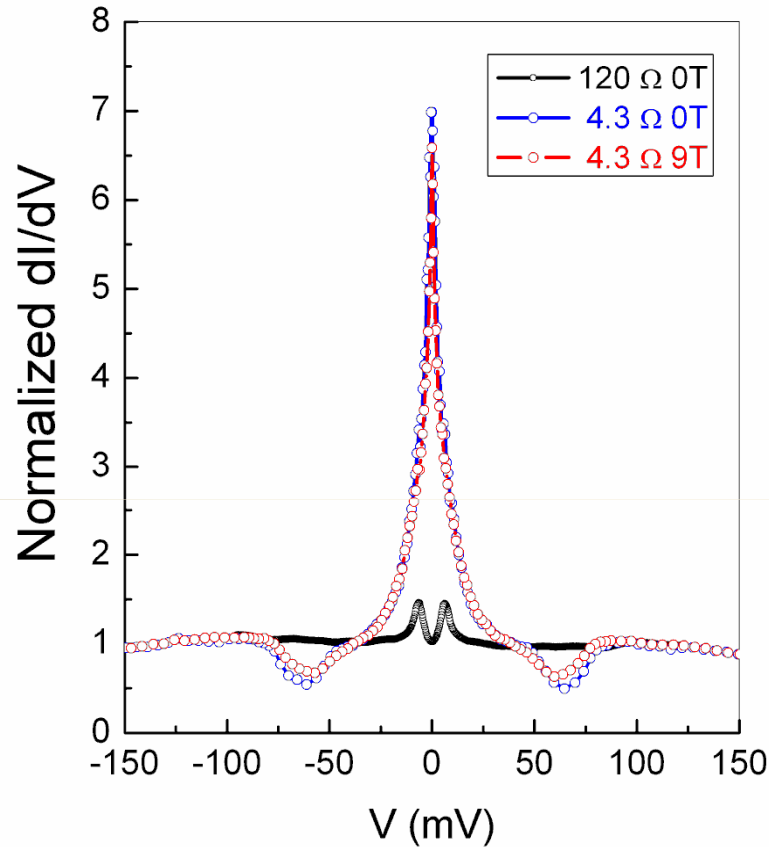
○ Extra features to ± 100 mV
due to SC

○ Slanted background
nothing to do with SC
mismatch of Fermi energy

○ Emergence of ZBA
due to SC

○ Small effect in 9 T field at 4.5K

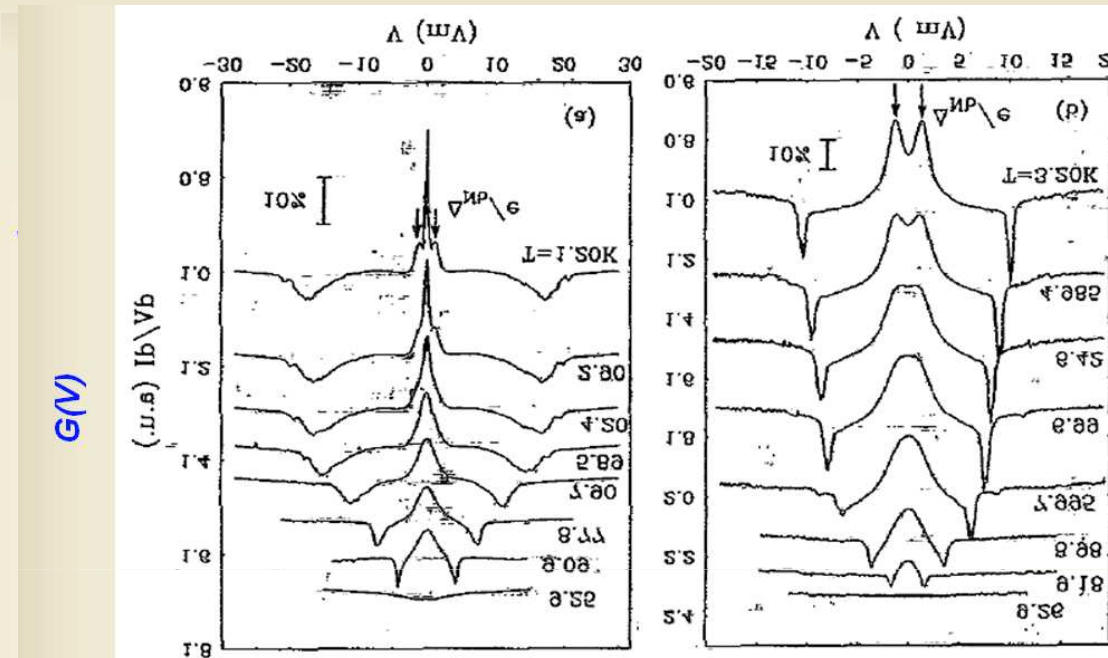
Emergence of Zero Bias Anomaly (ZBA)



- Emerges systematically
- ZBA obscures gap
- ZBA due to SC
- H field has negligible effect

- ❑ ZBA is a signature of *d*-wave pairing (???)
- ❑ ZBA depends strongly on *B* field, whereas gap does not (???)
(A scheme to obtain gap structure by subtraction)

ZBAs in s-wave Superconductors

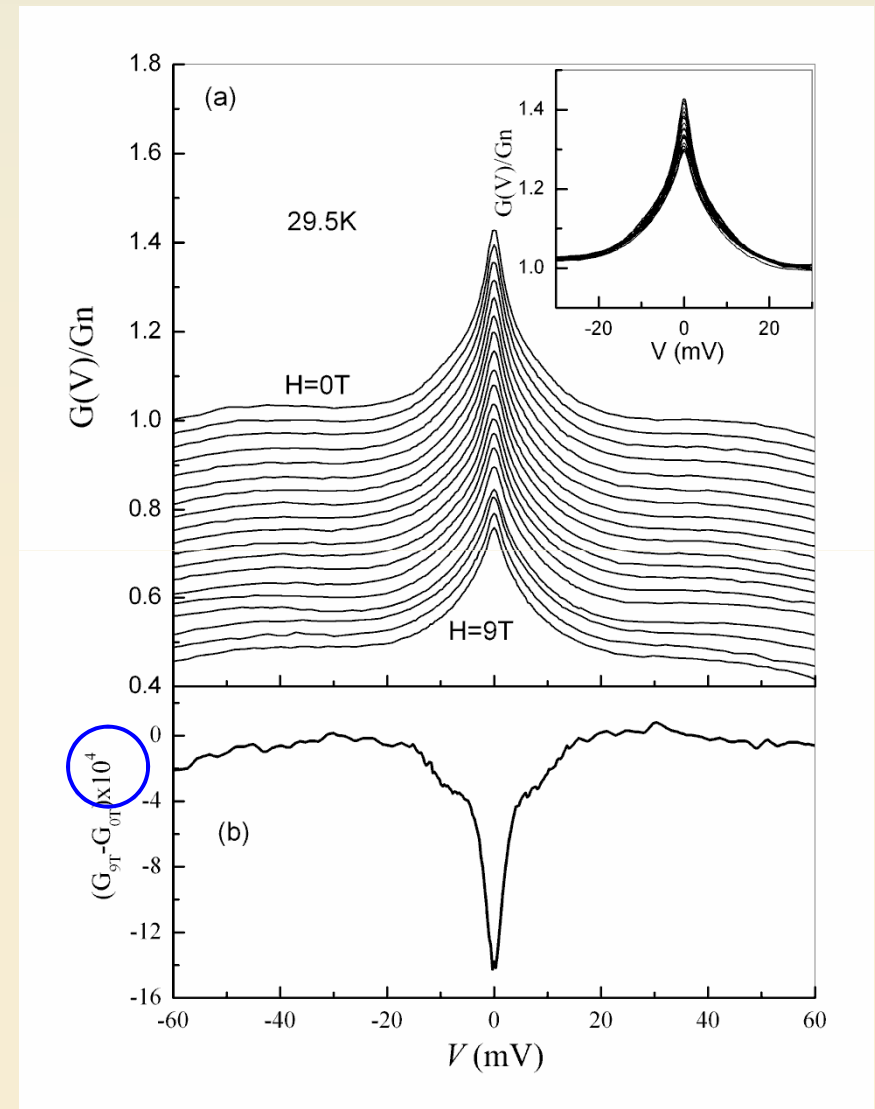
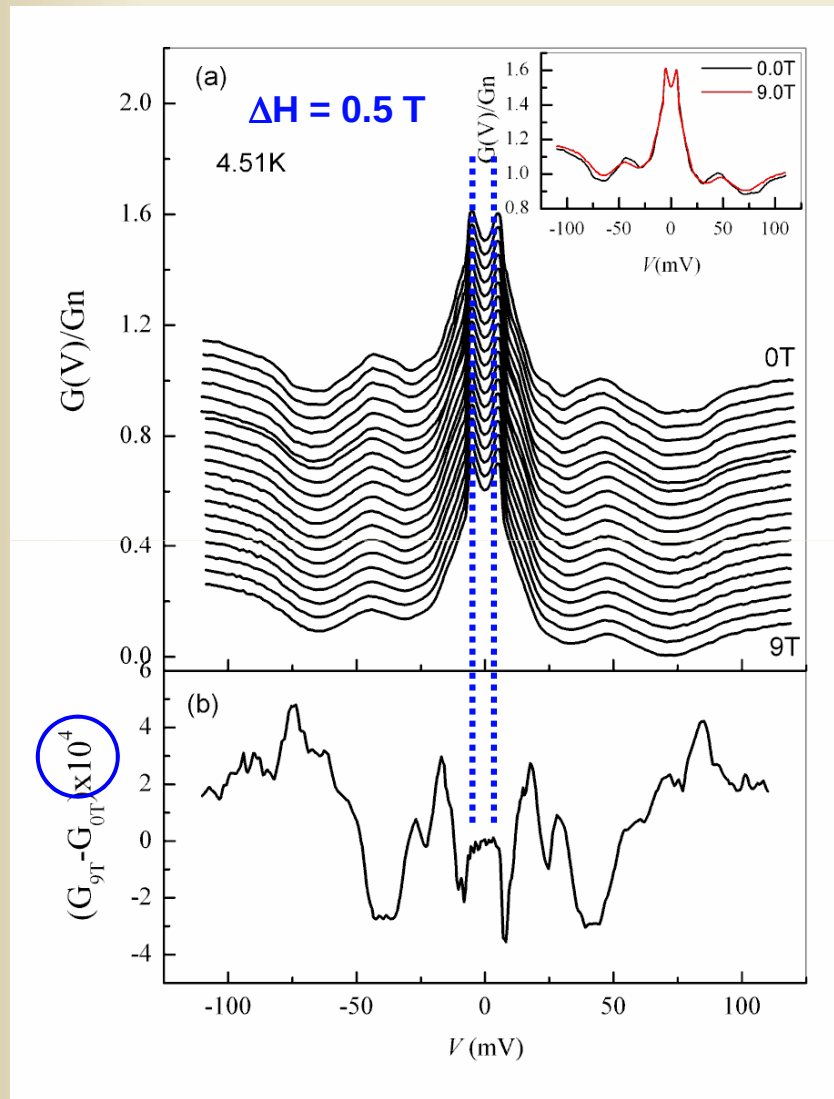


ZBA is a signature of *d*-wave pairing (???)

d-wave may give ZBA.

But, ZBA does not imply *d*-wave
since *s*-wave also gives ZBA.

Neither Gap nor ZBA depends strongly on H



Subtraction does not reveal the superconducting gap

Zero Bias Anomaly (ZBA)

- ❑ ZBA is a signature of *d*-wave pairing (???)

d-wave may give ZBA.

But, ZBA does not imply *d*-wave
since *s*-wave also gives ZBA.

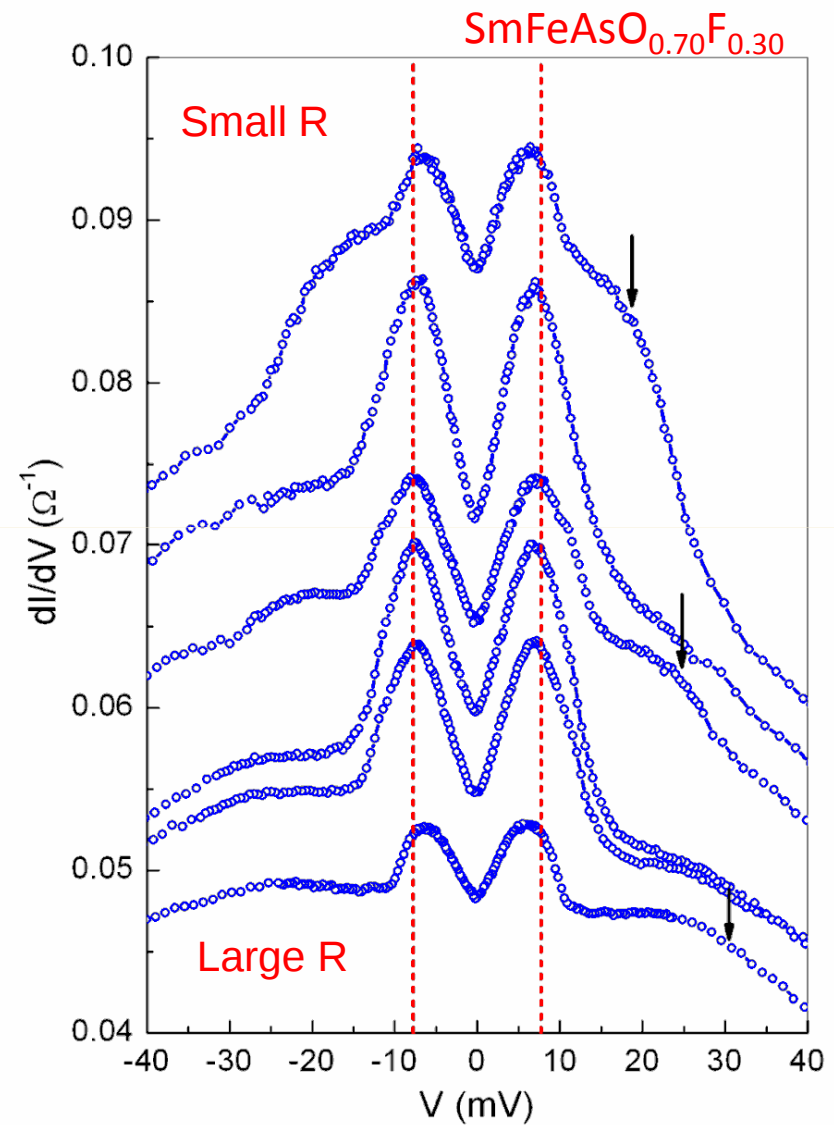
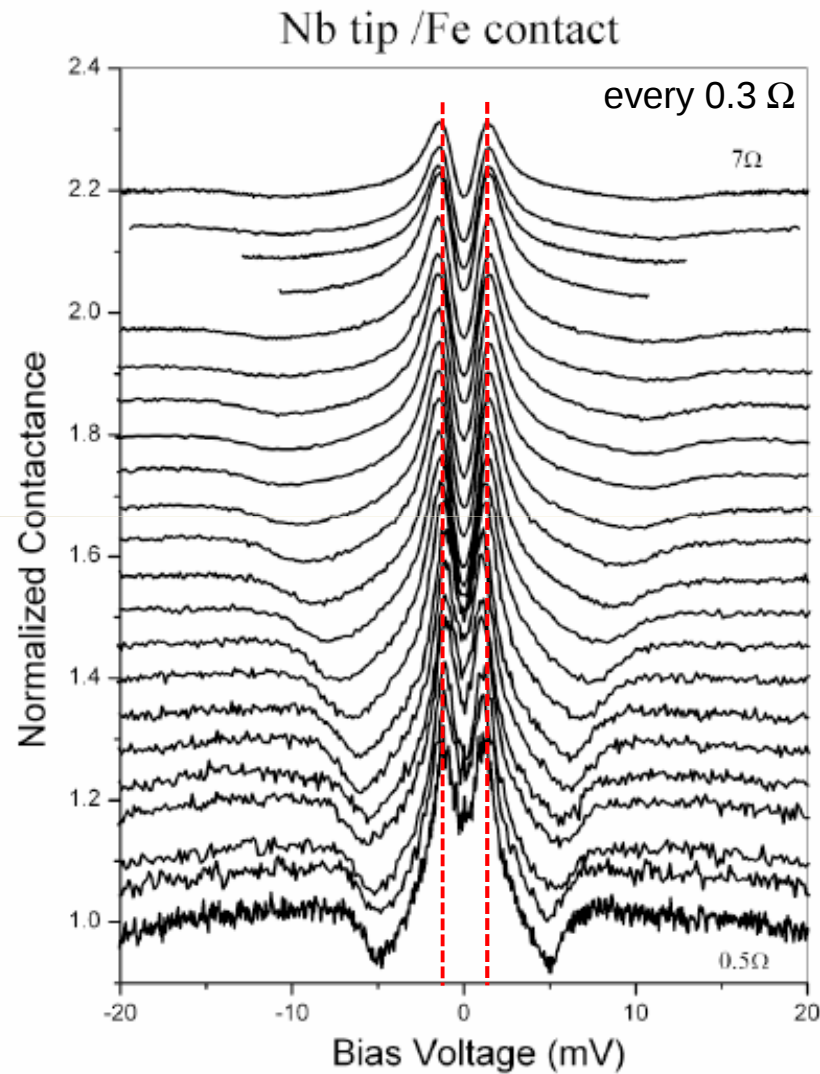
- ❑ ZBA depends strongly on *B* field, whereas gap does not (???)

(A scheme to obtain gap structure by subtraction)

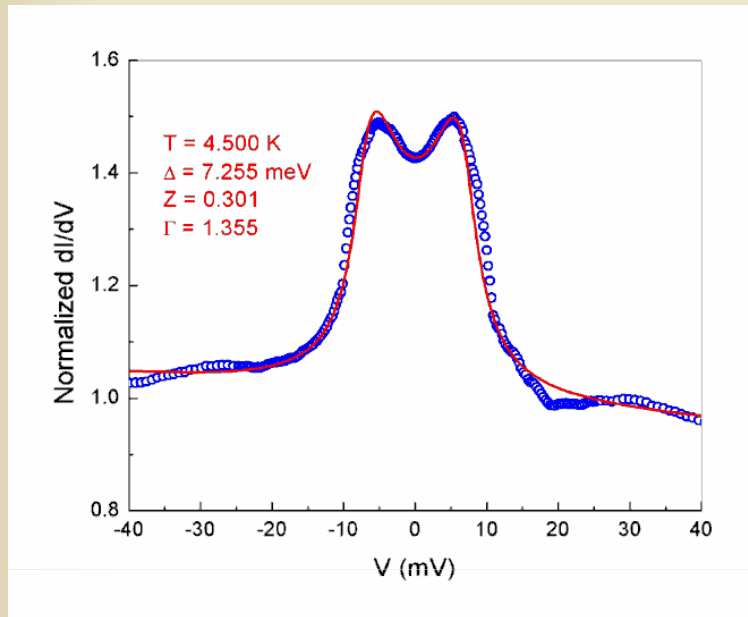
Neither ZBA nor gap depends strongly on *B* field.

Subtraction method does not work.

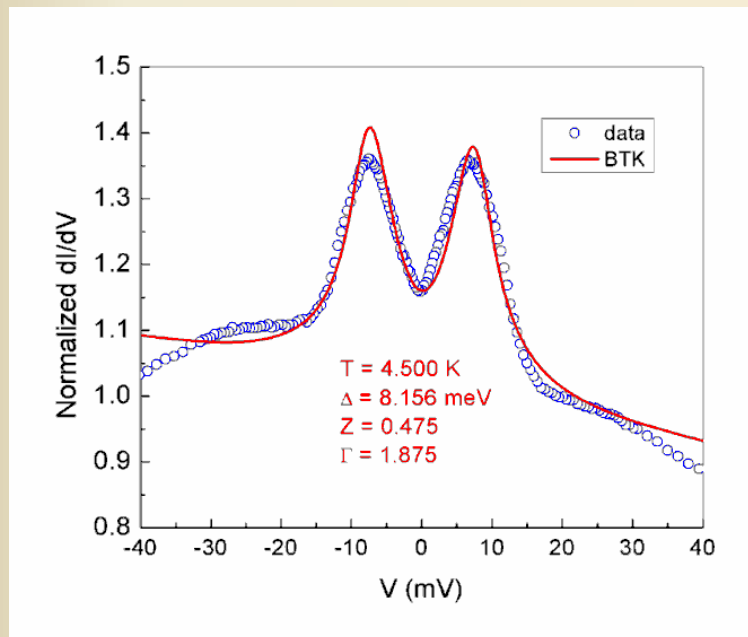
A real gap is independent of contacts



$\text{SmFeAsO}_{0.82}\text{F}_{0.18}$ ($T_C = 46.5$ K), $\text{SmFeAsO}_{0.70}\text{F}_{0.30}$ ($T_C = 53.75$ K)



$\text{SmFeAsO}_{0.82}\text{F}_{0.18}$:
 $T_C = 46.5$ K
 $\Delta = 7.26$ meV
 $2\Delta/k_B T_C = 3.62$



$\text{SmFeAsO}_{0.70}\text{F}_{0.30}$:
 $T_C = 53.75$ K
 $\Delta = 8.16$ meV
 $2\Delta/k_B T_C = 3.52$

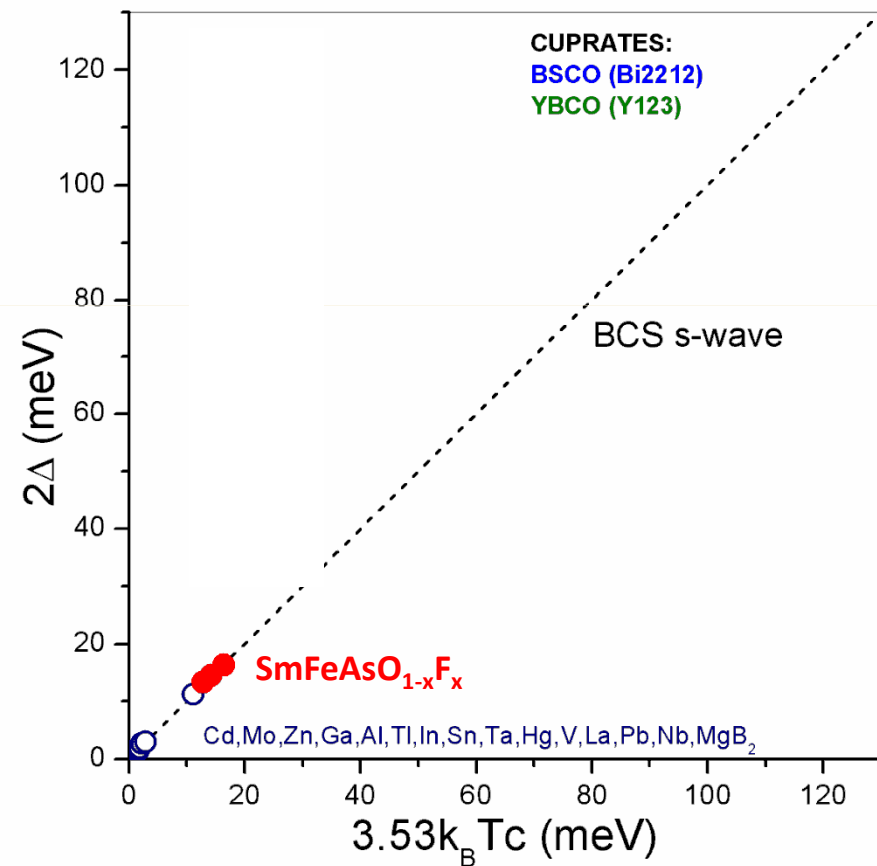
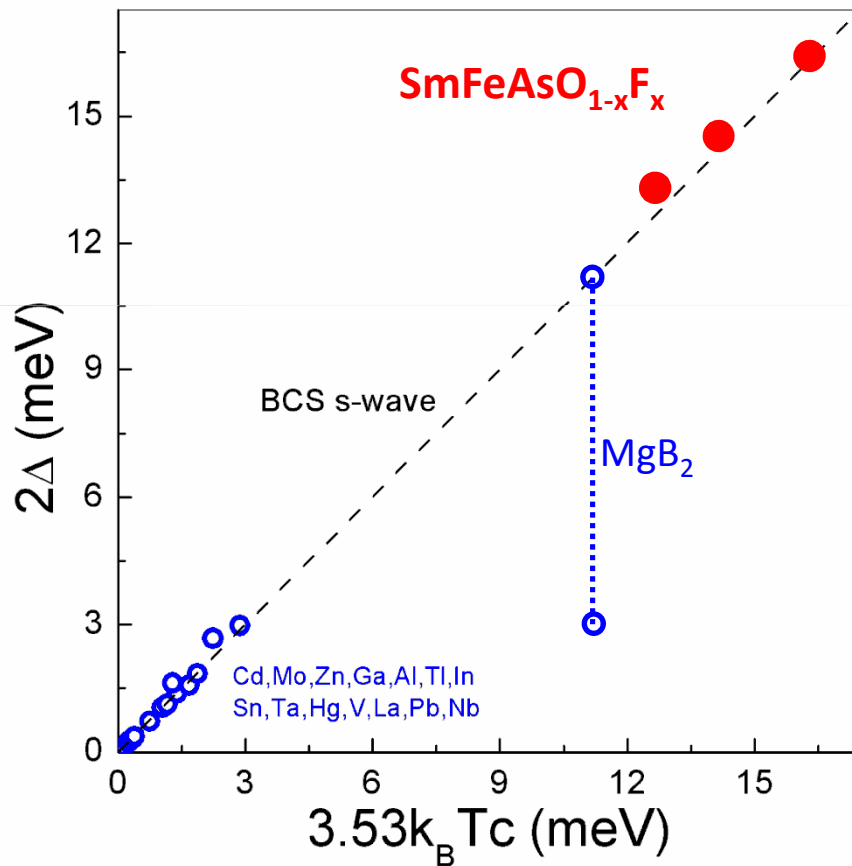
Superconducting Gap

Theory:

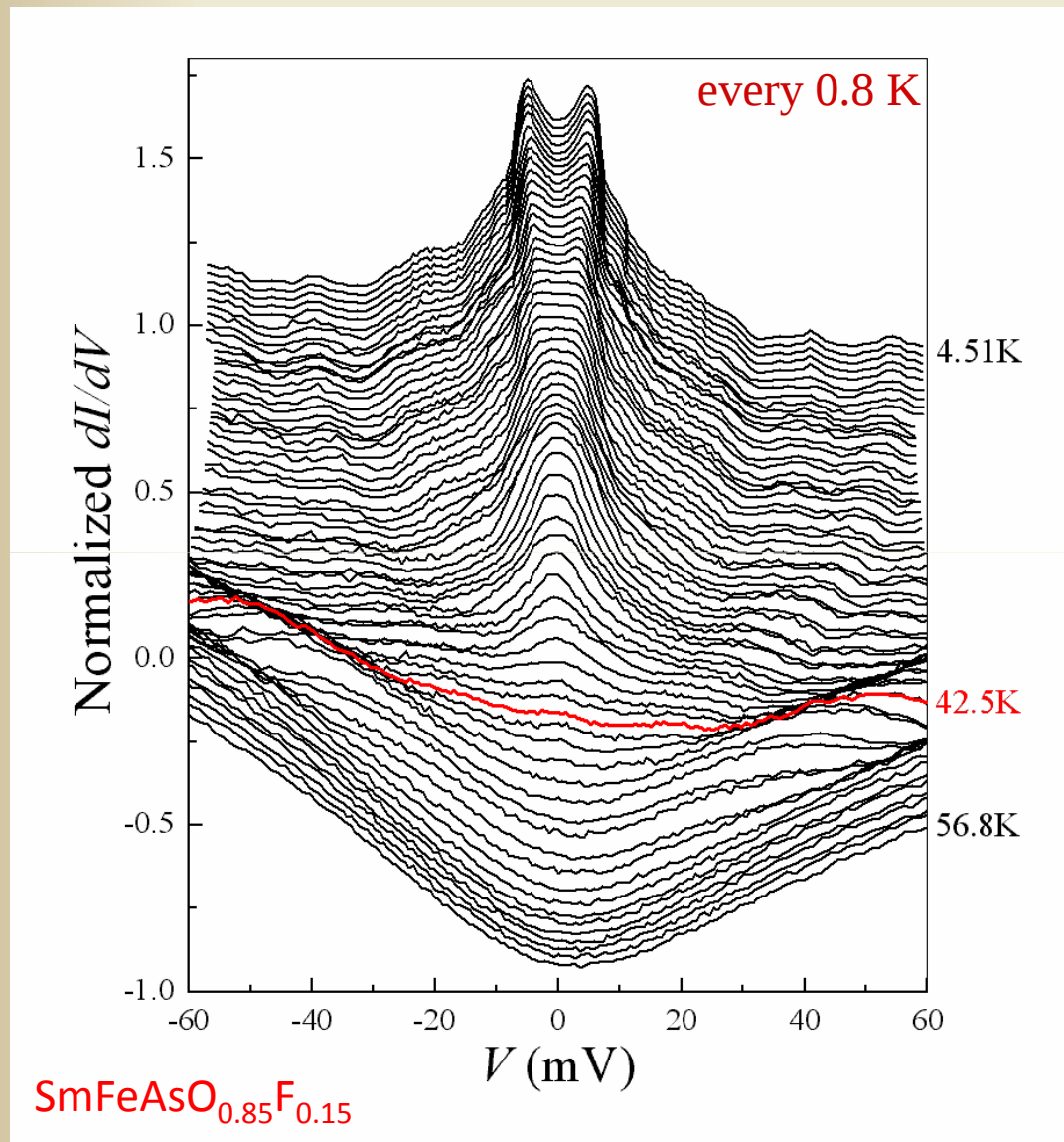
$$\frac{2\Delta(0)}{k_B T_c} = \begin{cases} 3.53 & \text{s-wave} \\ 4.06 & \text{p-wave} \\ 4.28 & \text{d-wave} \end{cases}$$

This work:

$$\frac{2\Delta}{k_B T_c} = \begin{cases} 3.68 & \text{SmFeAsO}_{0.85}\text{F}_{0.15} \\ 3.62 & \text{SmFeAsO}_{0.82}\text{F}_{0.18} \\ 3.52 & \text{SmFeAsO}_{0.70}\text{F}_{0.30} \end{cases}$$



Temperature-Dependence of Gap

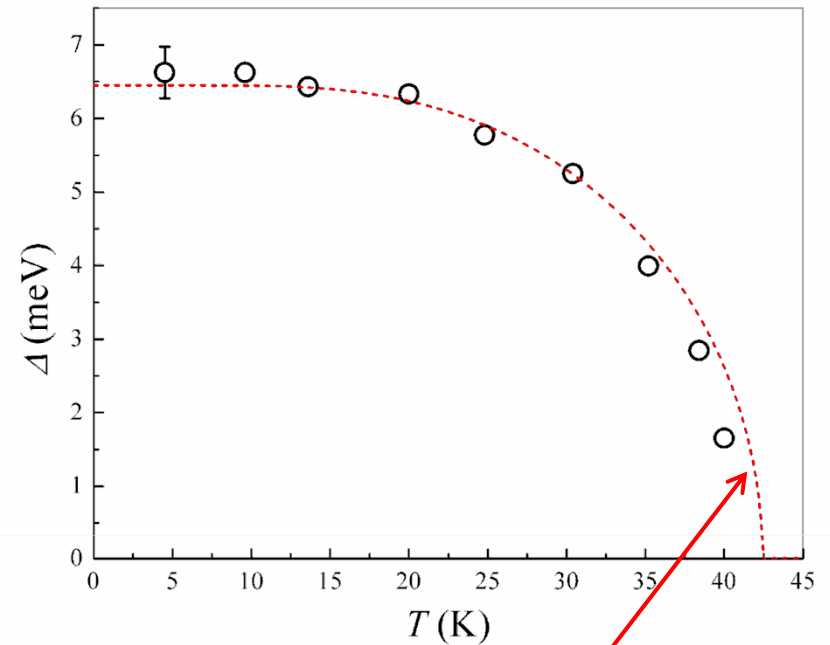
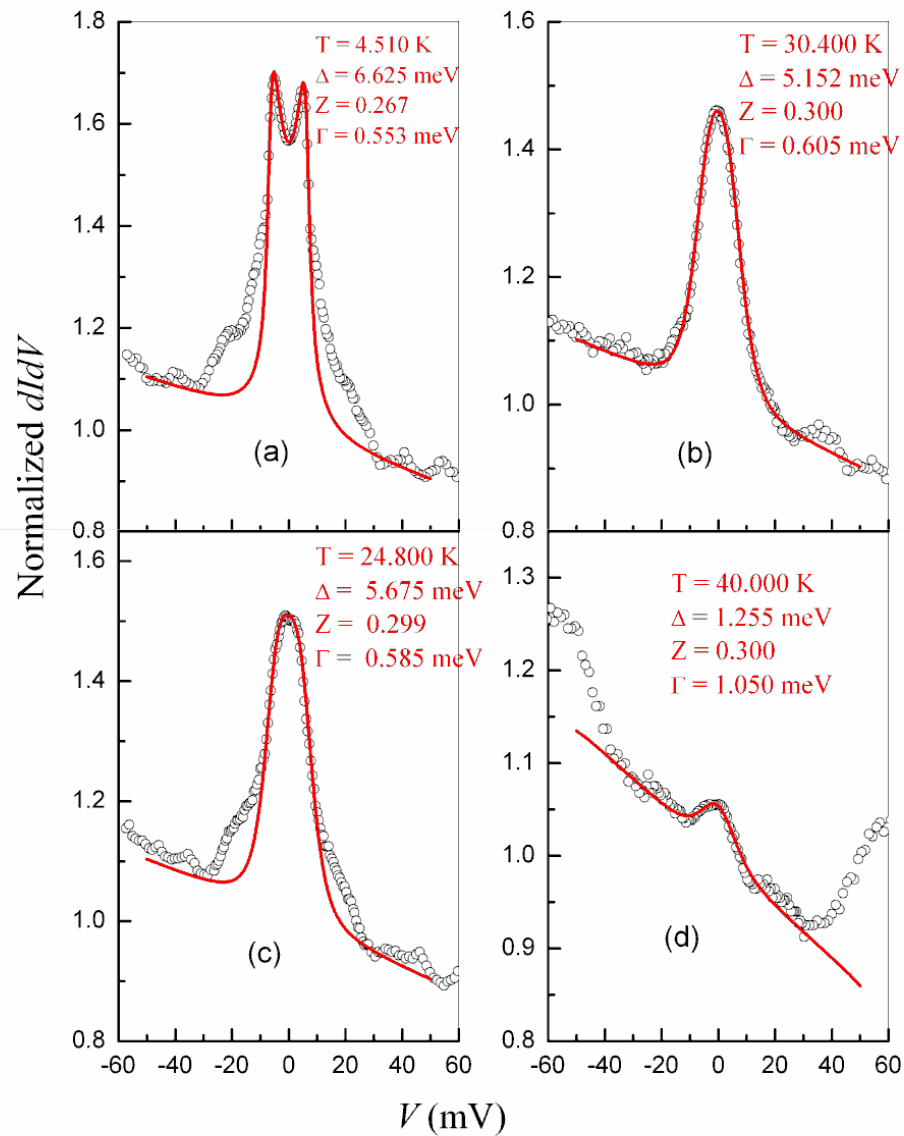


Double-peak (gap):
decreasing splitting
decreasing intensity
vanishes at $T_C = 42$ K

Extra features
(due to superconductivity)
vanishes at $T_C = 42$ K

Asymmetrical background
(due to normal state)
remains at $T > T_C$

BCS-Like Temperature Dependence

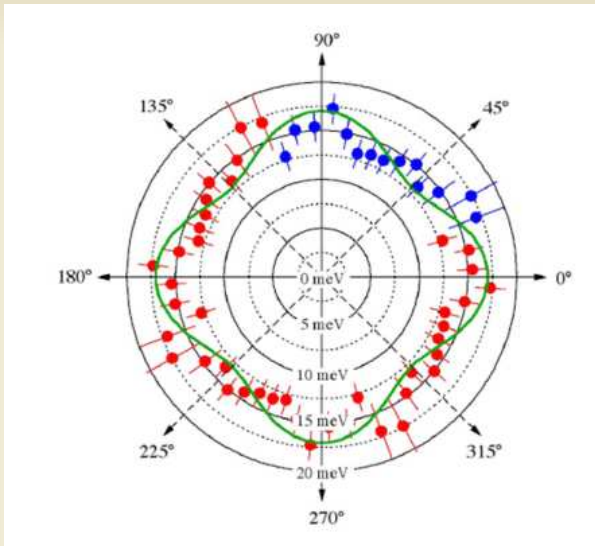


BCS s-wave

At each T :
 Gap value from analyzing data using
 modified BTK theory

Gaps from Other Measurements

ARPES

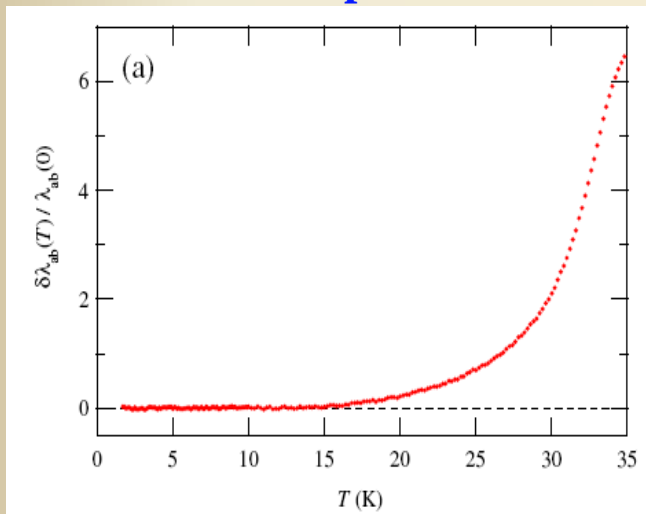


NdFeAsO_{0.9}F_{0.1} (0807.0815.Kondo)

Fully gapped, nearly isotropic **s-wave** symmetry

Nearly isotropic, fully gapped

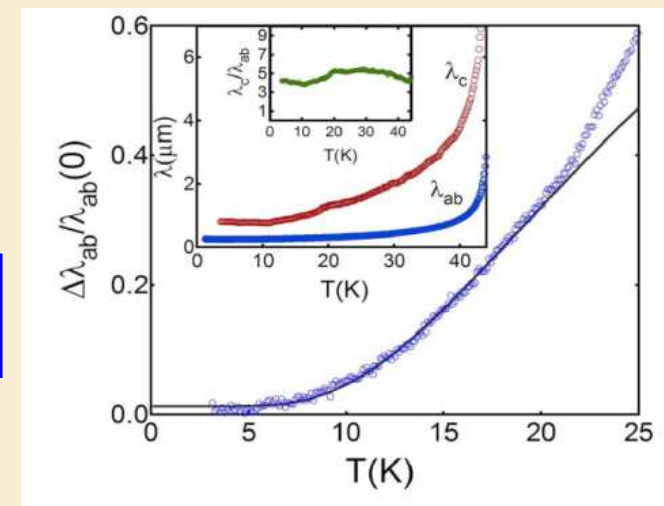
Penetration Depth



$$\frac{\Delta\lambda(T)}{\lambda(0)} = \sqrt{\frac{\pi\Delta_0}{2T}} \exp\left(-\frac{\Delta_0}{T}\right)$$

PrFeAsO_{1-y} (0806.3149.Hashimoto)

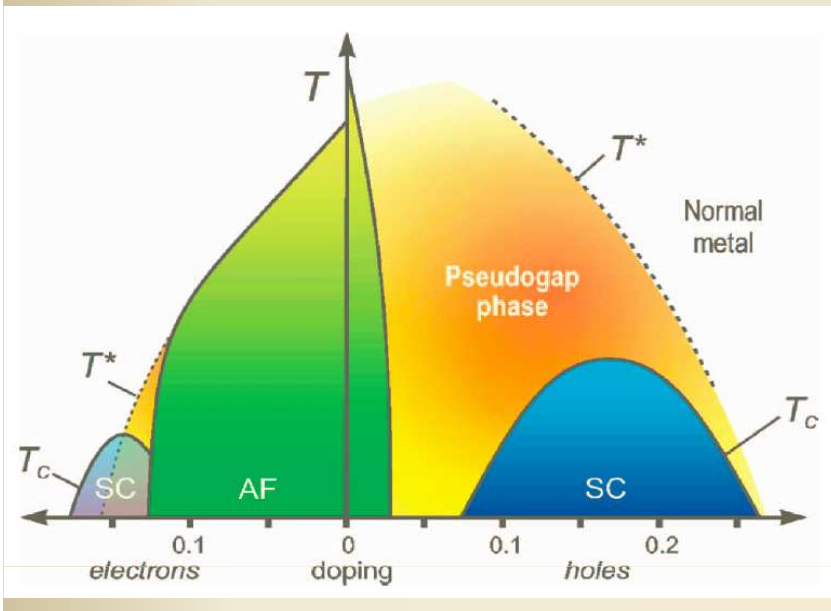
A single, fully gapped nearly isotropic gap



NdFeAsO_{0.9}F_{0.1} (0807.0876.Martin)

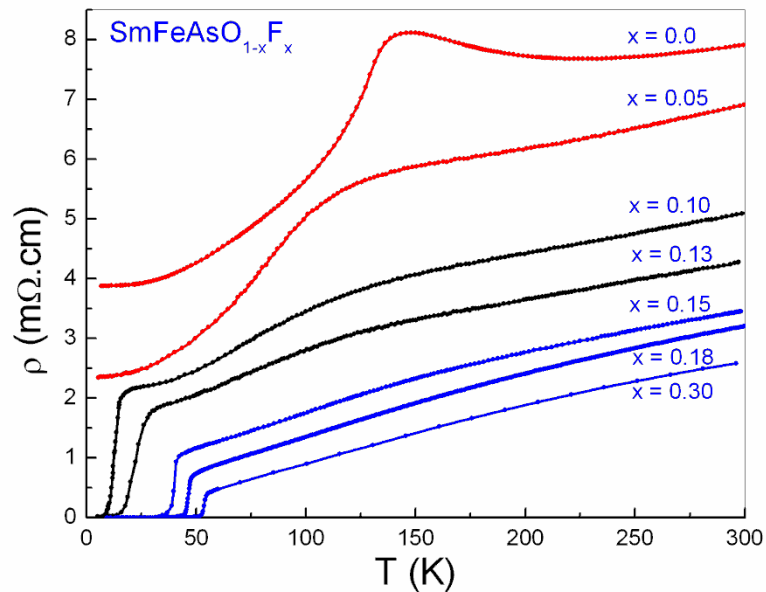
A single, fully gapped nearly isotropic gap

Pseudogaps ?



Pseudogap

in both **non-SC** and **SC** cuprates



No pseudogaps found in SC samples

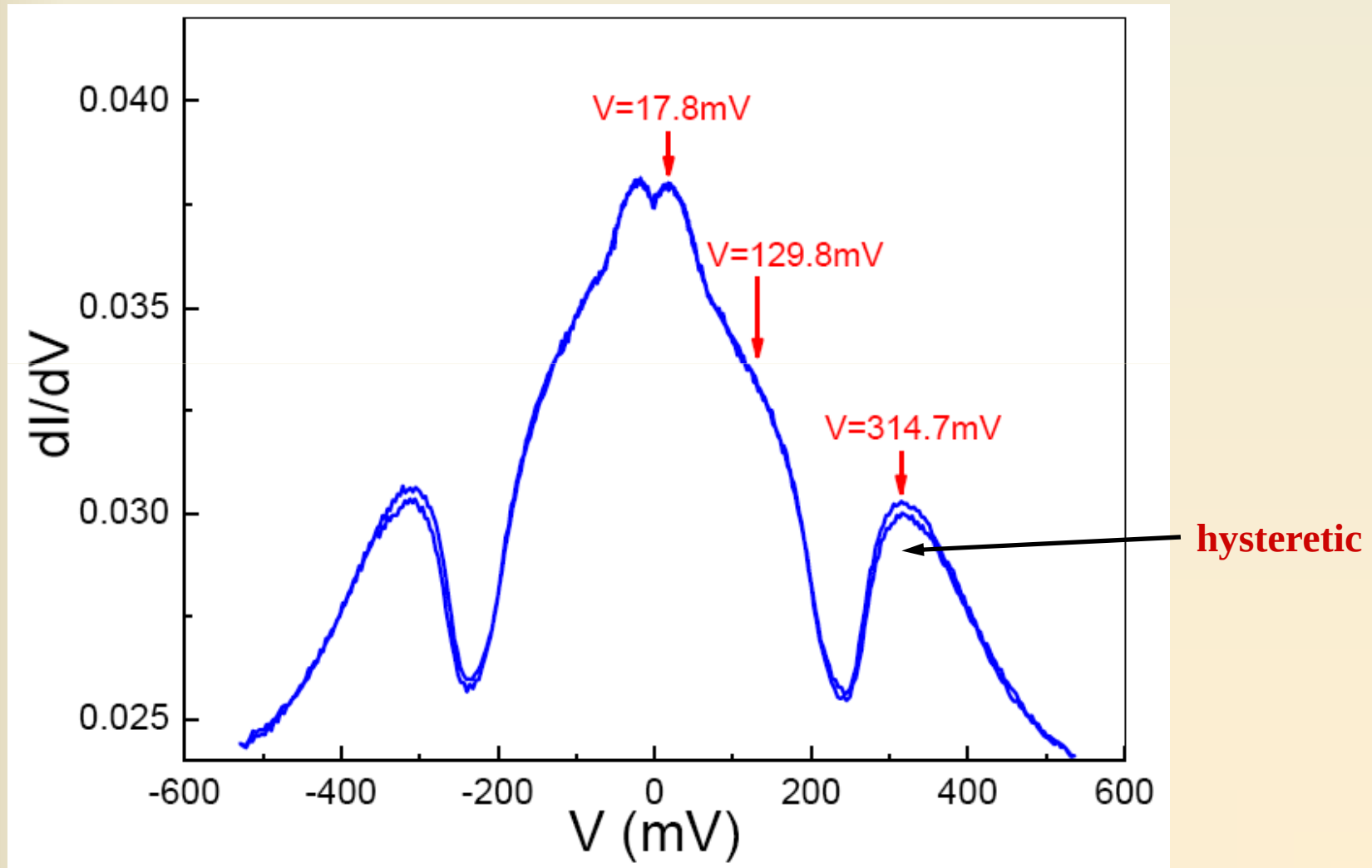
Non-SC samples ?

Characteristics of Pseudogaps

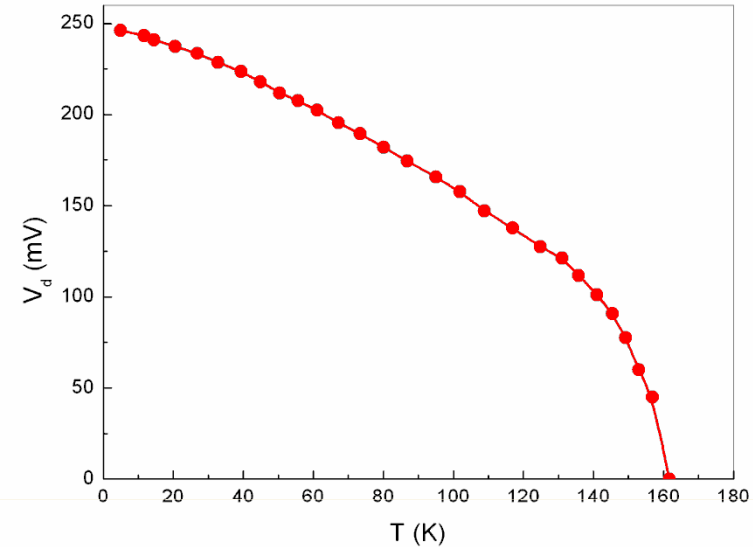
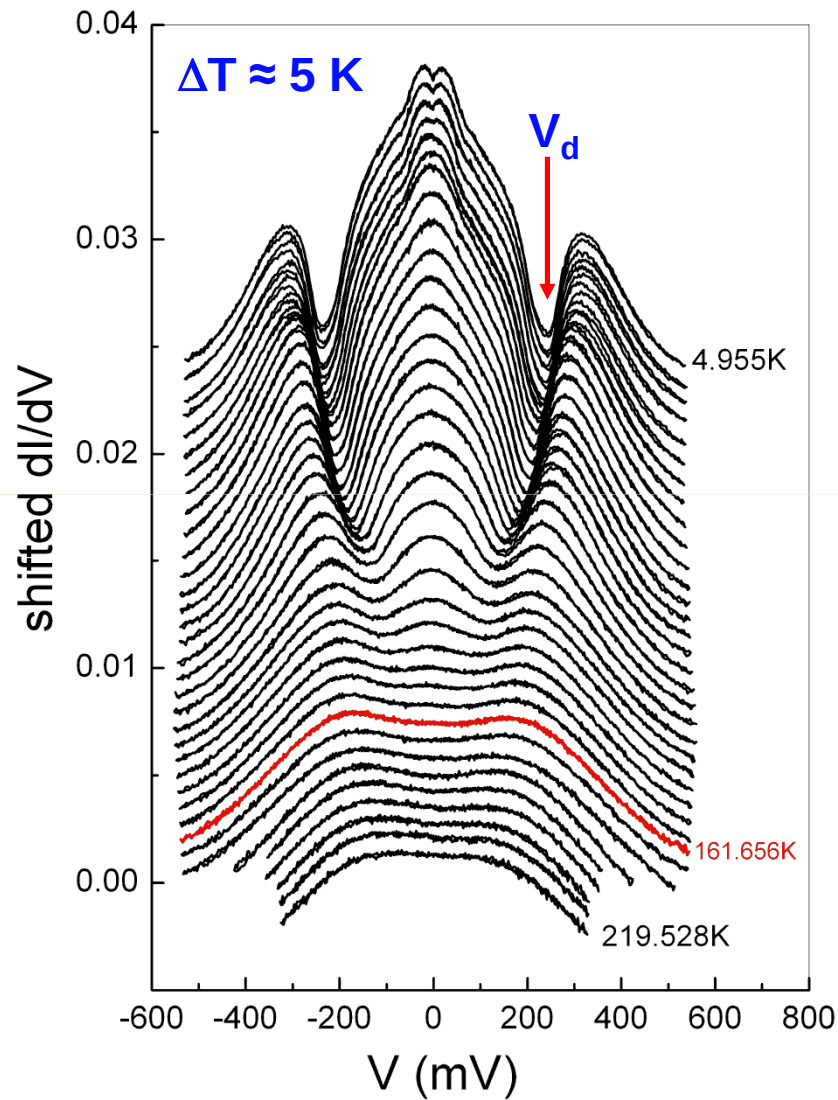
- ❑ Pseudogap may induce **conductance peaks**.
- ❑ Pseudogap should be **non-hysteretic**.
- ❑ Pseudogap must be **contact-independent**.

Conductance Peaks in parent SmFeAsO

Conductance peaks “suggesting” pseudogaps



Temperature Dependence of Conductance Peaks

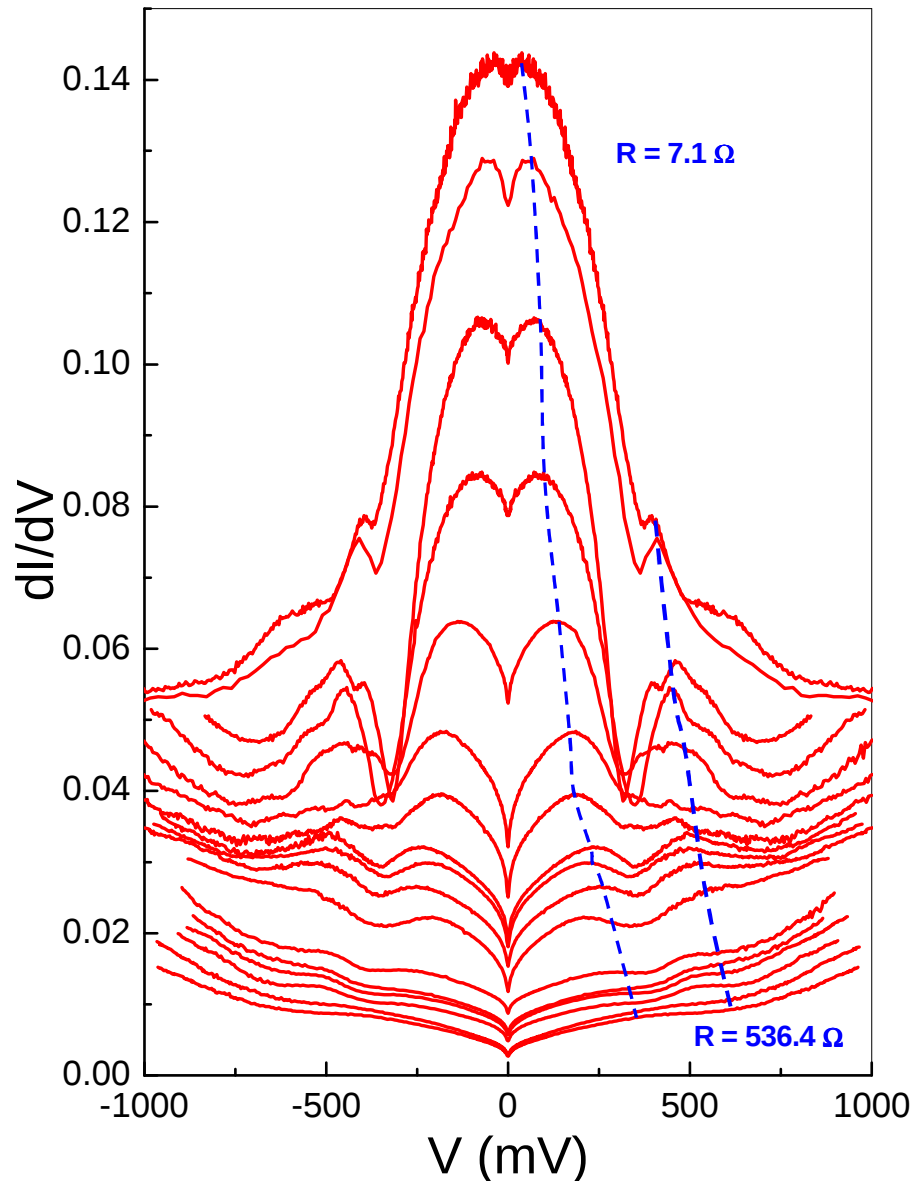


Structure disappears above 160 K

SmFeAsO with SDW at 160 K

Pseudogap for SDW ???

Systematic Contact-Dependent Conductance



Various contact resistance:
from **7 Ω** to **536 Ω**

Various conductance peaks:
from **3 mV** to **300 mV**

**Strongly depend on
contact resistance**

NOT pseudogaps!

Pseudogaps?

- ❑ **Pseudogap may induce conductance peaks**

But, **conductance peaks** do not imply **pseudogaps** since **resistance anomaly** also gives **conductance peaks**.

- ❑ **Pseudogap should be non-hysteretic.**

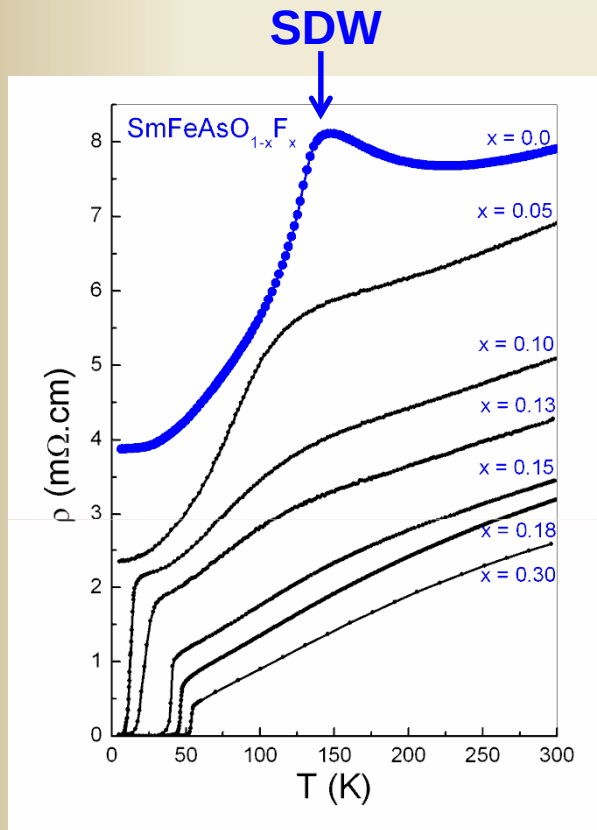
Spin density wave (**SDW**) results in *hysteretic behavior* in magnetic structure, crystalline structure, and resistance.

- ❑ **Pseudogap must be contact-independent**

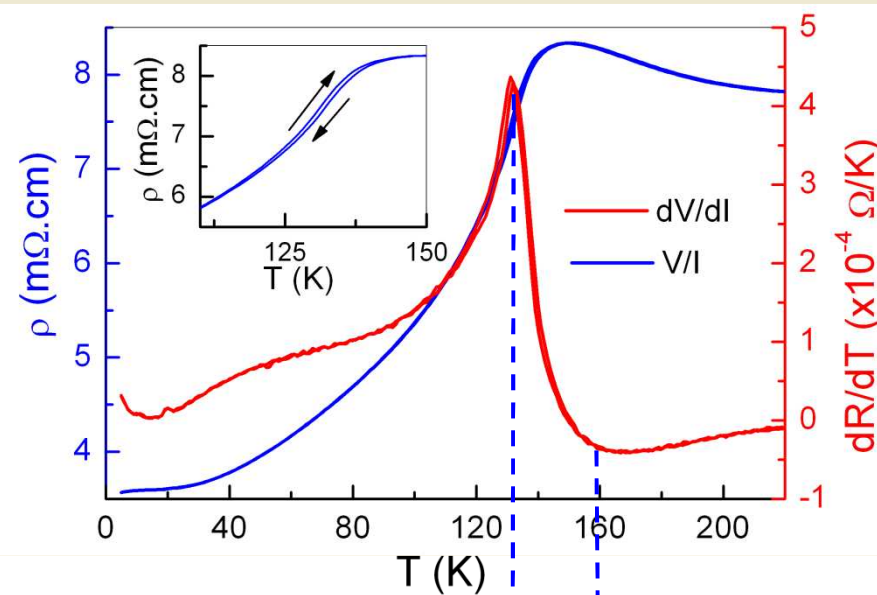
Contact-dependent peaks are not pseudogaps

Not pseudogaps, what are they ?

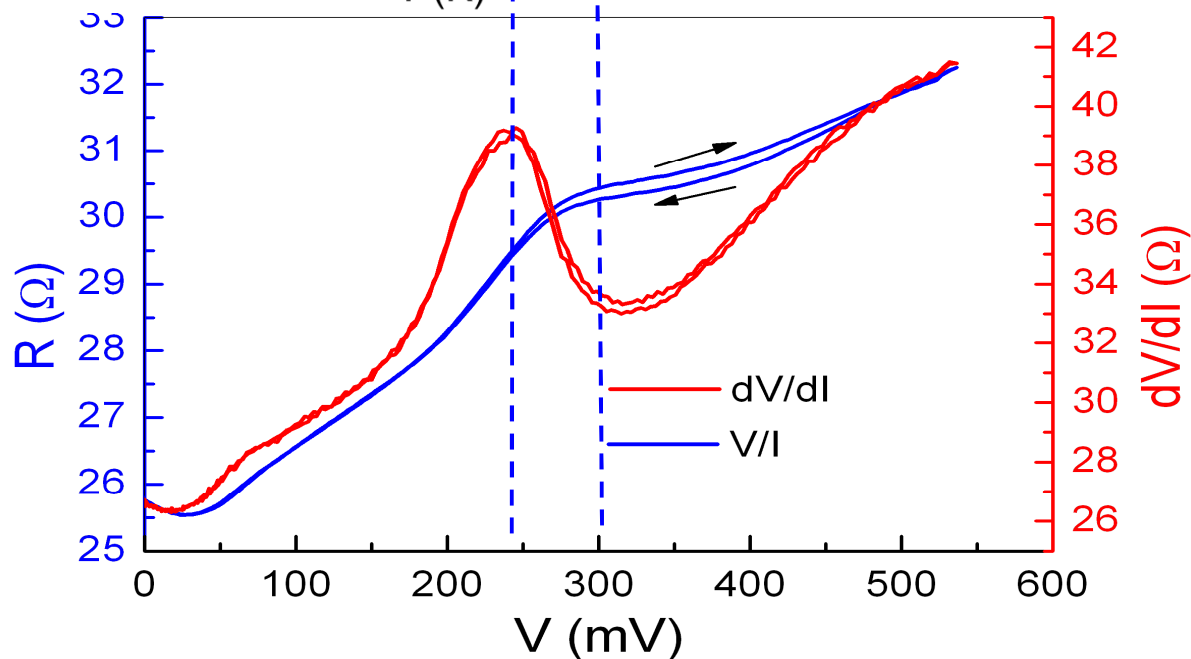
Conductance Peaks due to Ballistic Heating !



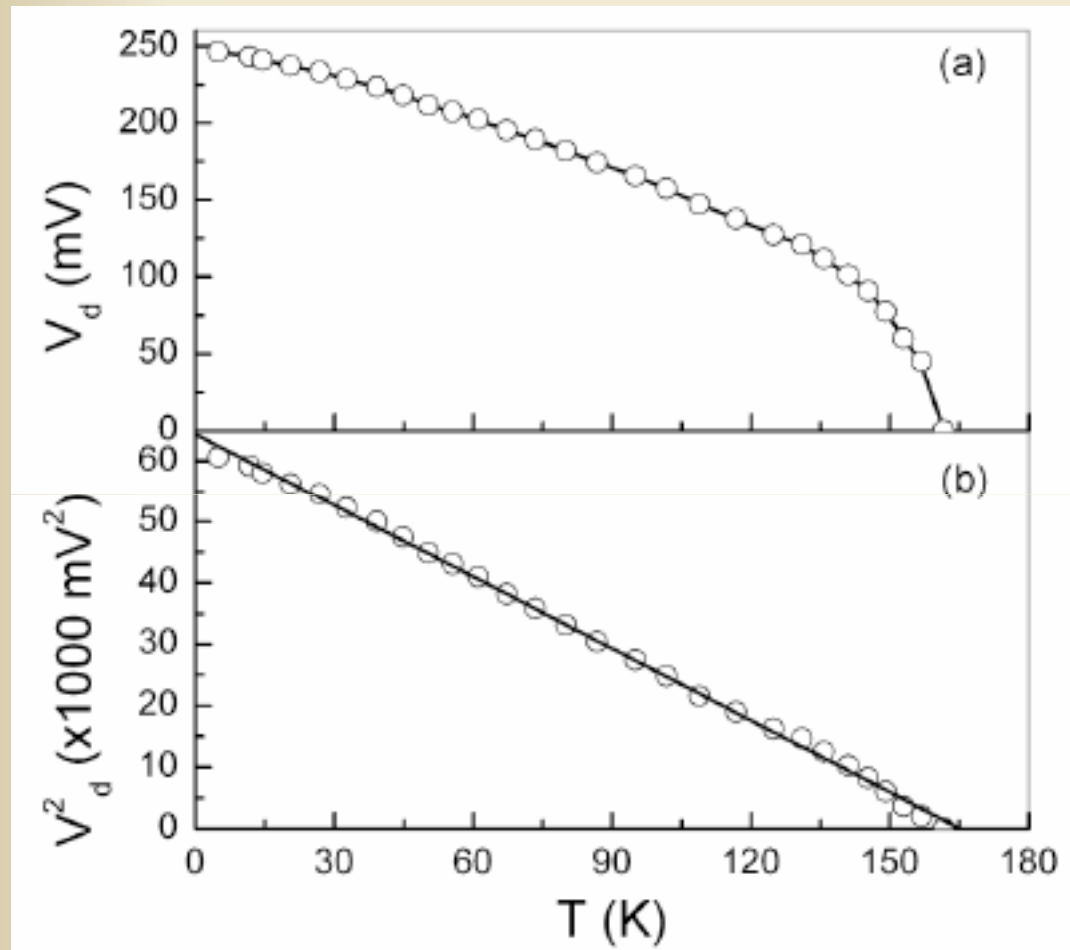
Point contact on
SmFeAsO



SmFeAsO with
SDW at 160 K



Temperature dependence of “Pseudogap”



**SmFeAsO with
SDW transition at 160 K**

$$V^2 \sim (T_{\text{SDW}} - T)$$

Ballistic heating effect

NOT pseudogaps, but SDW transition!

Summary

Superconducting gap (Δ) of **SmFeAsO_{1-x}F_x** (**x = 0.15, 0.18, 0.30**):

◆ Gap value

$$2\Delta/k_B T_C = 3.68, 3.62, 3.52$$

close to **3.53** (BCS s-wave)

◆ Gap structure

$\approx$ Isotropic single gap

◆ Gap temperature dependence

BCS-like

◆ No pseudogaps

Perhaps due to interband $S_{\pm}$ mechanism

Acknowledgements

Prof. **Chia-Ling Chien**, Advisor, Johns Hopkins University

Sunxiang Huang, JHU

Dr. **Mark Stiles**, NIST

Prof. **Zlatko Tesanovic**, Johns Hopkins University

Prof. **Xianhui Chen**, University of Science and Technology of China

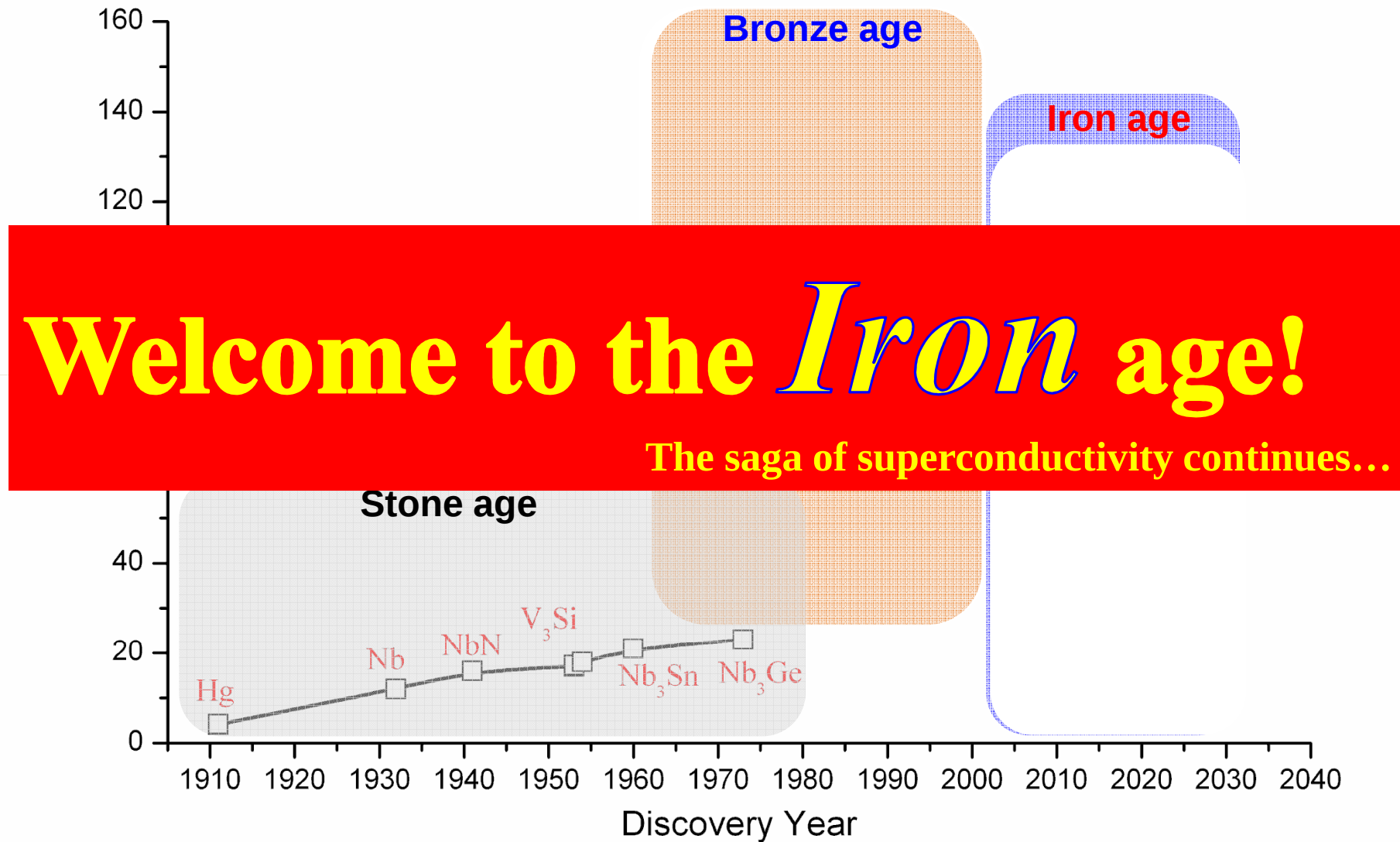
Prof. **Yi Ji**, University of Delaware

Fe Superconductors

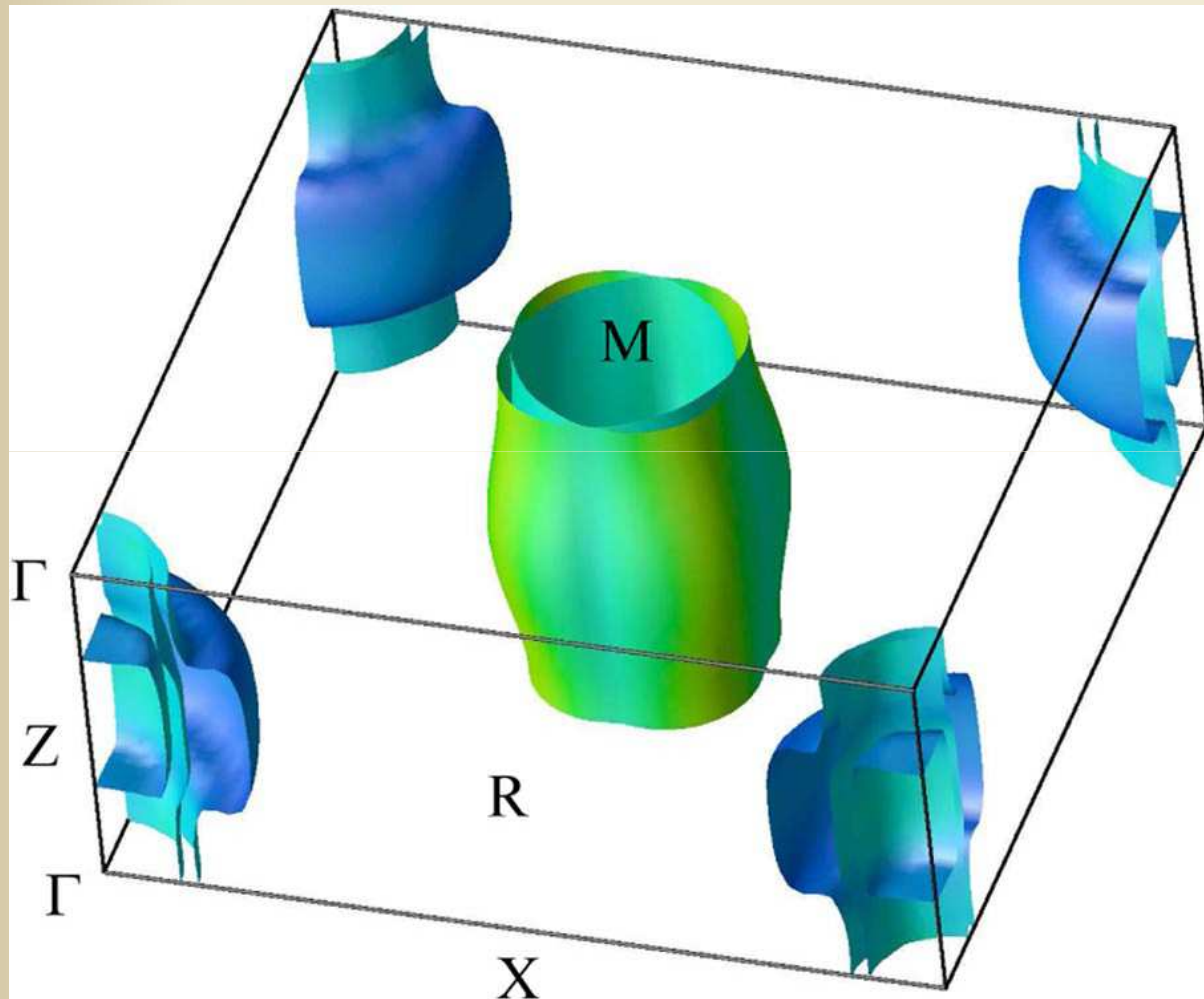
- ❑ Large varieties of materials
 - **1111**: LaFeAsO, CeFeAsO, SmFeAsO, ...
 - **122**: BaFe₂As₂, CaFe₂As₂, ...
 - **111**: LiFeAs, ...
 - **11**: FeSe, FeTe, ...
 - More to come...
- ❑ Highest **T_c~55K** achieved in **SmFeAsO_{1-x}F_x** (“single layer”)
- ❑ **Metallic** parent compound with common **SDW** feature
- ❑ Fully **gapped**, nearly **s-wave** symmetry
- ❑ FeAs plane can be doped with **Co**

Very different from Cuprates!

Saga of Superconductivity



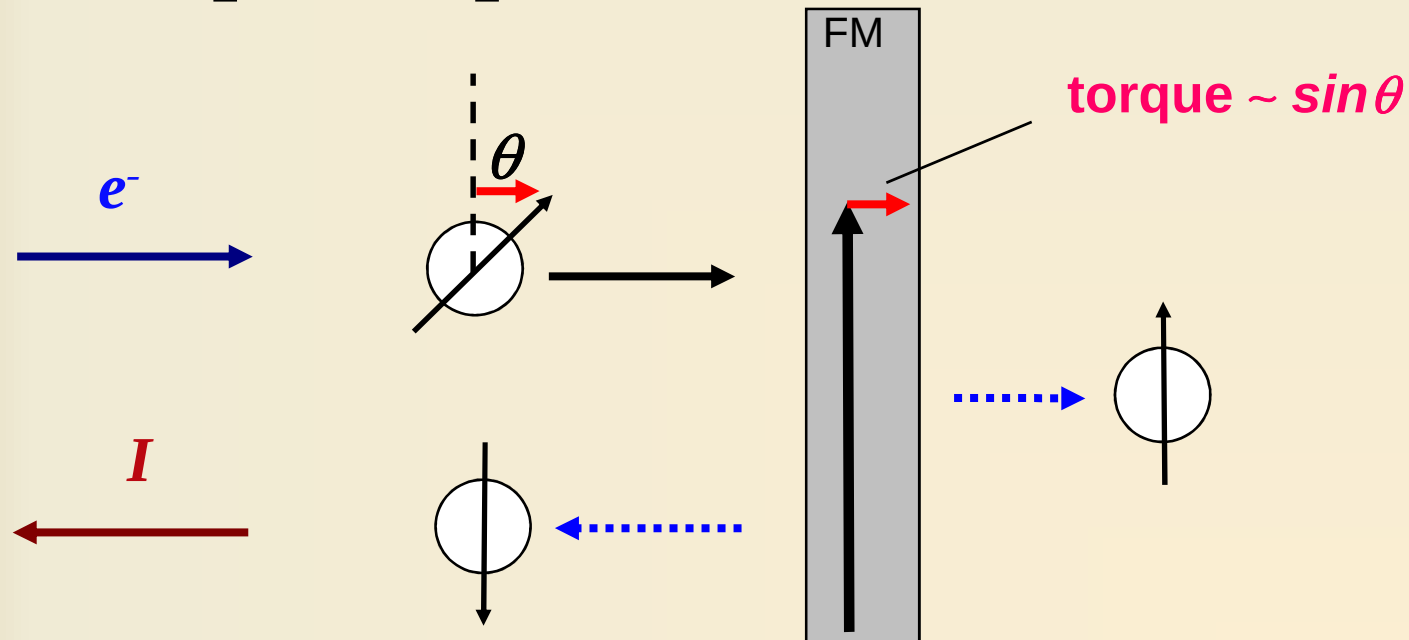
Fermi Surface of Fe-SC



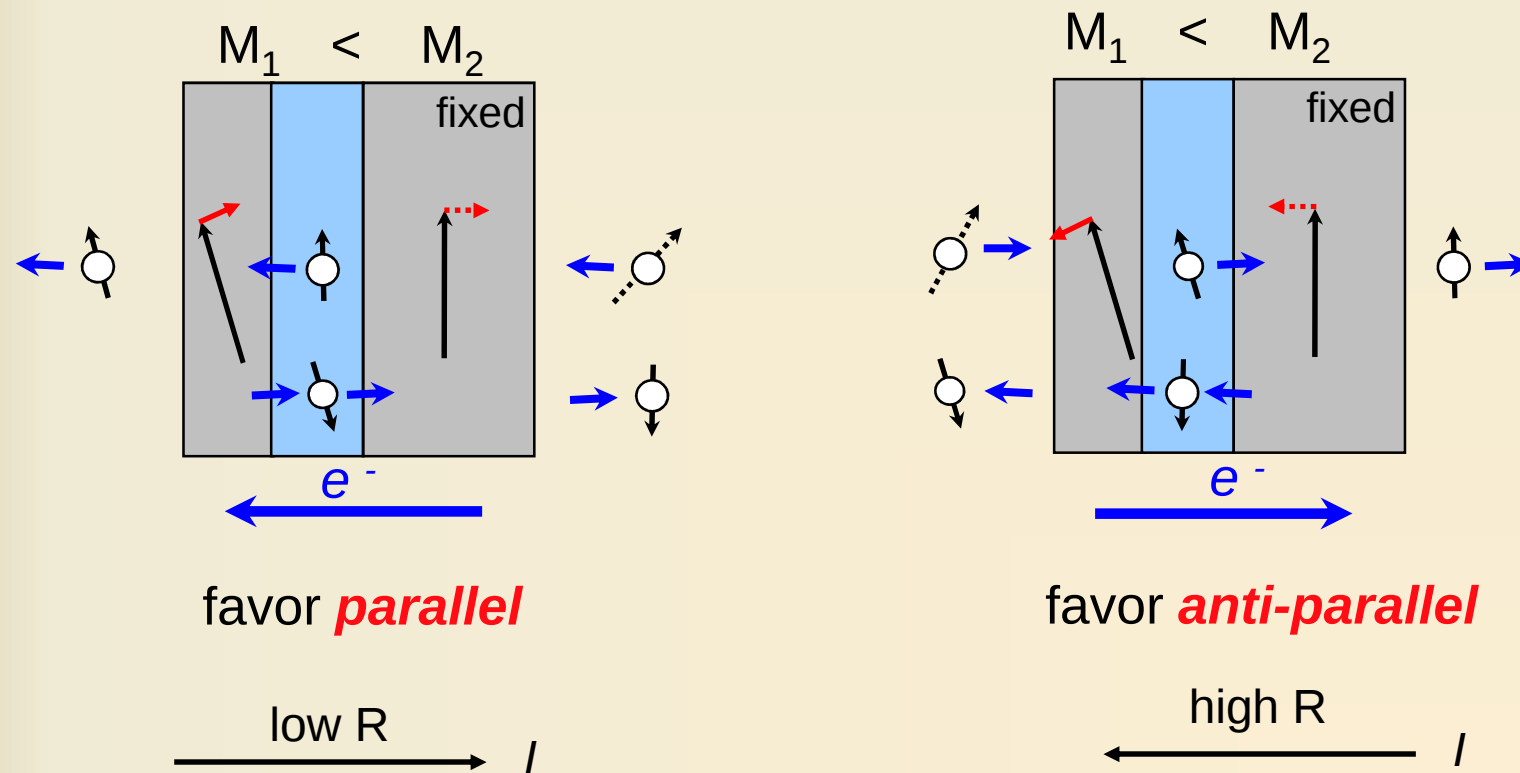
5 sheets:
2 electron sheets at MA
2 hole sheets at Γ
1 hole pocket at Z

Spin-Transfer Torque: Single Electron Scenario

$$\psi = \cos \frac{\theta}{2} |\uparrow\rangle + \sin \frac{\theta}{2} |\downarrow\rangle$$



Switching via Spin-Transfer Torque

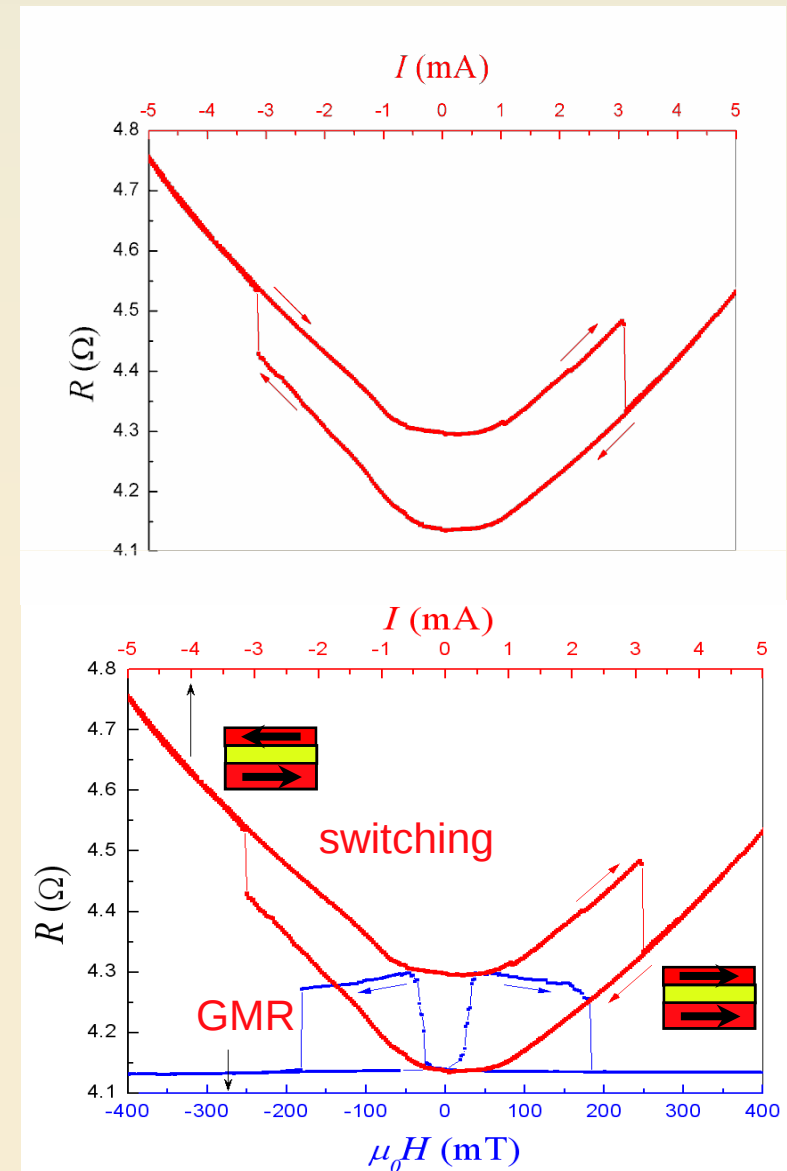
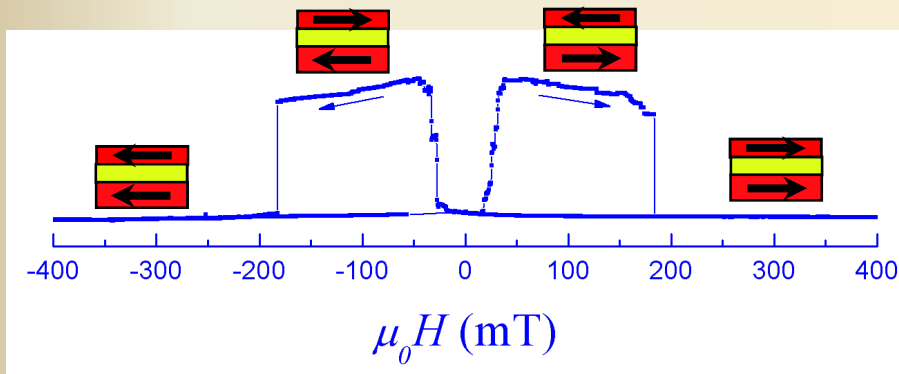
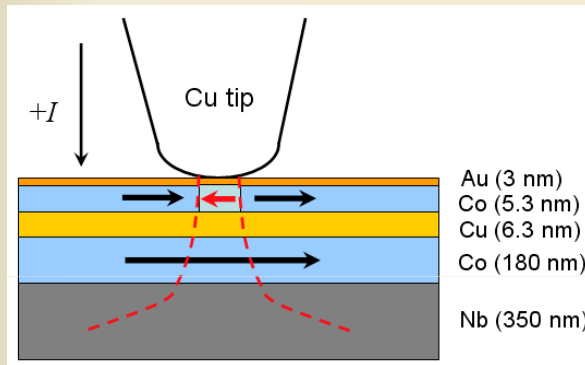


detect switching through GMR effect

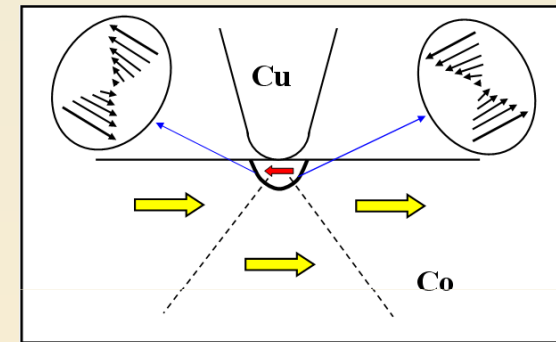
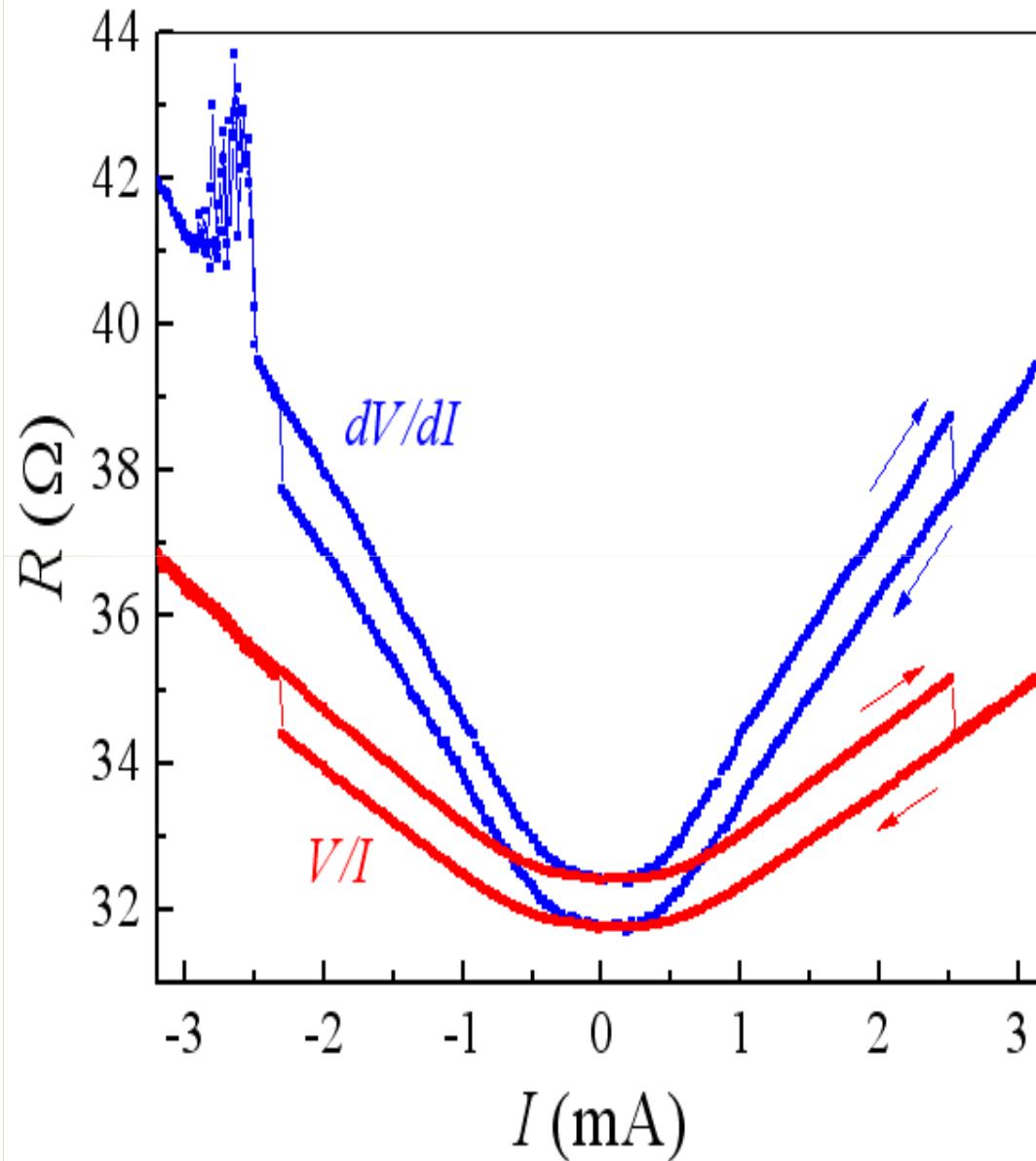
Switching via Spin Transfer Torque (STT) Effect

$$j_C \sim 10^7 \text{ A/cm}^2$$

Nano-pillar, point contact

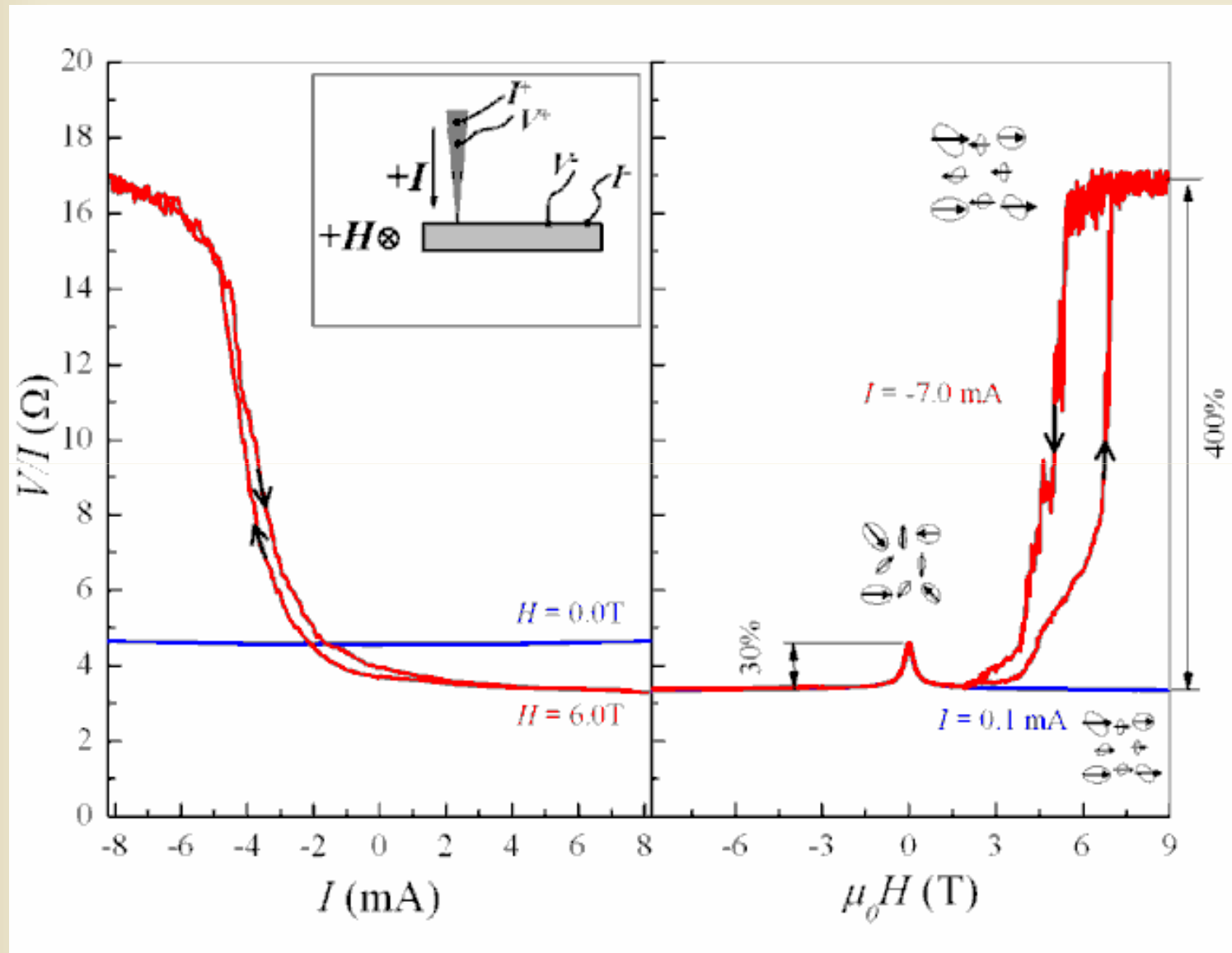


STT in Magnetic Nanostructures



New form of STT effect
Inverse of the domain wall MR

New Spin Structure Realized by STT



New spin structure, largest MR of over 400%

T. Y. Chen, S. X. Huang, C. L. Chien and M. D. Stiles, *Phys. Rev. Lett.* **96**, 207203 (2006).

Cryogenic Dewar

- ☐ Temperature control: 1.5 K to 400 K
- ☐ Magnetic field: 0 - 9 T
- ☐ Rotating magnetic field without mechanical motion
- ☐ Nitrogen jacket

- ☐ Sensitive transport system
 - Two lock-in from Stanford research
 - 10-channel switching box
 - Current source, nano-voltmeter

Purposes

- ☐ Point contact spectroscopy
- ☐ Transport study of nanostructure
- ☐ Spin injection and detection studies

Sputtering System

Specific features:

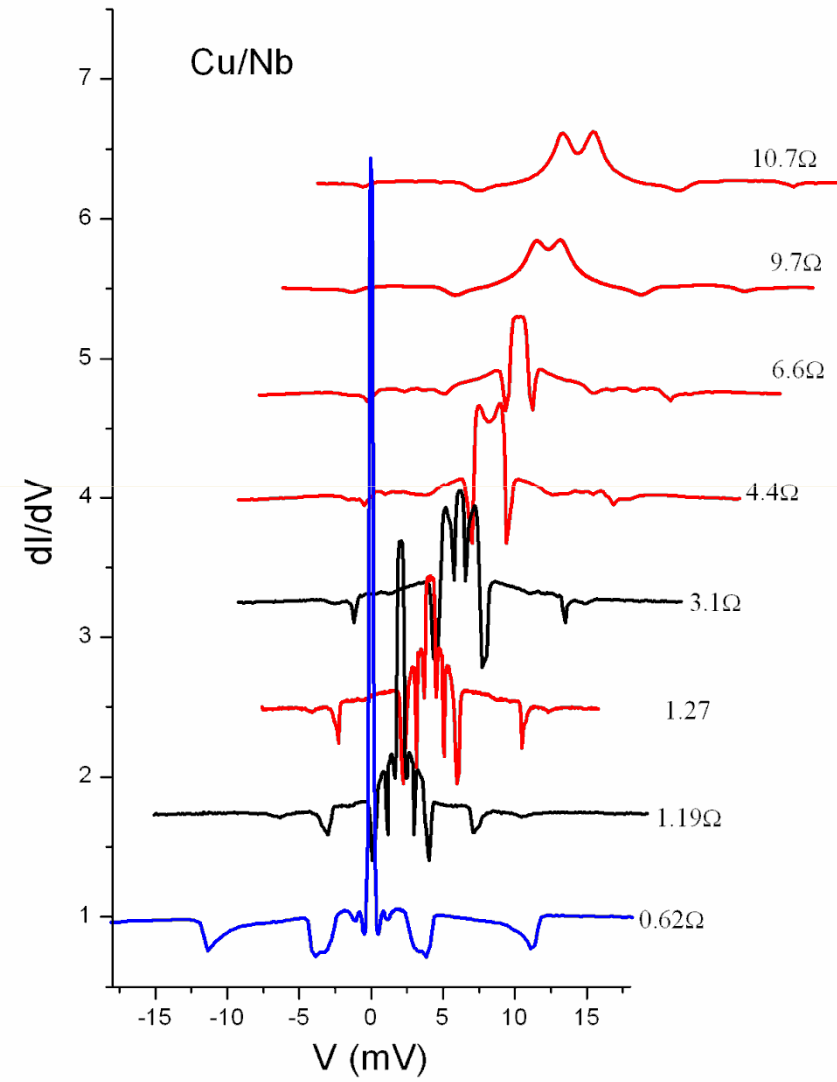
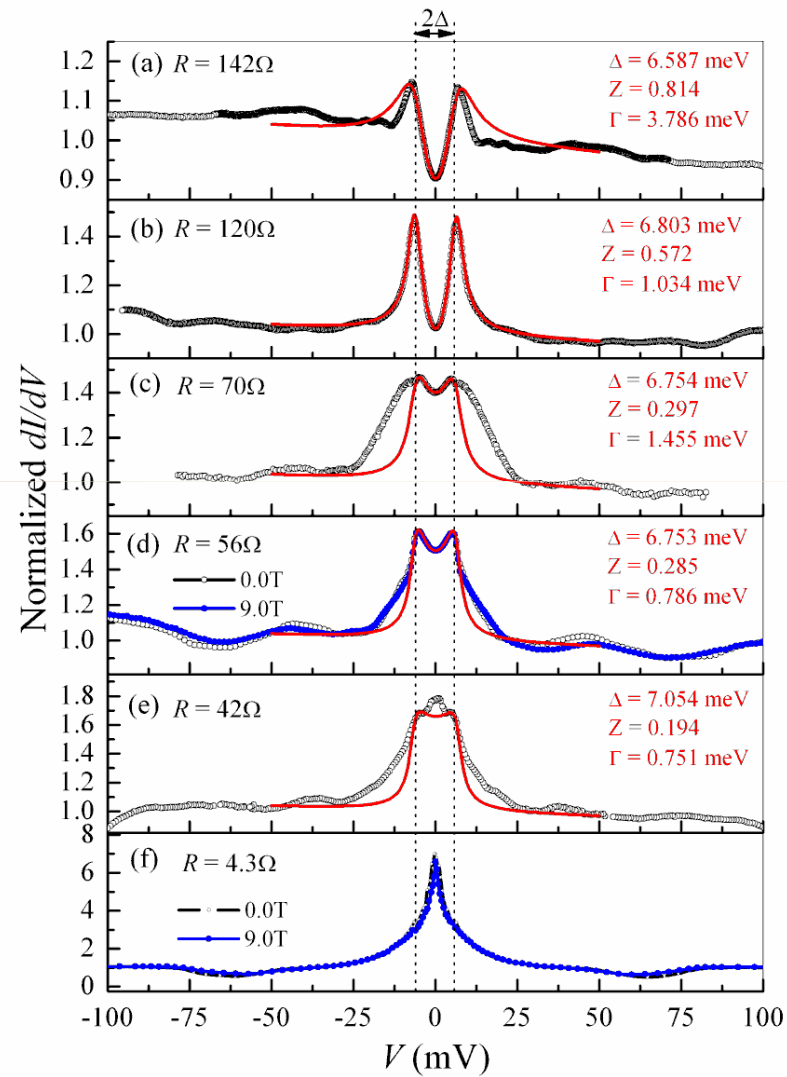
Tilt gun, co-sputtering

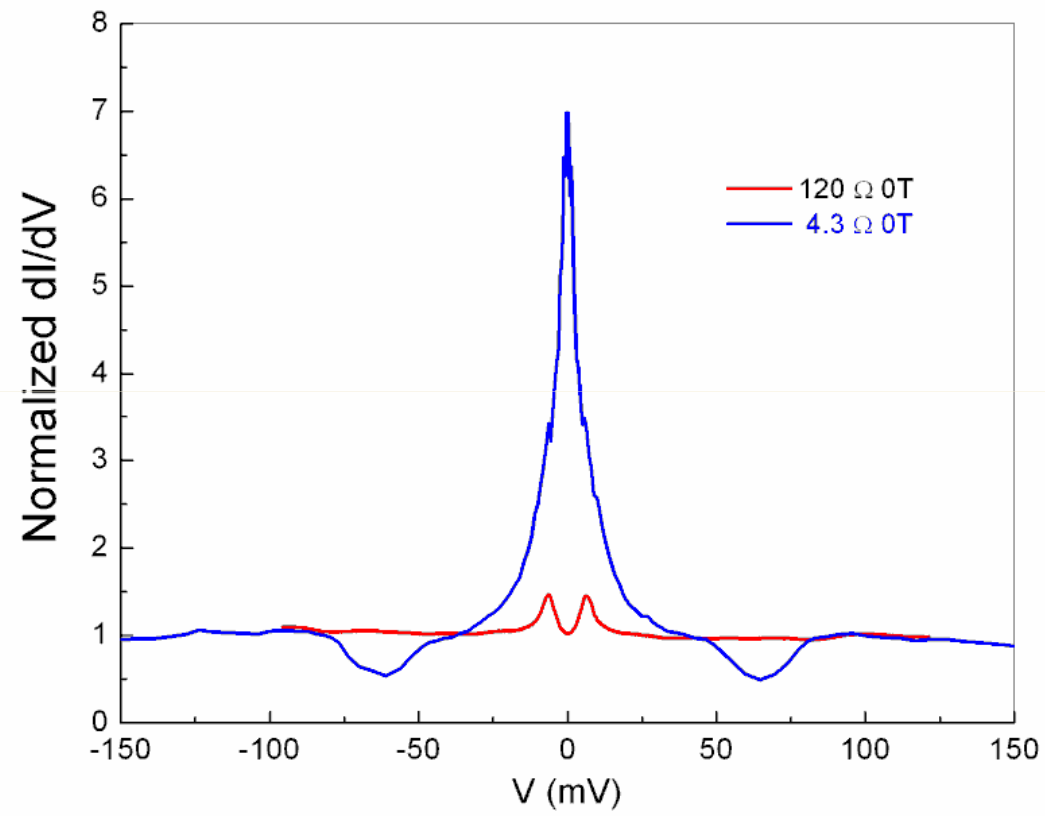
Evaporator

Vibrating Sample Magnetometer

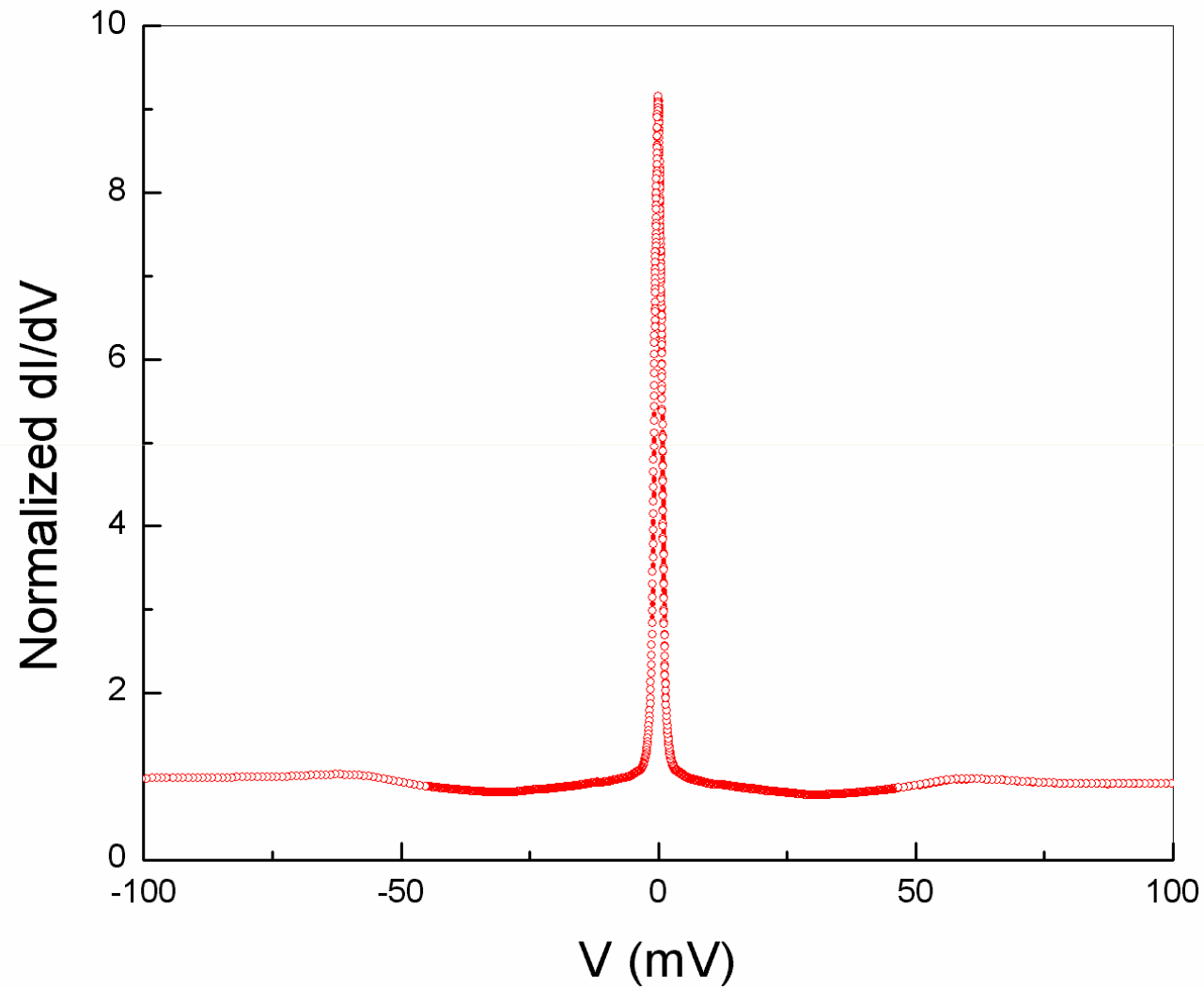
Thin film
Study oxides for energy

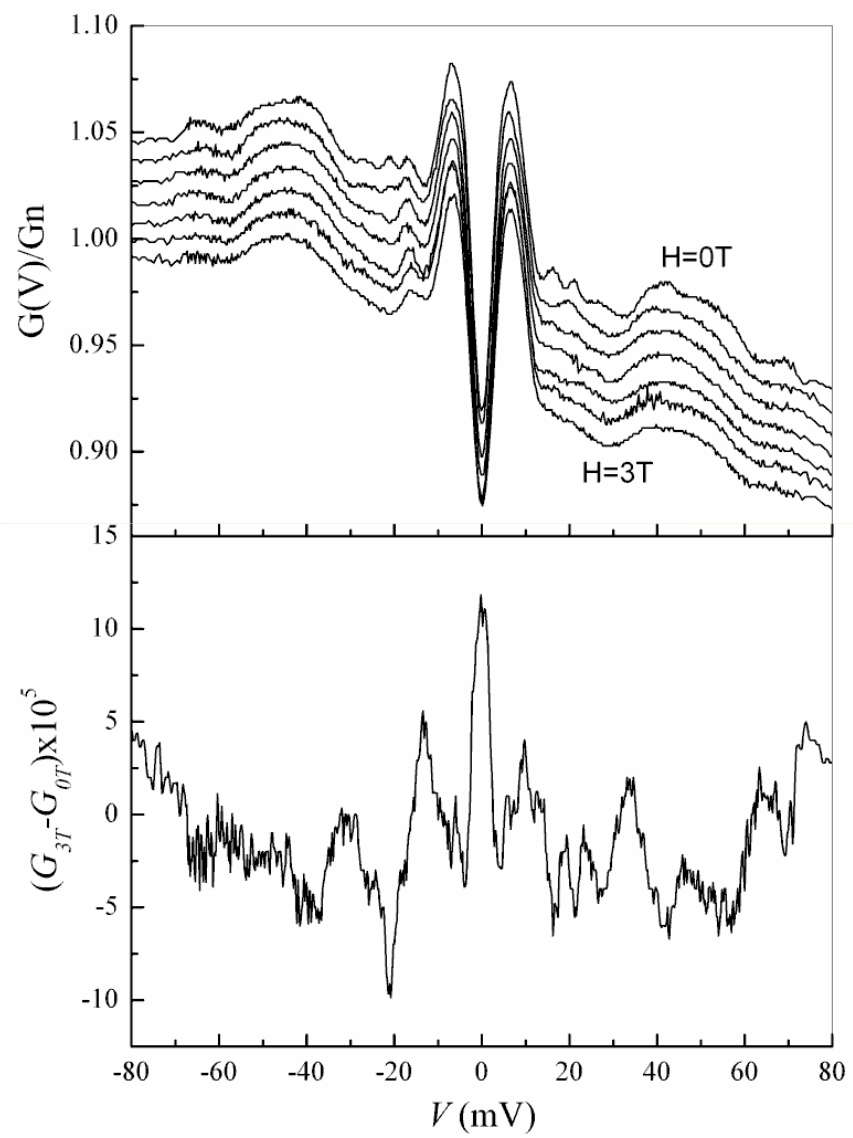
Dependence of Contact Resistance

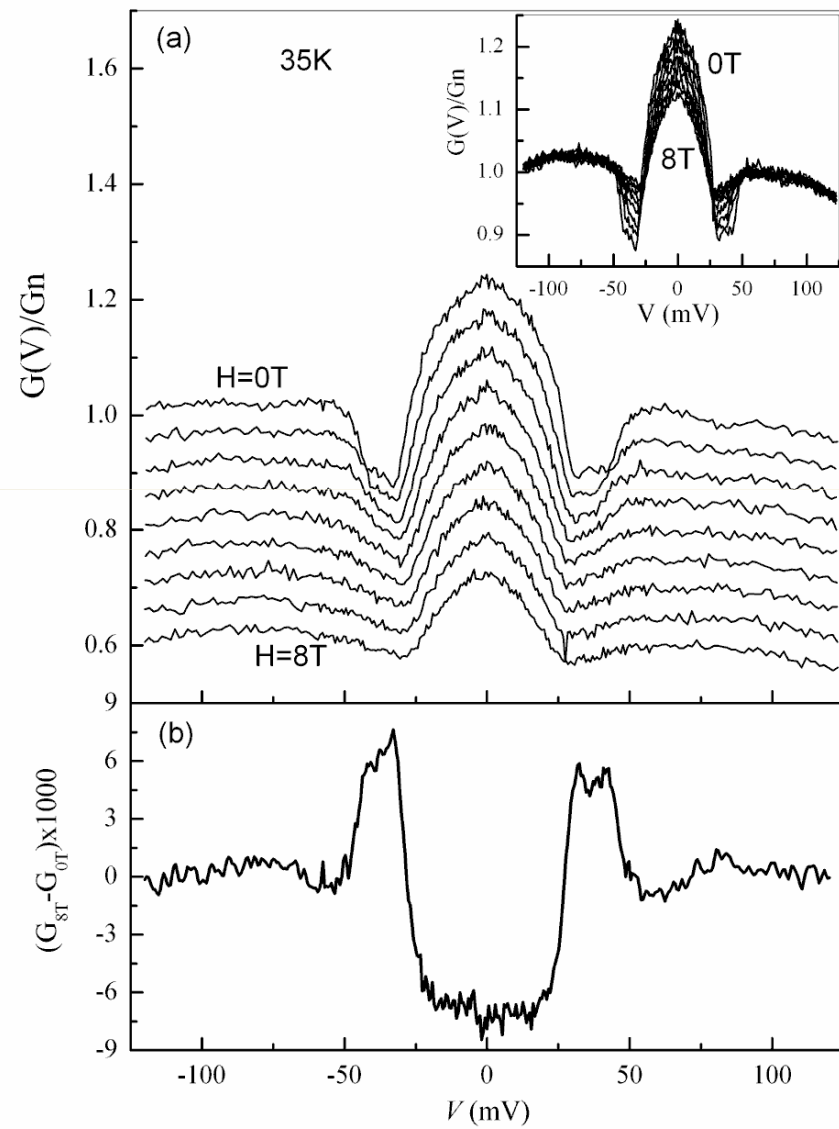


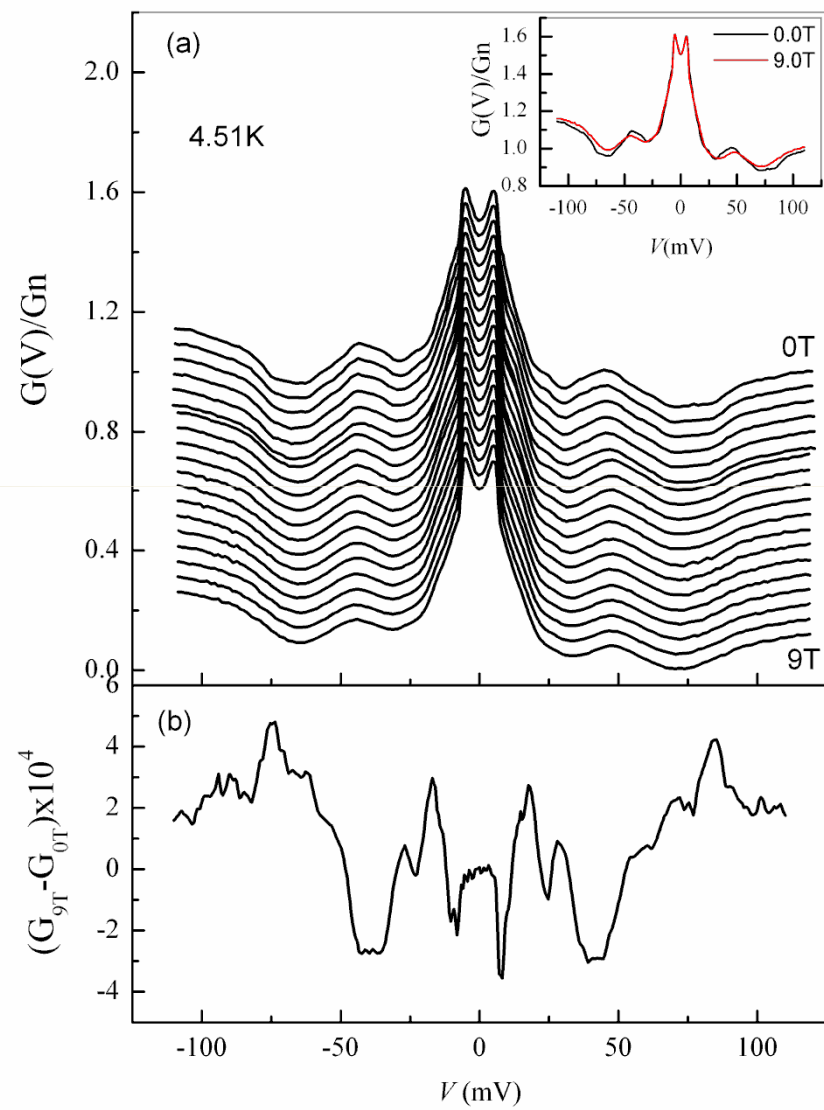


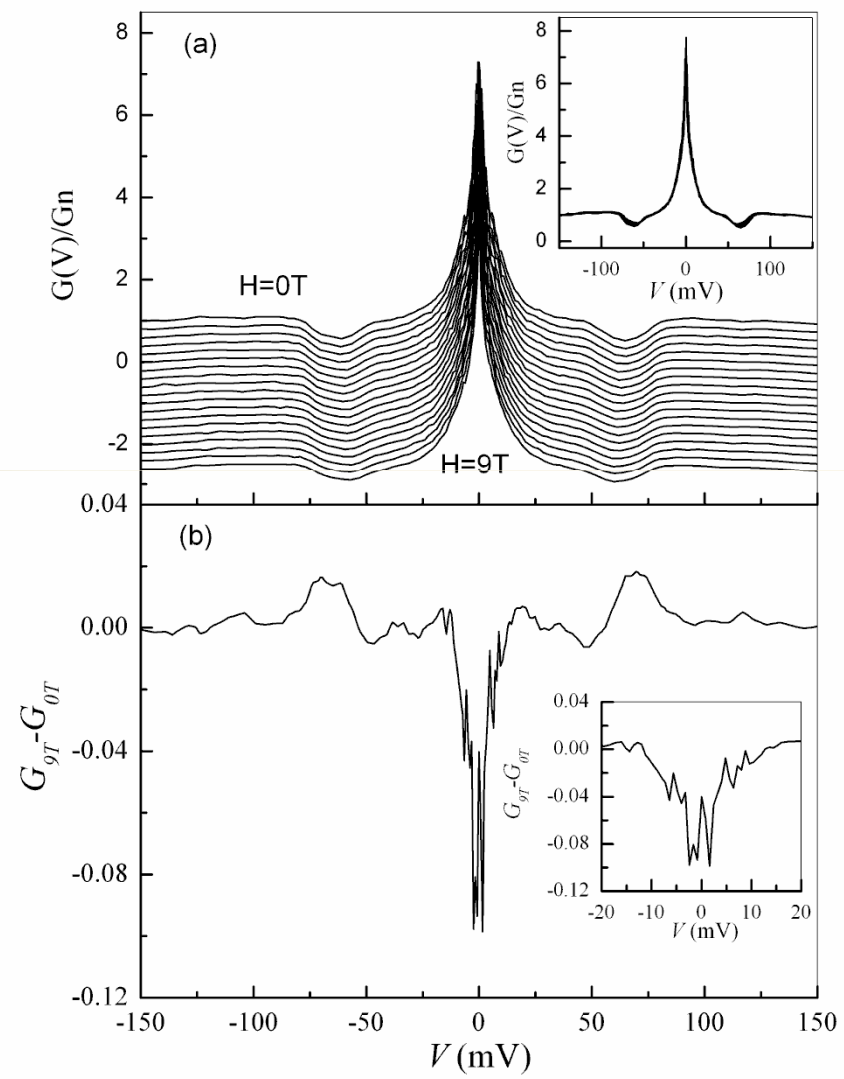
ZBA shape is different for different contacts

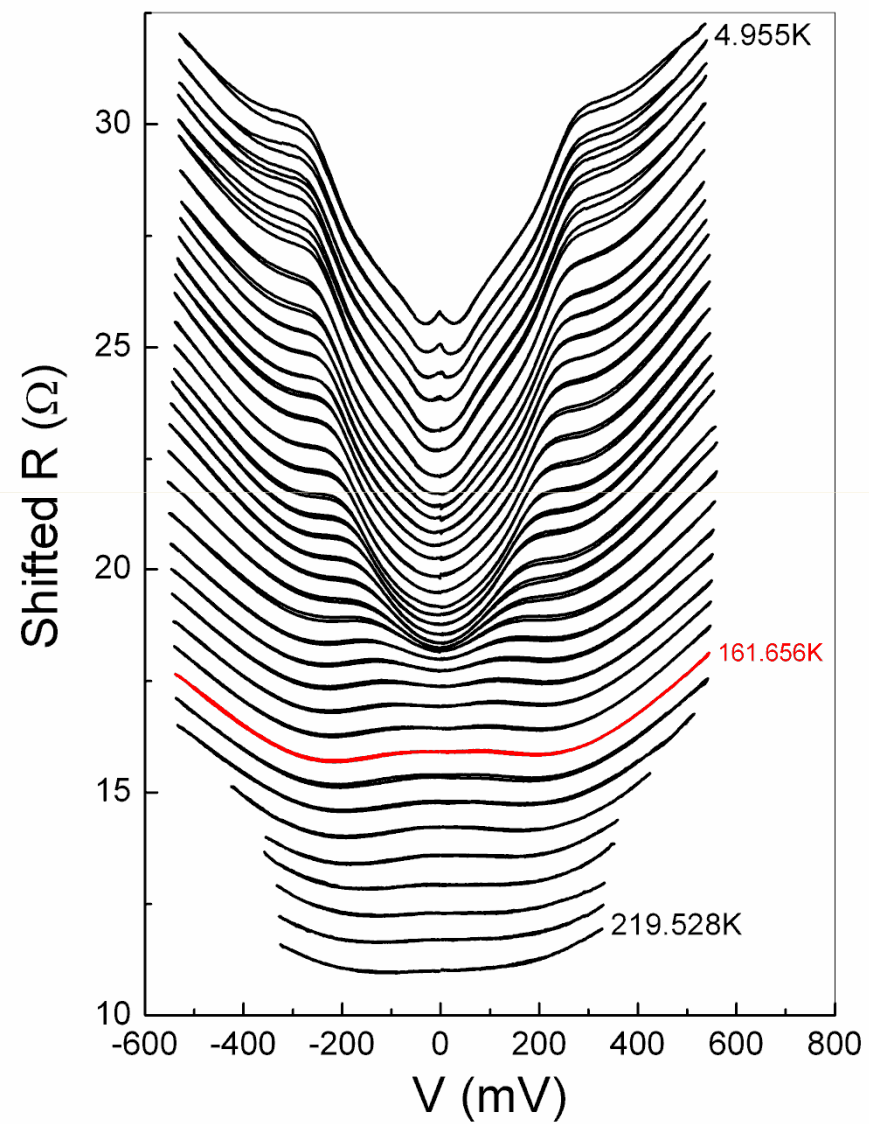




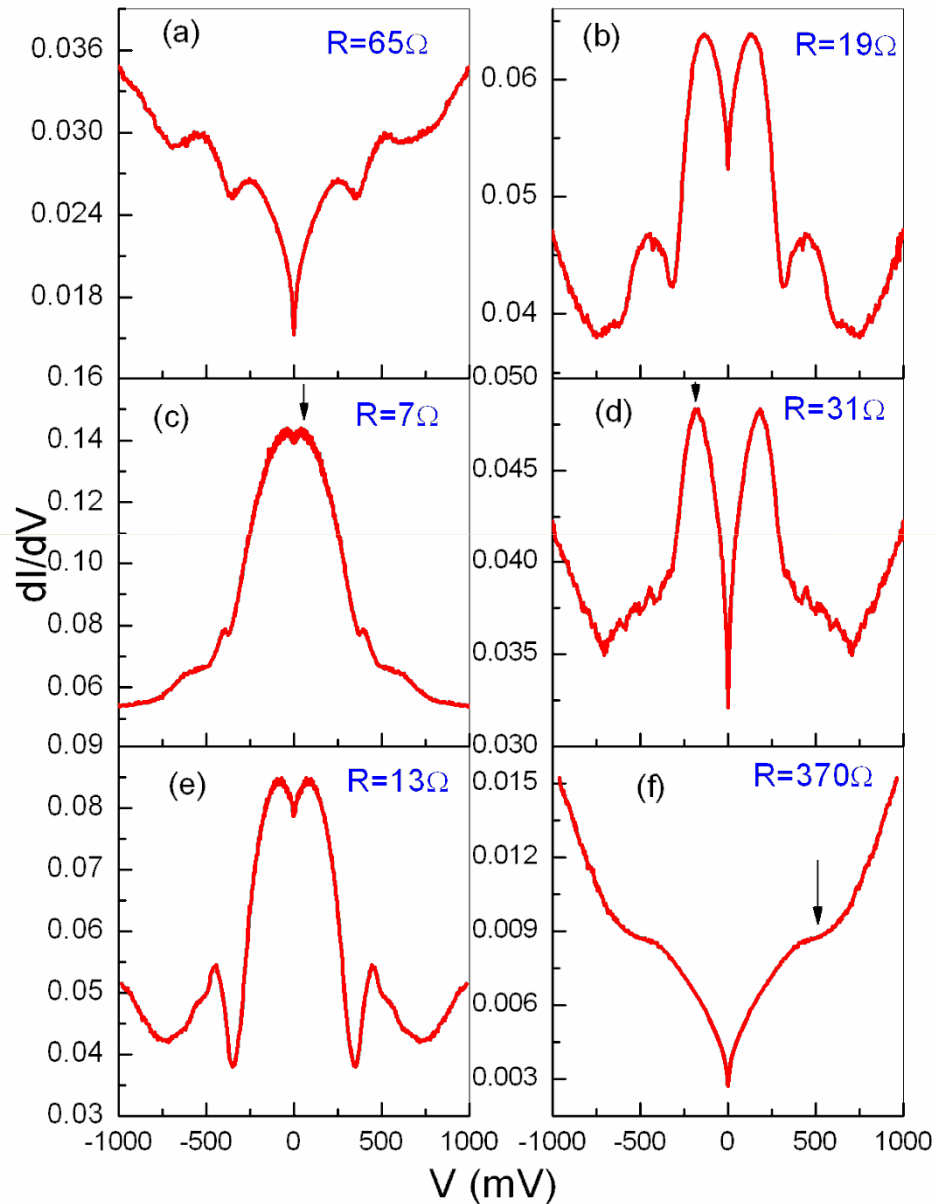








Various Contacts



Various contact, **various** peak values
from **5 mV** to **500 mV**

A real pseudogap must be **intrinsic**
independent of contact