

Measurement of the Polarization Observables I^s and I^c for $\vec{\gamma}p \rightarrow p \pi^+ \pi^-$ using the CLAS Spectrometer

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Outline

1 Theory and Motivation

- Theory: QCD
- Constituent Quark Models and Missing Resonances

2 Experiment

- Jefferson Lab
- Hall B

3 Data Analysis

- Polarization Observables
- Kinematic Fitting
- Extraction of Polarization Observables

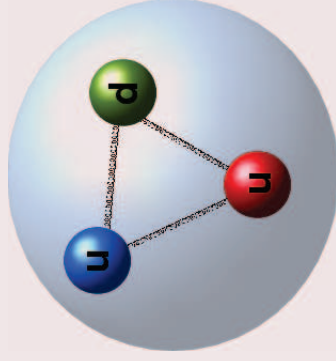
What is the big picture?

Hadrons:

- Particles that are held together and interact via the strong force.
- Exist as bound states of quarks.
- Are the most fundamental of systems found in nature.
- Members of this group are Baryons and Mesons.

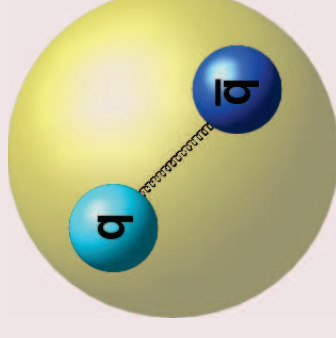
Quark Structure:

- Baryon: qqq



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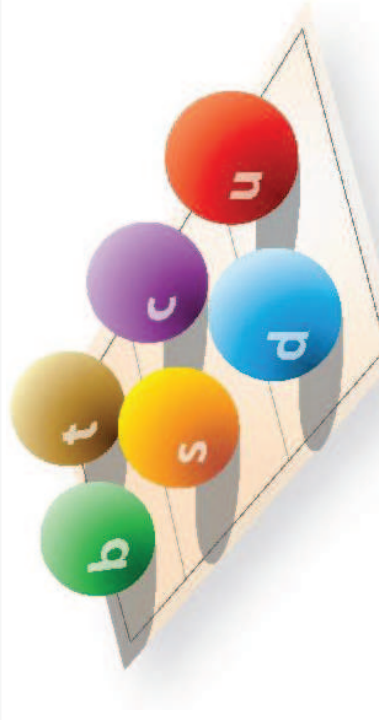
- Meson: $q\bar{q}$



Quantum Chromodynamics (QCD)

Properties of QCD

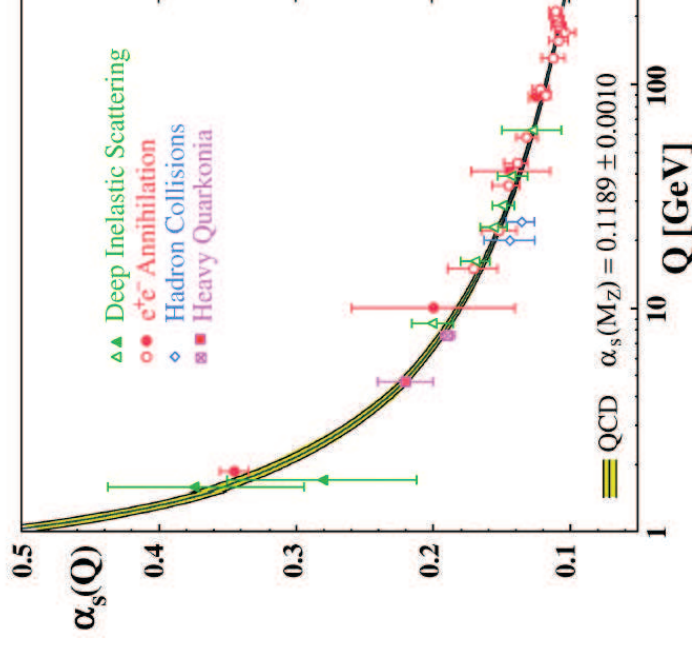
- Describes the strong force.
- Introduces a “color charge” analogous to the electric charge (**red, blue, green**).
- Gluons and quarks possess a color charge, antiparticles possess “anticolor”.
- All observed particles are “white” or colorless states.
- Interactions are mediated by gluons.
- Since gluons carry a color charge, they mediate AND participate in the interactions.



Quantum Chromodynamics (QCD): Features

Two main features of QCD

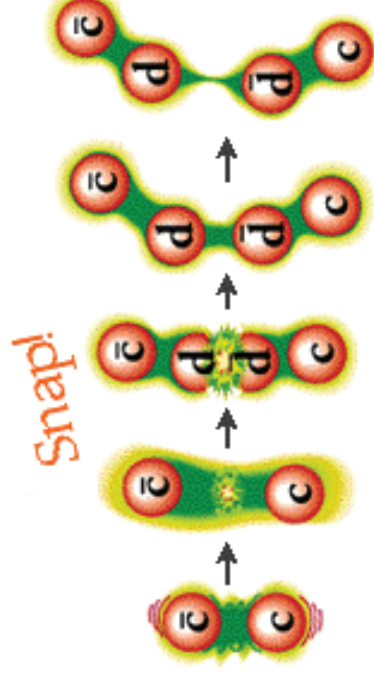
- Asymptotic Freedom: When the energy of the system is high (like in a high energy reaction $\rightarrow$ high momentum transfer), the quarks and gluons interact very weakly. Baryon can then be treated as being comprised of three non-interacting quarks $\rightarrow$ pQCD.



Quantum Chromodynamics (QCD): Features

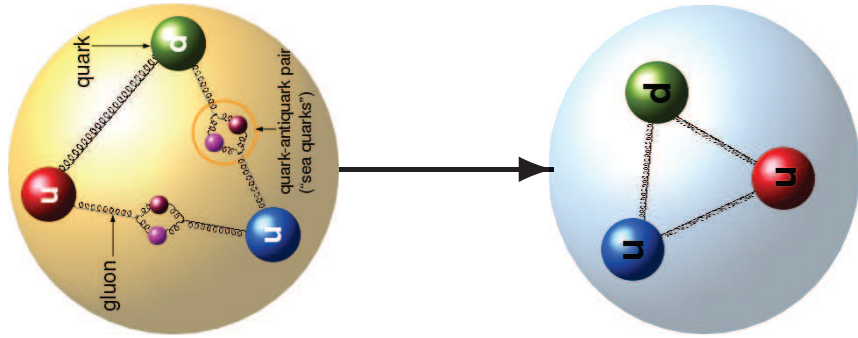
Two main features of QCD

- Asymptotic Freedom: When the energy of the system is high (like in a high energy reaction $\rightarrow$ high momentum transfer), the quarks and gluons interact very weakly. Baryon can then be treated as being comprised of three non-interacting quarks $\rightarrow$ pQCD.
- Confinement: The attractive force between quarks grows with the distance between them. This means that the quarks are confined to be in the hadron $\rightarrow$ there are no free quarks (non-pQCD).

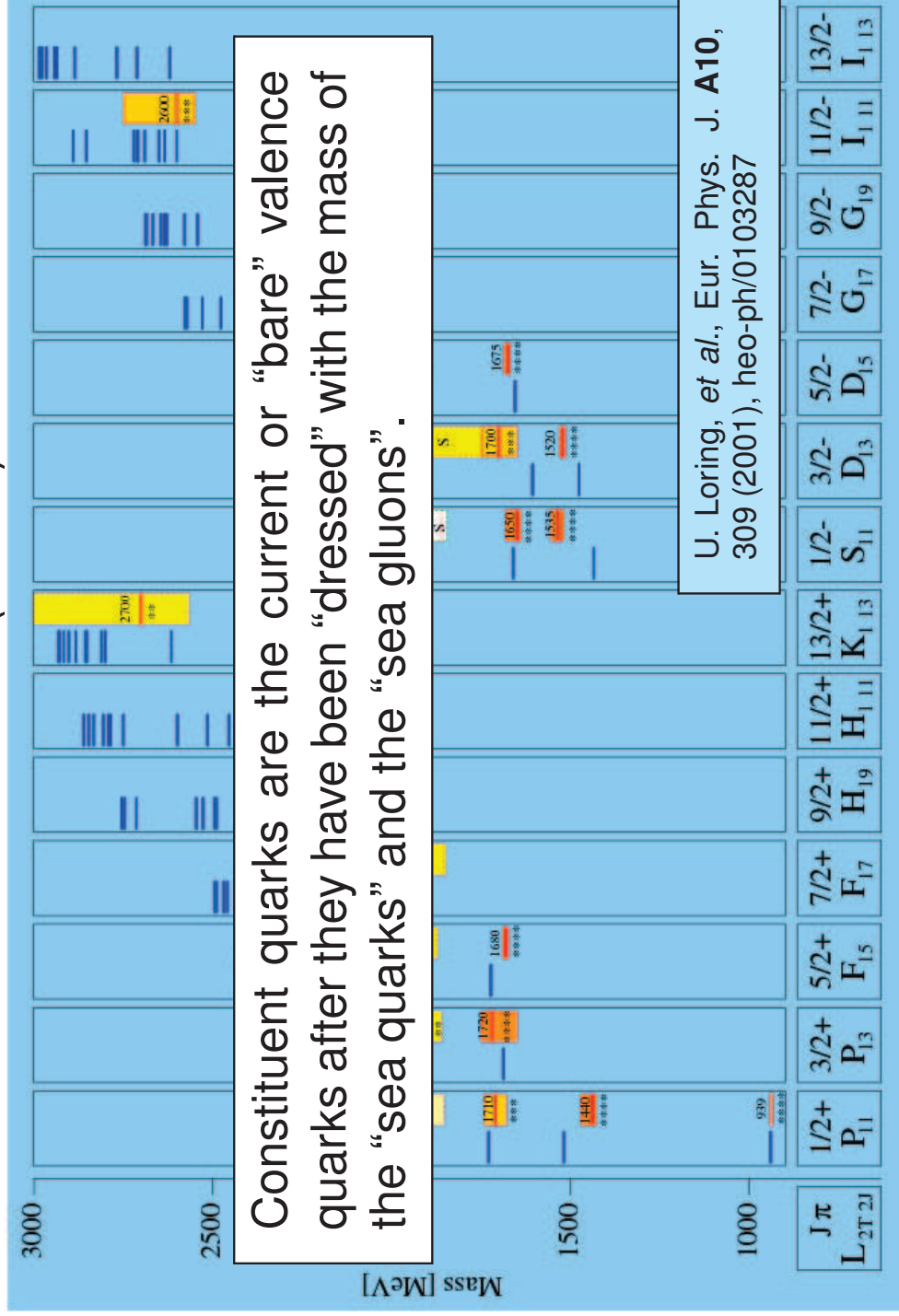


Constituent Quark Models (CQMs)

N* Resonances ($I = 1/2$)

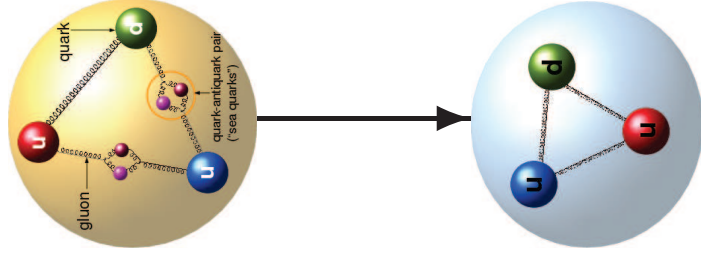


Constituent quarks are the current or “bare” valence quarks after they have been “dressed” with the mass of the “sea quarks” and the “sea gluons”.

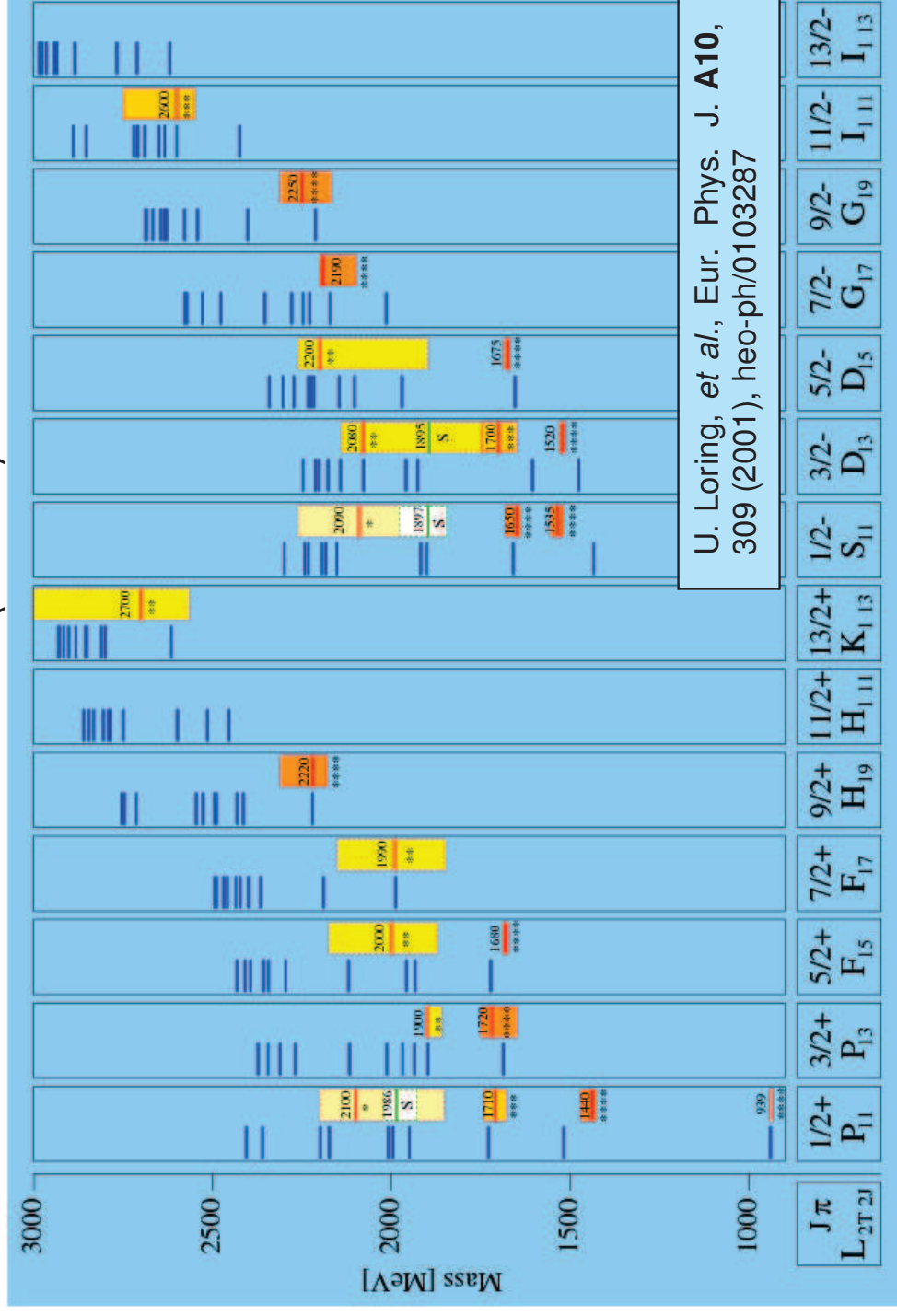


- ⇒ Generated using predictions according to a model view of the baryon.
- ⇒ Assumes three quark degrees of freedom.
- ⇒ Various models differ in how they treat short-range interactions.

Constituent Quark Models (CQMs)



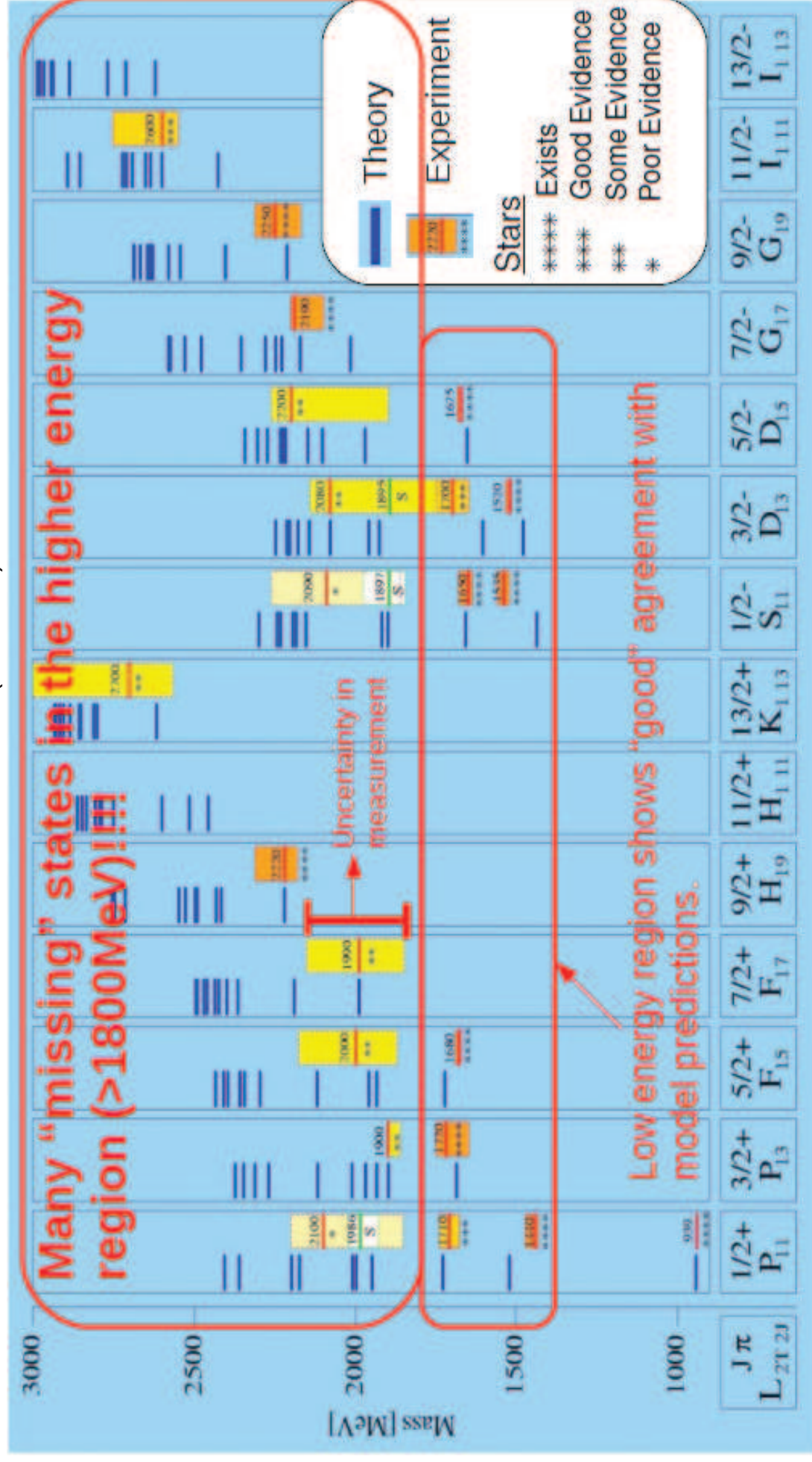
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CQMs: Missing Resonances

N* Resonances ($I = 1/2$)



Why so many missing resonances?

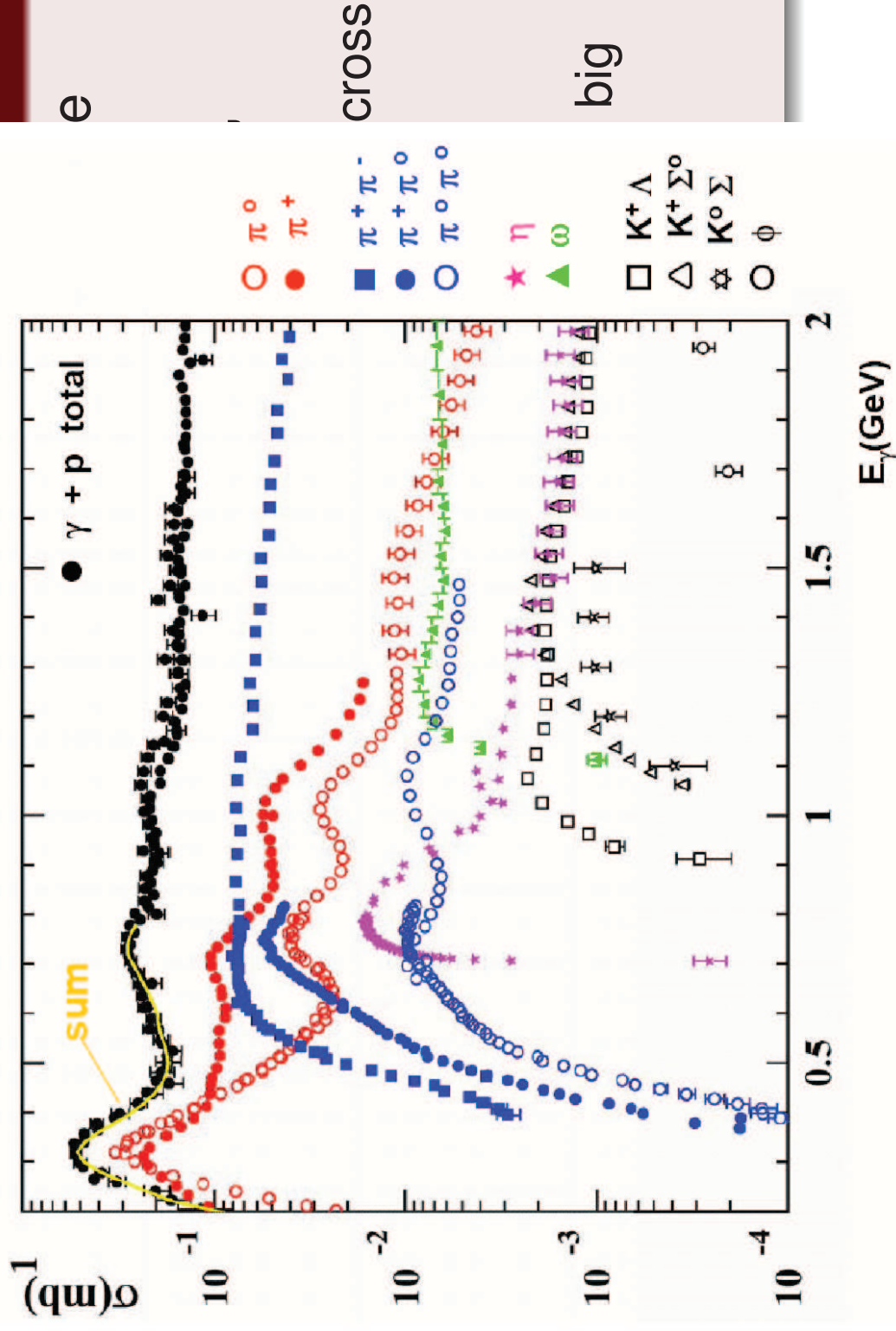
Possible answers and reasons...

- Most of what is found in the PDG handbooks has come from πN and $K N$ scattering $\rightarrow \gamma N$.
- Analyses involved a single meson final state (i.e. $K^+ \Lambda$, $K^+ \Sigma^0$, $N \pi$, $N \eta$).
 - Analyze a double-meson final state $\rightarrow$ has the largest cross section.
 - Sequential decay to final ground state particles:
ex. $\gamma p \rightarrow N^* \rightarrow \Delta^{++} \pi^- \rightarrow p \pi^+ \pi^-$.
 - Particle mass widths expected to be ≈ 150 MeV $\rightarrow$ too big for a two-body (single meson) final state.
- Possible quark-diquark structure of the baryon.

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Possible answers and reasons

- Mos from
- Anal $K^+\Sigma$
- Poss



Why so many missing resonances?

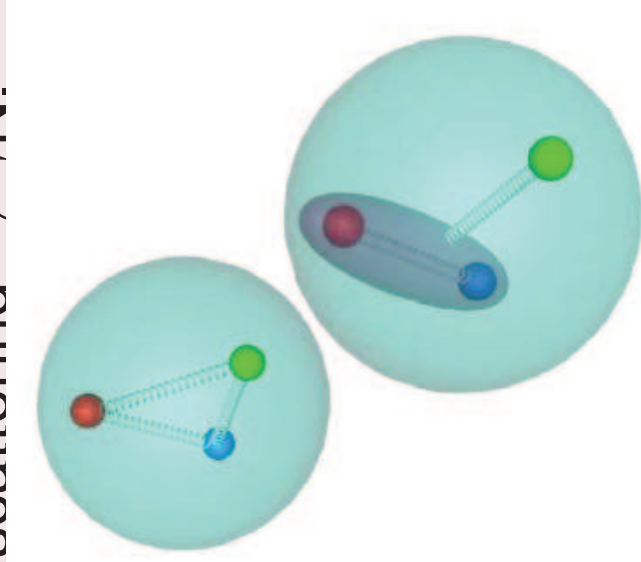
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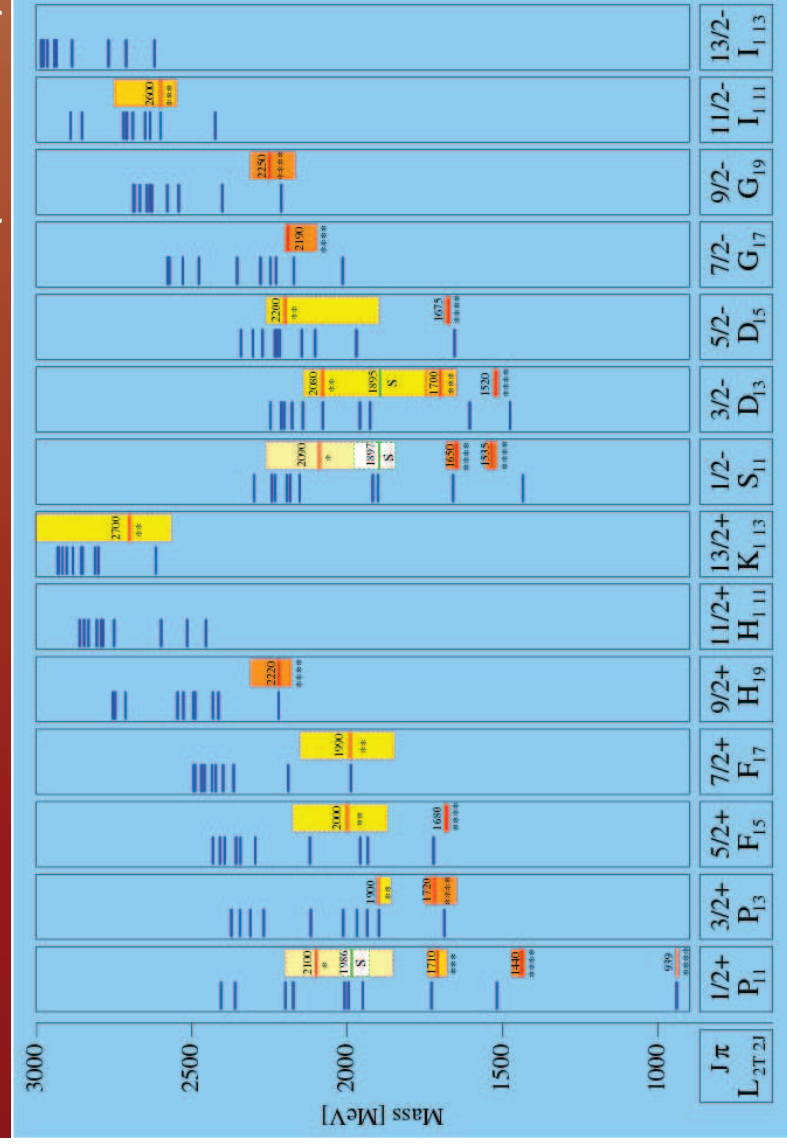
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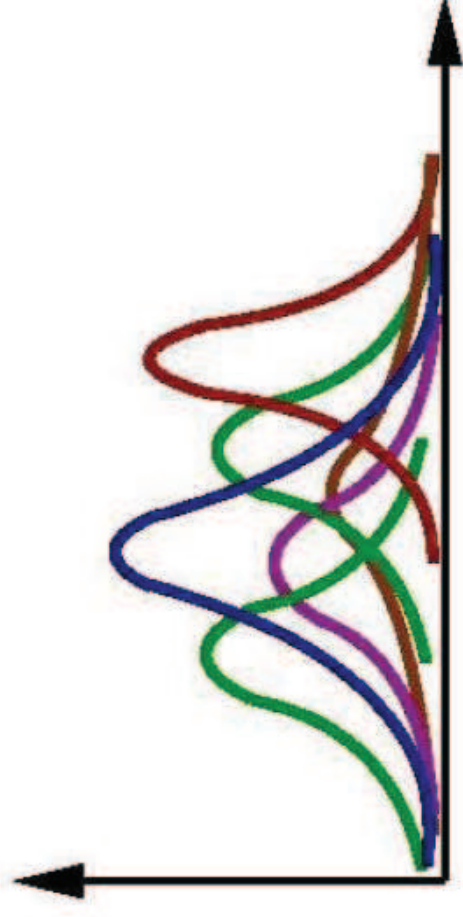
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 - Analyze a dc section.
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 - Particle mass for a two-body
- Possible quark-quark structure of the baryon.



Constituent Quark Models (CQMs)



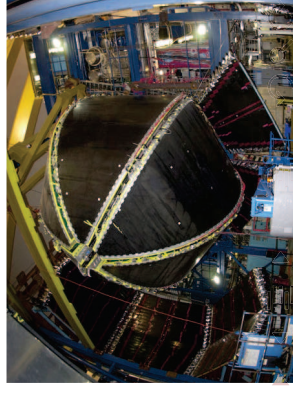
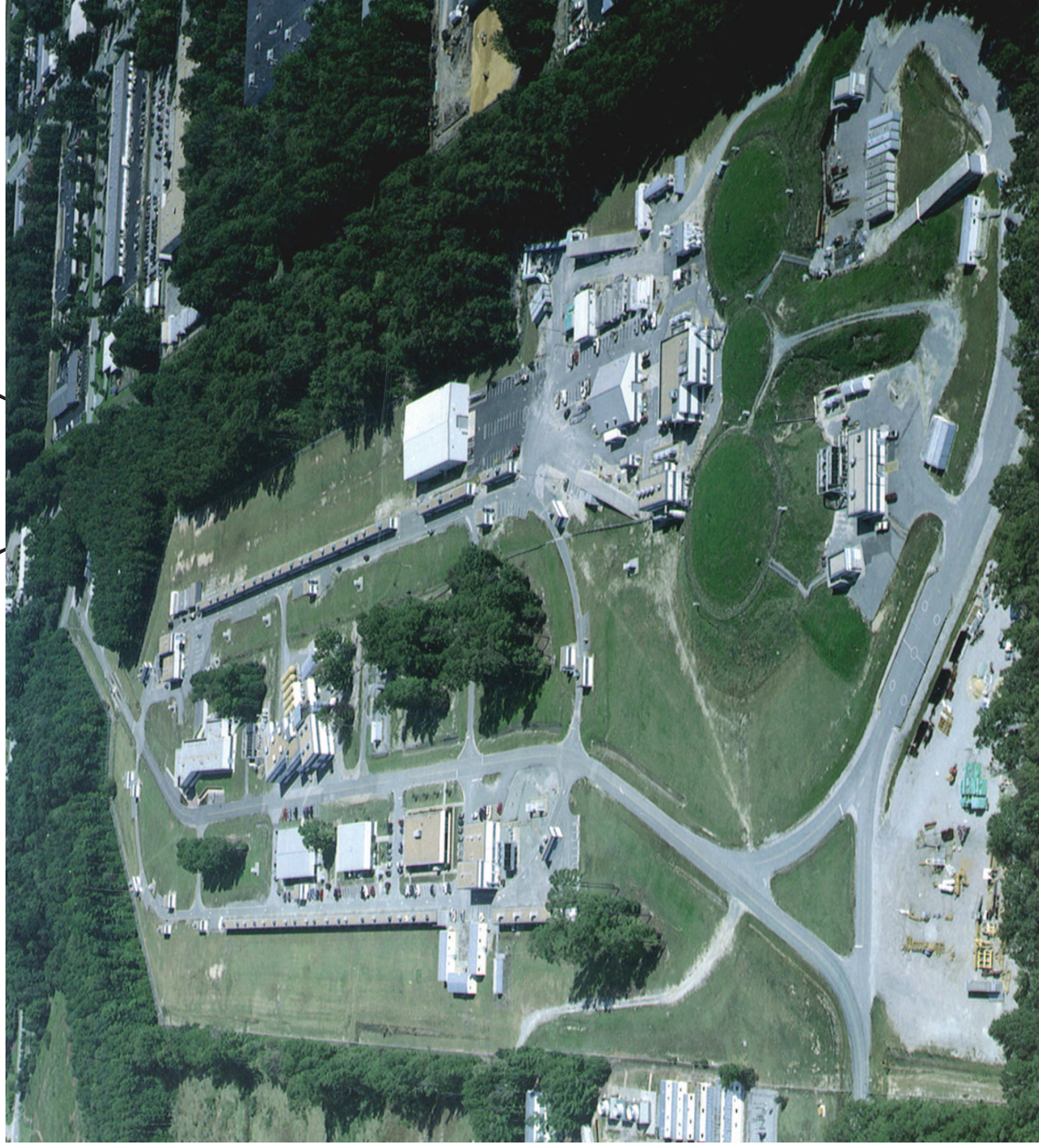
We need a way to disentangle these resonances and determine which resonances are produced!



- Resonances are both broad and overlapping!
- Makes it very difficult to determine singular resonance contributions.

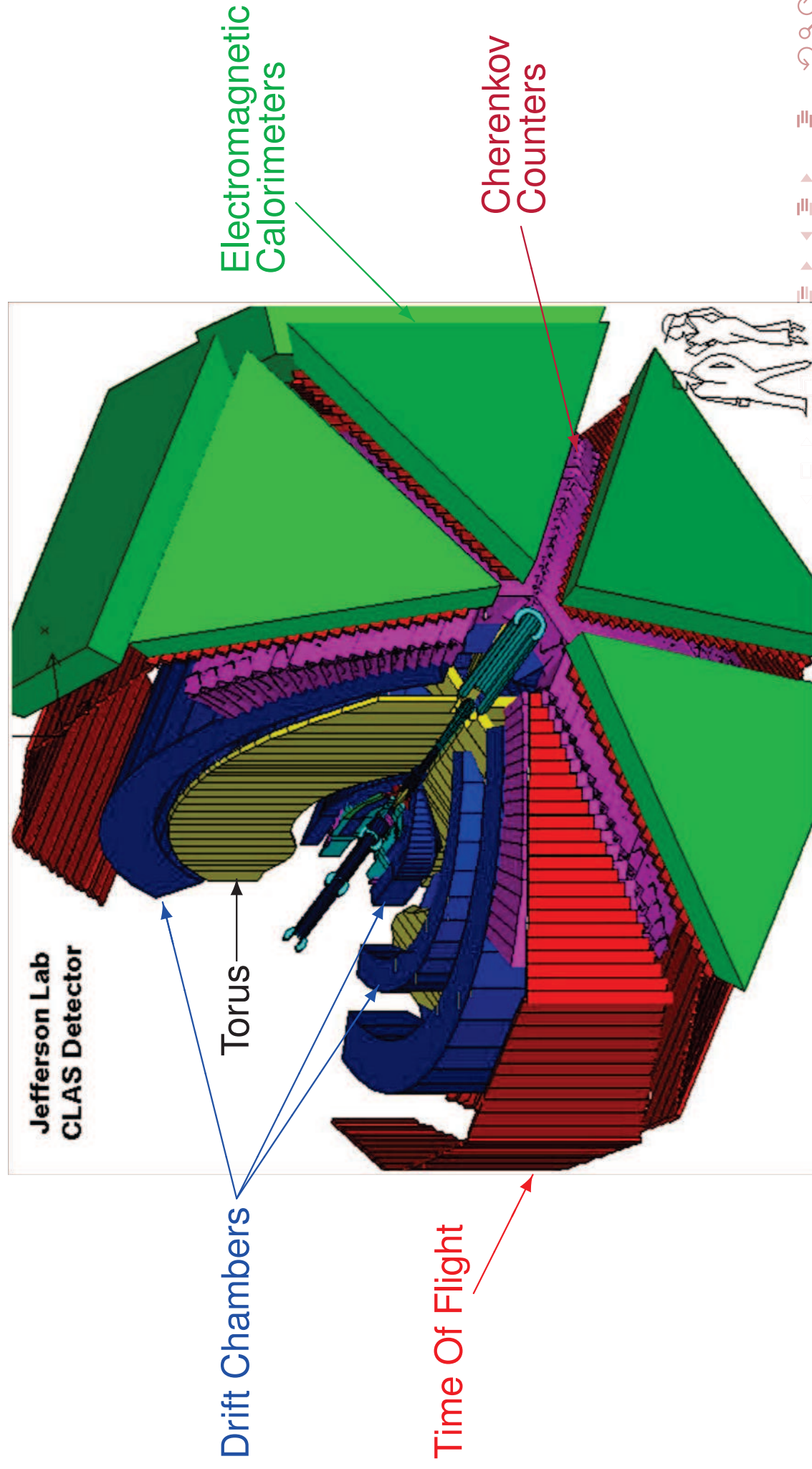
The Facility

THOMAS JEFFERSON NATIONAL ACCELERATOR FACILITY (JLAB)



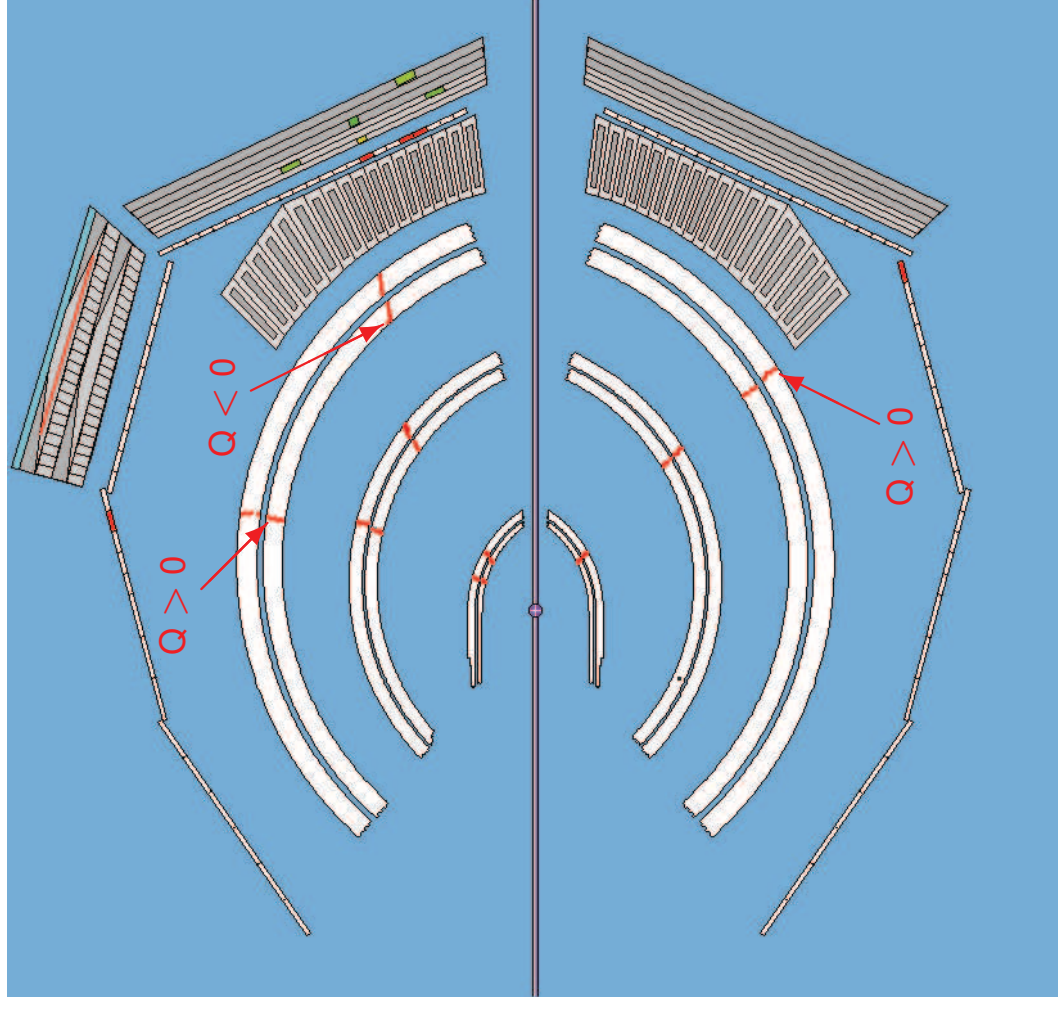
Hall B: The CLAS Detector

CEBAF Large Acceptance Spectrometer



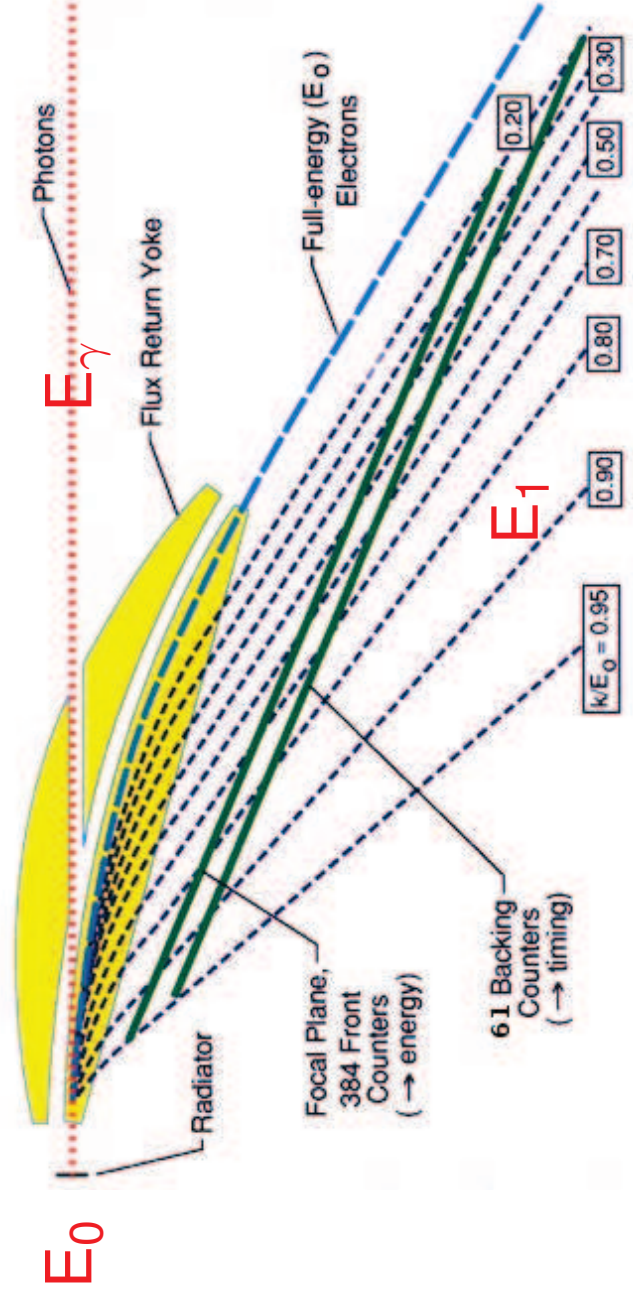
Hall B: the CLAS Detector

- Magnetic field of Torus affects the particle's trajectory (bending it toward or away from the beamline).
- Particle's path through CLAS is tracked.
- If trigger requirements are met, event is recorded.
- Reconstruction code determines particle's trajectory through CLAS.
- Information regarding the incident particle and final state particles are determined.



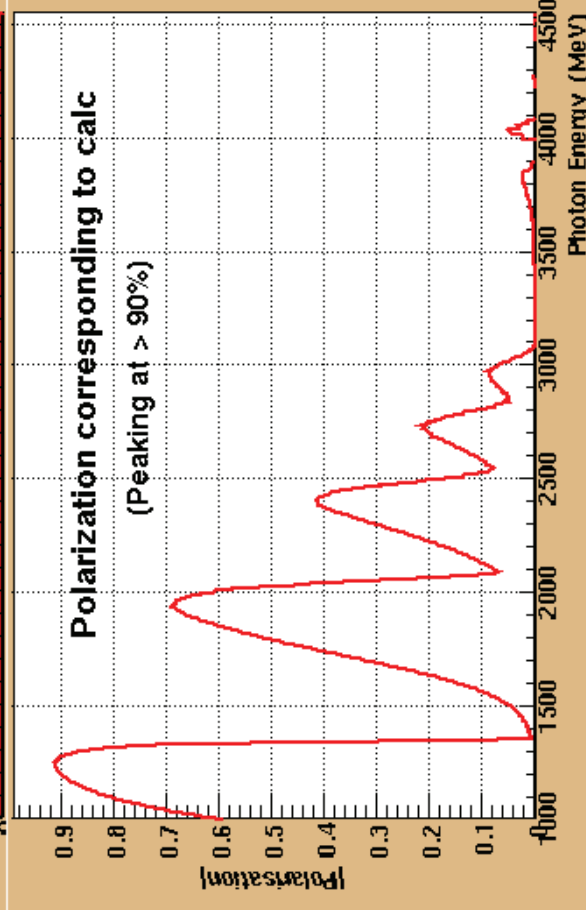
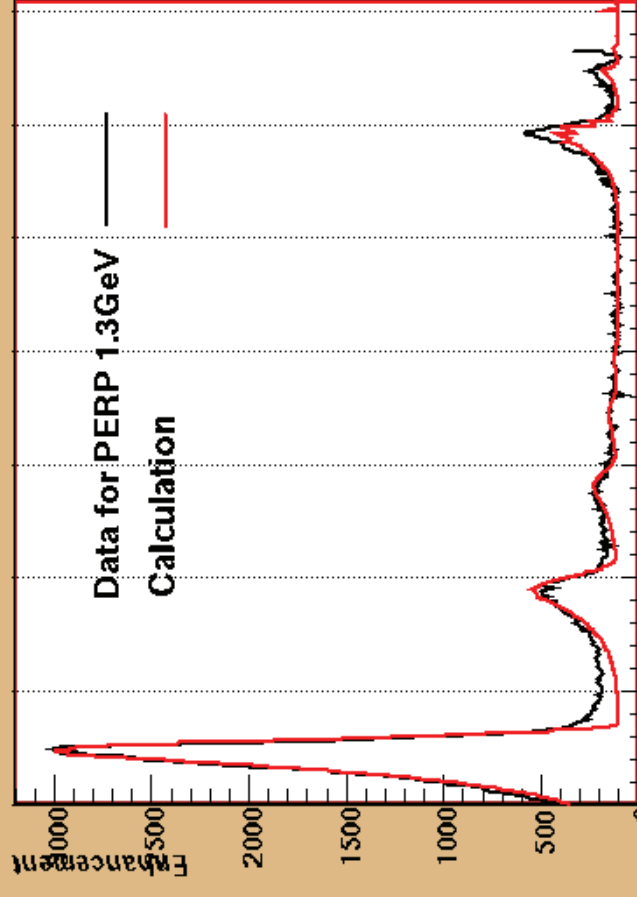
Hall B: Photon Tagger

- Electrons from accelerator interact with the radiator, producing photons via bremsstrahlung.
- Electrons are bent into hodoscope via the Tagger Magnet.
- Time stamp and energy of electron is measured.
- Photon energy determined via $E_\gamma = E_0 - E_1$
- Can tag photons with energies ranging from $(0.2)E_0$ to $(0.95)E_0$.

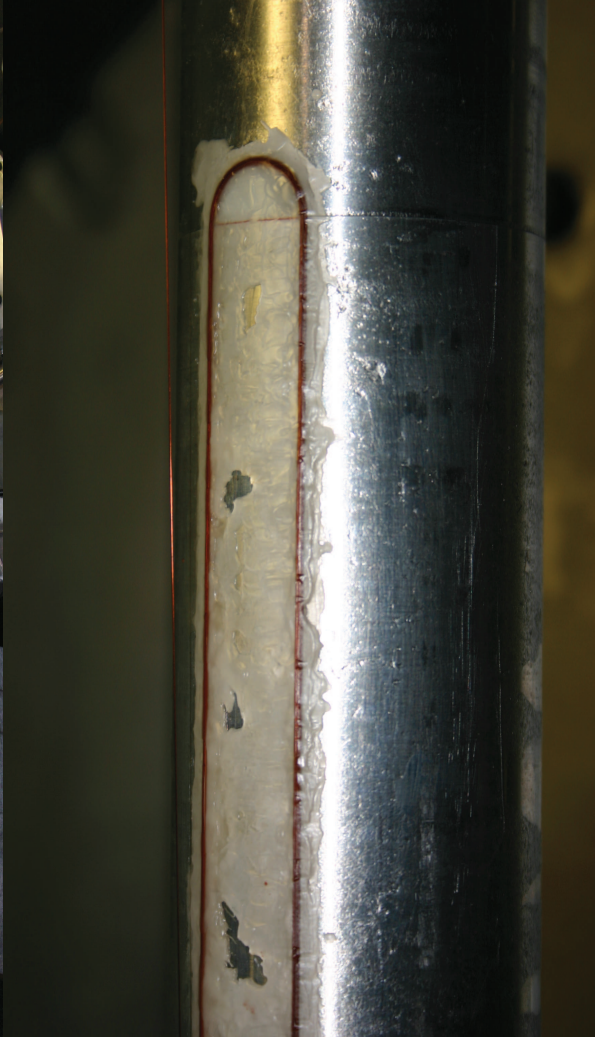
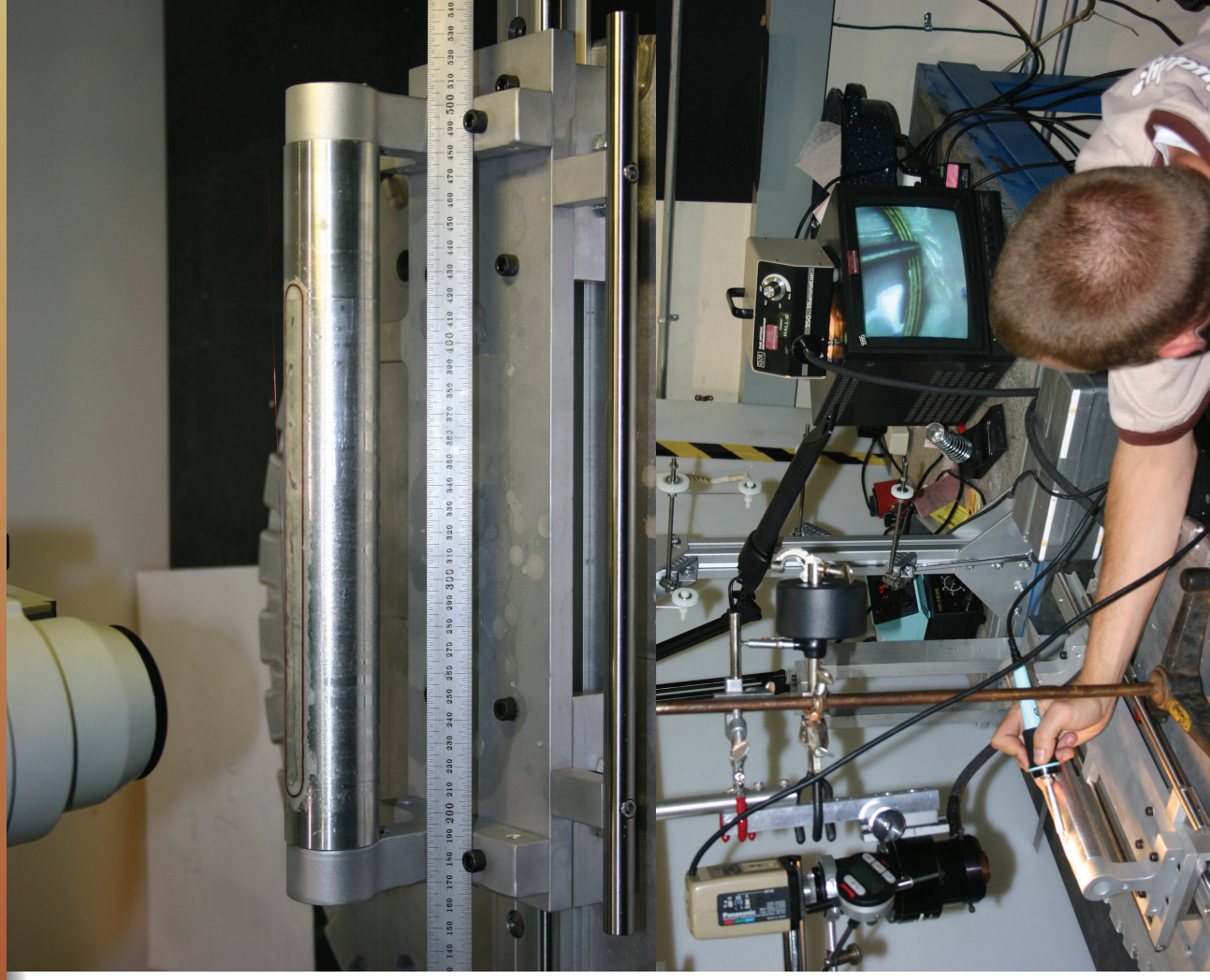


Polarized Photons

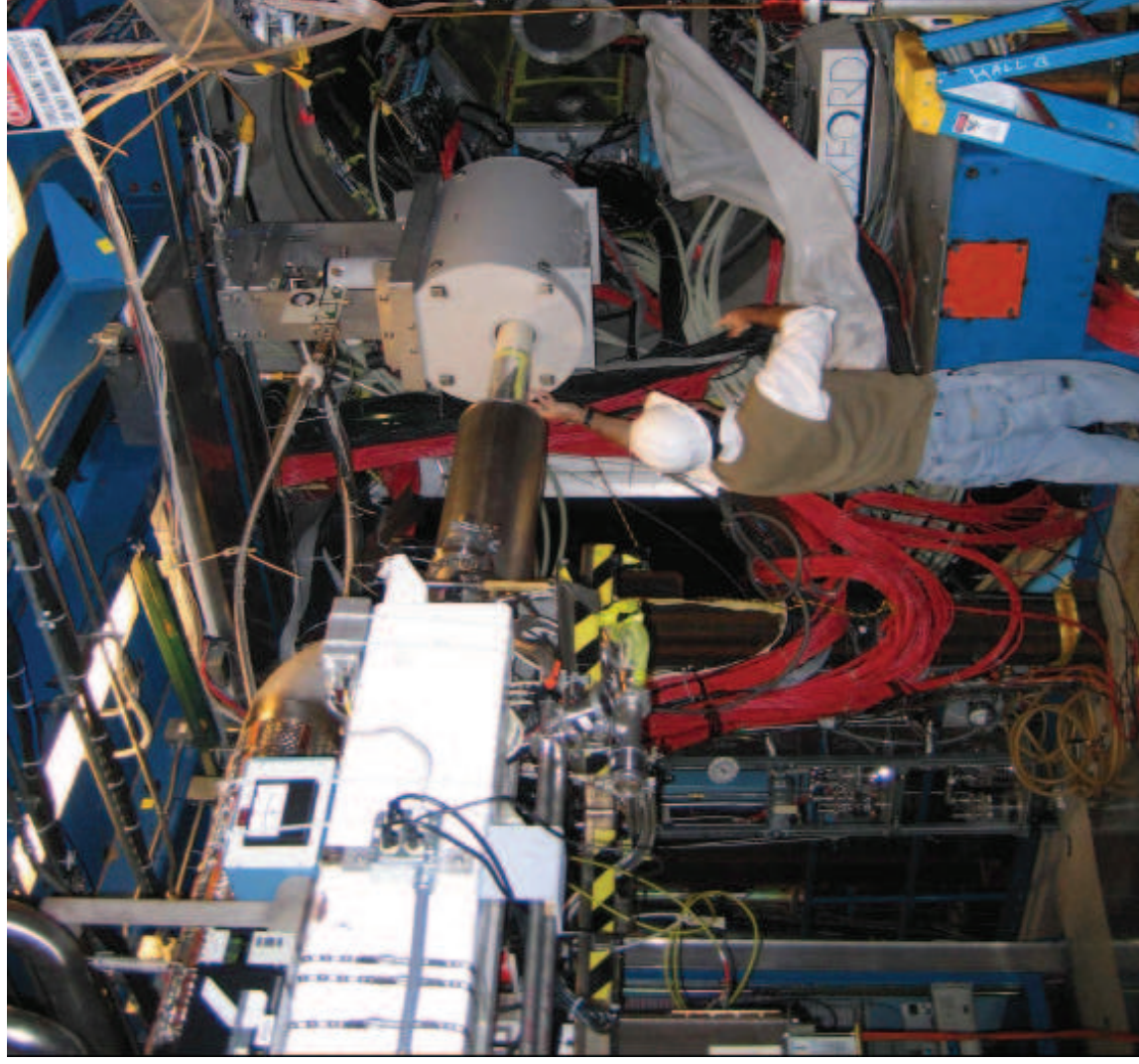
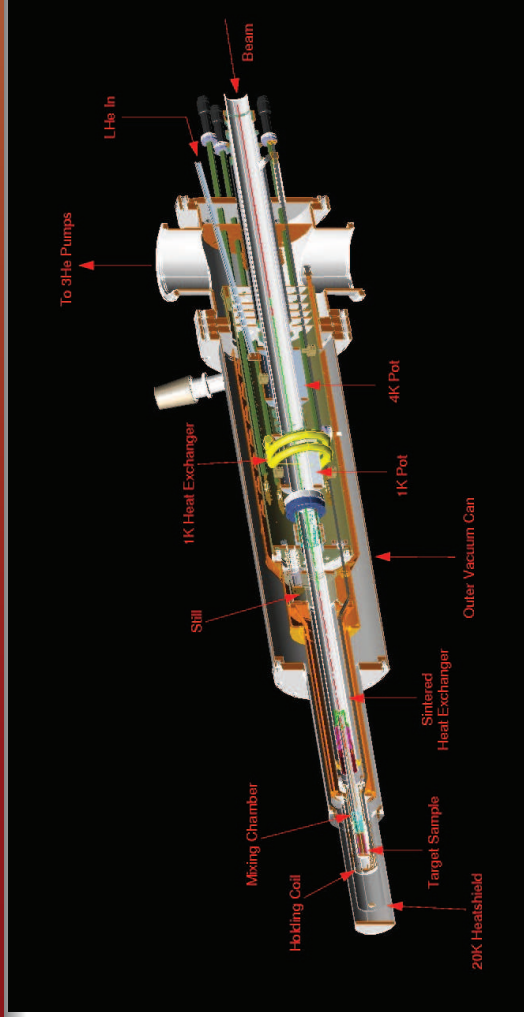
- Circularly polarized beam:
 - longitudinal electron beam + amorphous radiator
- Linearly polarized beam:
 - unpolarized electron beam + oriented diamond radiator
 - Can obtain 90% polarization
- g8b $\rightarrow$ highly polarized photons at five energies: 1.3 GeV, 1.5 GeV, 1.7 GeV, 1.9 GeV and 2.1 GeV.



FROST: Transverse Holding Coil



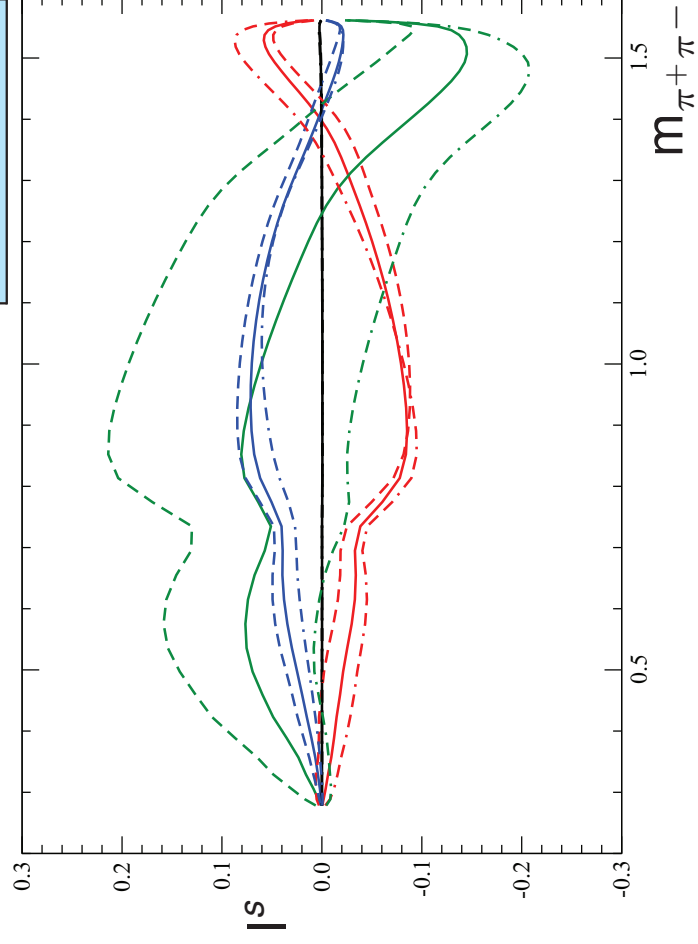
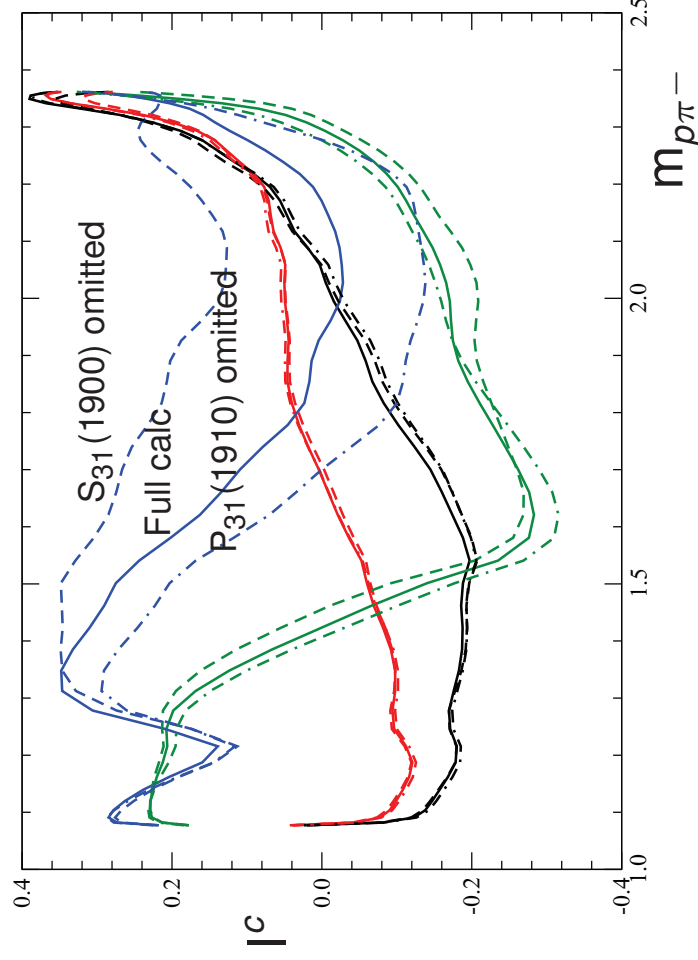
FROST: Target Assembly, Testing, Installation



Polarization Observables

- Measurable quantities called Polarization Observables are sensitive to resonance contributions.
- Occur when the constraint of polarization (beam and/or target) is imposed on the reaction(s).
- Can be predicted according to model calculations.

Dr. W. Roberts



$$\phi^* = 0.00356138 \text{ rad}$$

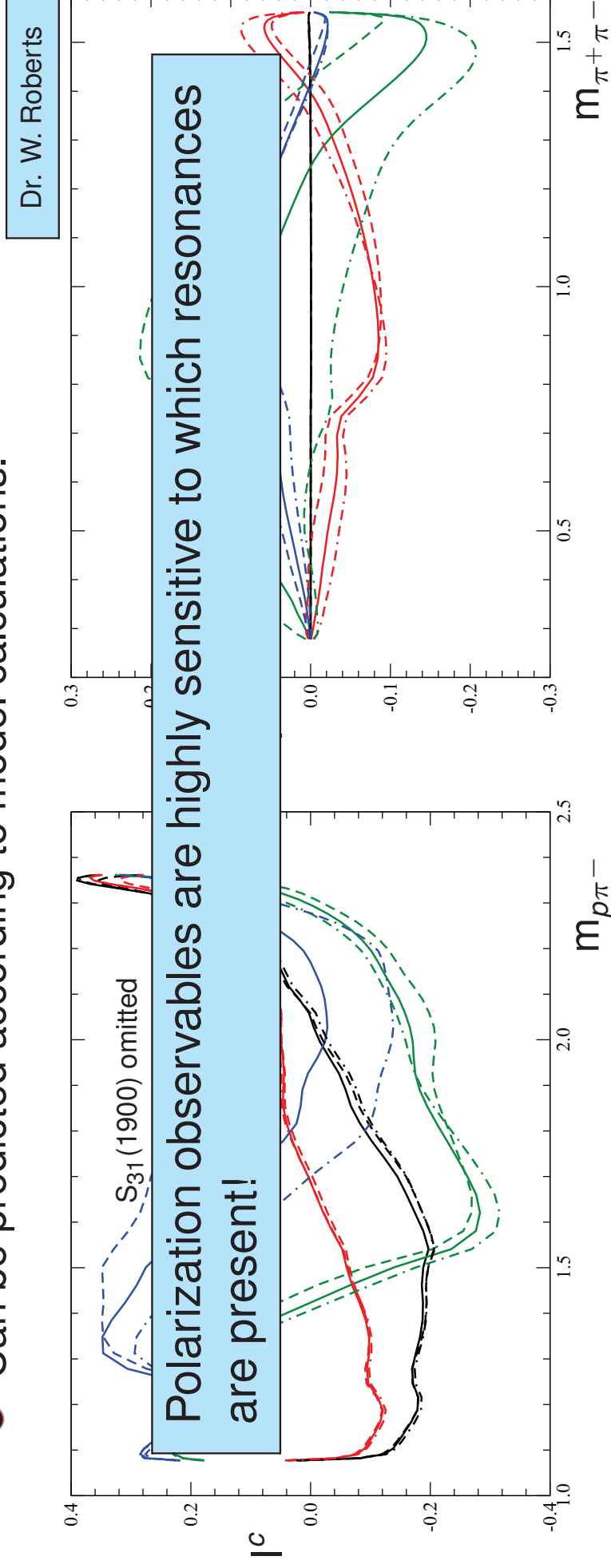
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$$\phi^* = 2.0876 \text{ rad}$$

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Polarization Observables: Single Meson Final State

- The final state equation relates the differential cross section, $\frac{d\sigma}{d\Omega}$, to the polarization observables.
- For a final state with one meson, there are a total of 7 observables.

$$\begin{aligned} \frac{d\sigma}{d\Omega} = \sigma_0 \{ & 1 - \delta_I \Sigma \cos(2\phi) \\ & + \Lambda_x (-\delta_I \mathbf{H} \sin(2\phi) + \delta_{\odot} \mathbf{F}) \\ & - \Lambda_y (-\mathbf{T} + \delta_I \mathbf{P} \cos(2\phi)) \\ & - \Lambda_z (-\delta_I \mathbf{G} \sin(2\phi) + \delta_{\odot} \mathbf{E}) \} \end{aligned}$$

Definitions

- $\frac{d\sigma}{d\Omega}$ = differential cross section
- σ_0 = unpolarized cross section
- $\delta_{I,\odot}$ = degree and orientation of photon beam polarization
- $\Lambda_{x,y,z}$ = direction and degree of polarization of the target
- $\Sigma, \mathbf{H}, \mathbf{F}, \mathbf{T}, \mathbf{P}, \mathbf{G}, \mathbf{E}$ = polarization observables.



Polarization Observables: Double Meson Final State

- For a final state with two mesons, there are a total of 15 observables!

$$I = I_0 \{ (1 + \vec{\Lambda}_i \cdot \vec{P}) + \delta_{I^\odot} (I^\odot + \vec{\Lambda}_i \cdot \vec{P}^\odot) \\ + \delta_I [\sin(2\beta) (I^s + \vec{\Lambda}_i \cdot \vec{P}^s) \\ + \cos(2\beta) (I^c + \vec{\Lambda}_i \cdot \vec{P}^c)] \}$$

Definitions

- I_0 = unpolarized reaction rate
- $\vec{\Lambda}_i$ = direction and degree of polarization of the target
- $\delta_{I^\odot}$ = degree and orientation of photon beam polarization
- $\vec{P}$ = observable arising from the use of a polarized target
- $I^{\odot,s,c}$ = observables arising from the use of polarized photons
- β = orientation of polarization w.r.t. a final state particle
($\beta = \phi_{lab} + \phi_{polarization}$)



Reducing the final state equation

- Applying the run conditions of g8b simplifies the final state equation and reduces the number of observables.
- Linearly polarized photon beam incident on an unpolarized LH₂ target. ($\vec{\Lambda}_i = 0$, $\delta_{\odot} = 0$)

$$\begin{aligned}
 I = I_0 \{ & (1 + \vec{\Lambda}_i \cdot \vec{\mathbf{P}}) + \delta_{\odot} (\mathbf{I}^{\odot} + \vec{\Lambda}_i \cdot \vec{\mathbf{P}}^{\odot}) \\
 & + \delta_I [\sin(2\beta) (\mathbf{I}^{\mathbf{s}} + \vec{\Lambda}_i \cdot \vec{\mathbf{P}}^{\mathbf{s}}) \\
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Reducing the final state equation

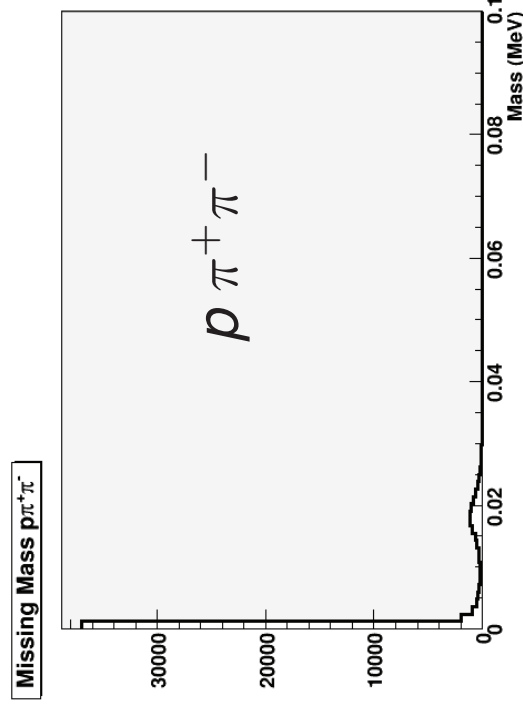
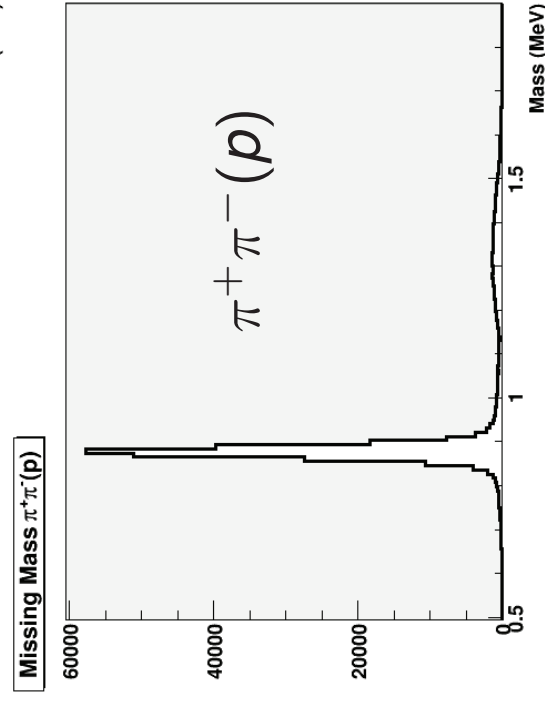
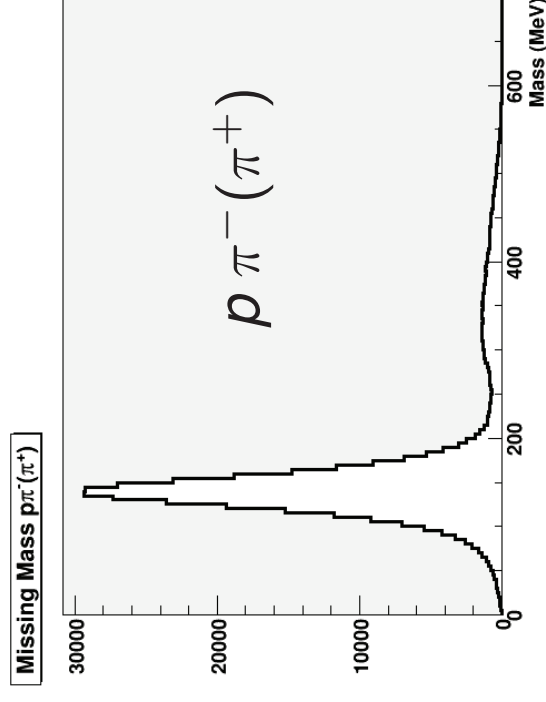
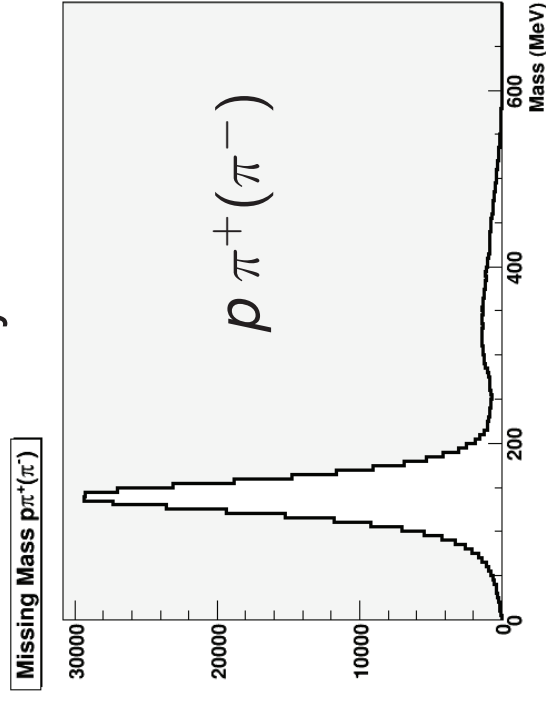
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$$I = I_0 \{ 1 + \delta_I [I^S \sin(2\beta) + I^C \cos(2\beta)] \}$$

- I^S
- I^C (Σ in the single-meson final state equation)
- Measuring both for $\vec{\gamma}p \rightarrow p \pi^+ \pi^-$

Analysis of $\vec{\gamma}p \rightarrow p \pi^+ \pi^-$

- While $\vec{\gamma}p \rightarrow p \pi^+ \pi^-$ is studied, a total of four topologies are kinematically fitted.



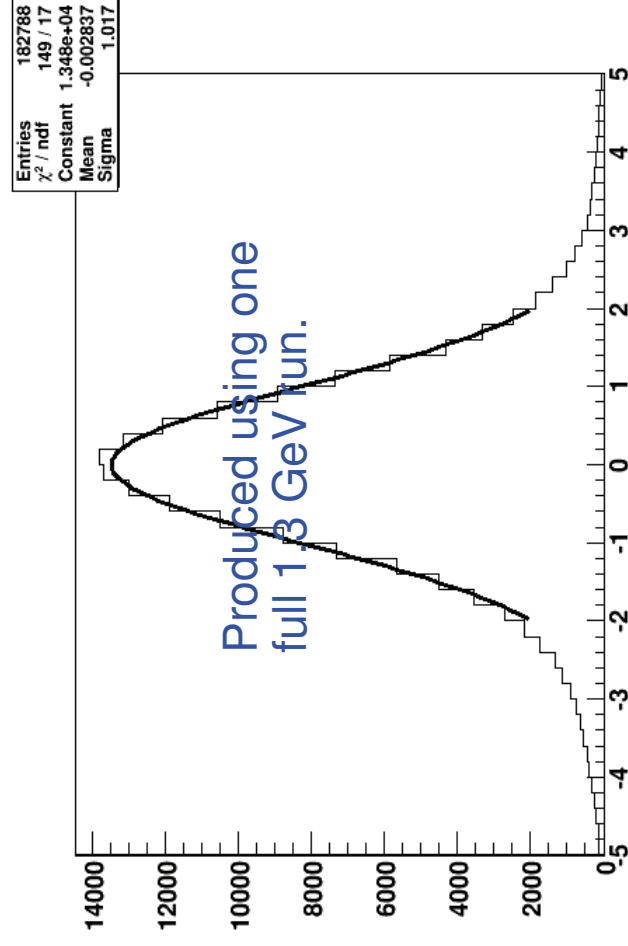
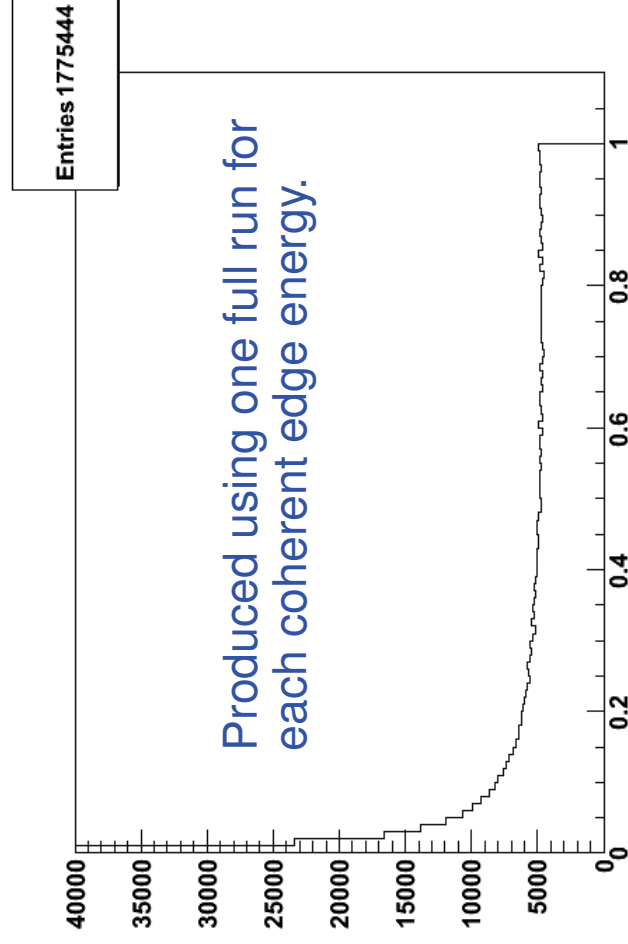
Kinematic Fitting

- Enforces energy-momentum conservation.
- Varies final state measurements (4-vectors) within resolution limits (errors) of said measurements.
 - Errors (as calculated by the reconstruction code) are extracted from data and used to build the tracking covariance matrix (event by event).
 - Errors are then scaled to actual run-dependent resolutions to make the full covariance matrix.

$$\begin{pmatrix} \sigma_{E_\gamma}^2 & 0 & 0 & 0 & \dots & 0 & 0 & 0 \\ 0 & C_1^{pp} & C_1^{p\lambda} & C_1^{p\phi} & \dots & 0 & 0 & 0 \\ 0 & C_1^{p\lambda} & C_1^{\lambda\lambda} & C_1^{\lambda\phi} & \dots & 0 & 0 & 0 \\ 0 & C_1^{p\phi} & C_1^{\lambda\phi} & C_1^{\phi\phi} & \dots & 0 & 0 & 0 \\ \vdots & \vdots & \vdots & \vdots & \ddots & \vdots & \vdots & \vdots \\ 0 & 0 & 0 & 0 & \dots & C_k^{pp} & C_k^{p\lambda} & C_k^{p\phi} \\ 0 & 0 & 0 & 0 & \dots & C_k^{p\lambda} & C_k^{\lambda\lambda} & C_k^{\lambda\phi} \\ 0 & 0 & 0 & 0 & \dots & C_k^{p\phi} & C_k^{\lambda\phi} & C_k^{\phi\phi} \end{pmatrix}$$

Kinematic Fitting

- Quality of the fit is judged by examining the confidence level and pull distributions.
- Every diagonal element of the covariance matrix has a pull value associated with it.
- Systematics affect pull means (μ 's), errors contained in the covariance matrix affect pull widths (σ 's) and the slope of confidence level plots.

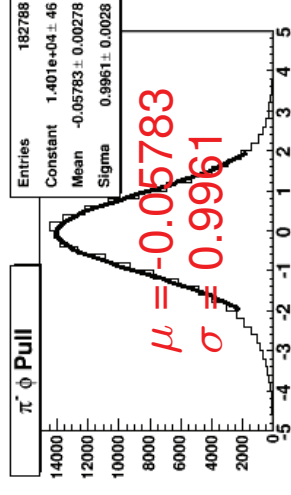
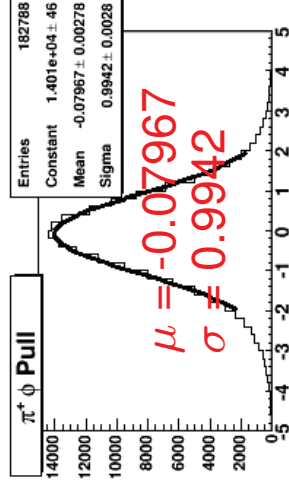
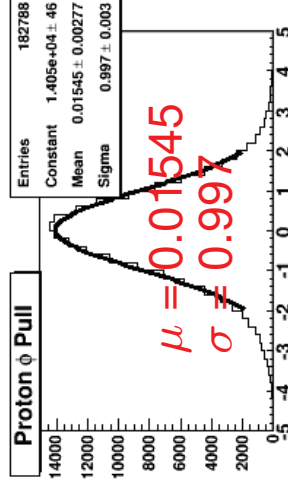
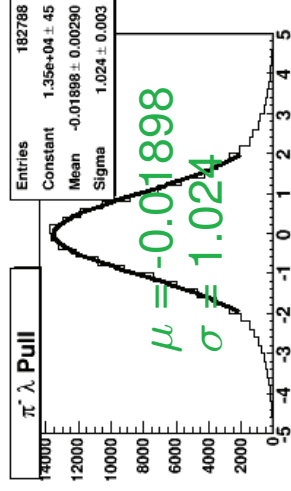
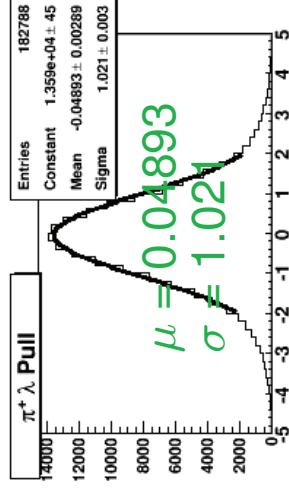
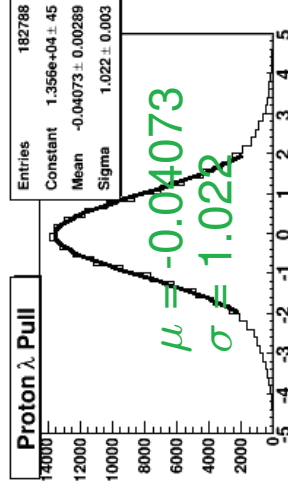
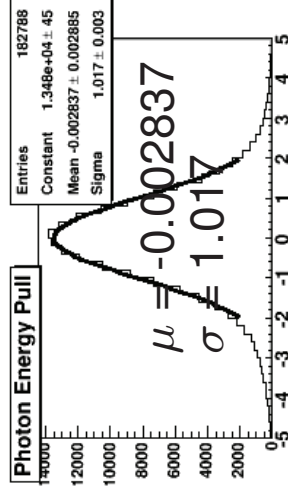
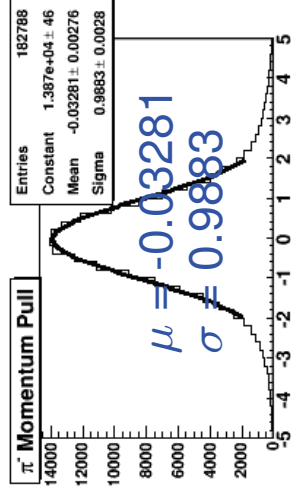
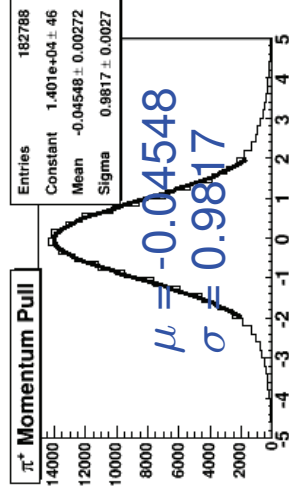
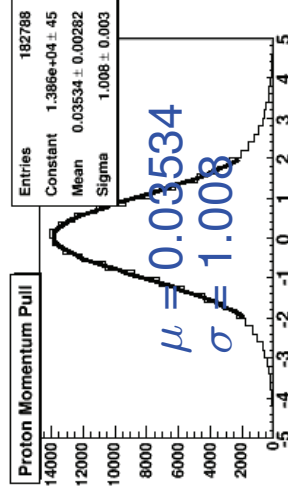


Momentum Corrections and Tagger Sag

- Momentum Corrections:
 - Kinematic fitter is highly sensitive to systematic effects $\rightarrow$ good tool for determining momentum corrections.
 - Corrections are needed to account for variations in torus $\vec{B}$ field and misalignments of Drift Chambers.
 - Determined by examining momentum pull distributions for the proton for $\vec{\gamma}p \rightarrow p \pi^+ \pi^-$ fits/events.
 - Correction factor and pulls binning:
 - 7 bins in θ_{lab} ($10^\circ \leq \theta_{lab} \leq 70^\circ$), 18 bins in ϕ_{lab}
 - 6 bins in momentum ($0.2 \text{ GeV} \leq |\vec{p}| \leq 1.7 \text{ GeV}$)
- Tagger Sag
 - Occurs due to a physical sagging of the support structure for Tagger Hodoscope.
 - Affects the determination of photon energy.
 - Has been seen in several runs in Hall B.
 - Requires energy-dependent correction

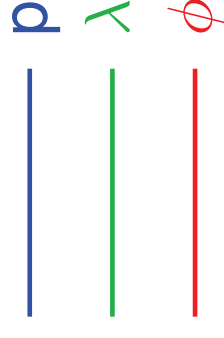
$\rightarrow$ M. Dugger, C. Hanretty, CLAS-Note 2009-030

Check the pull distributions



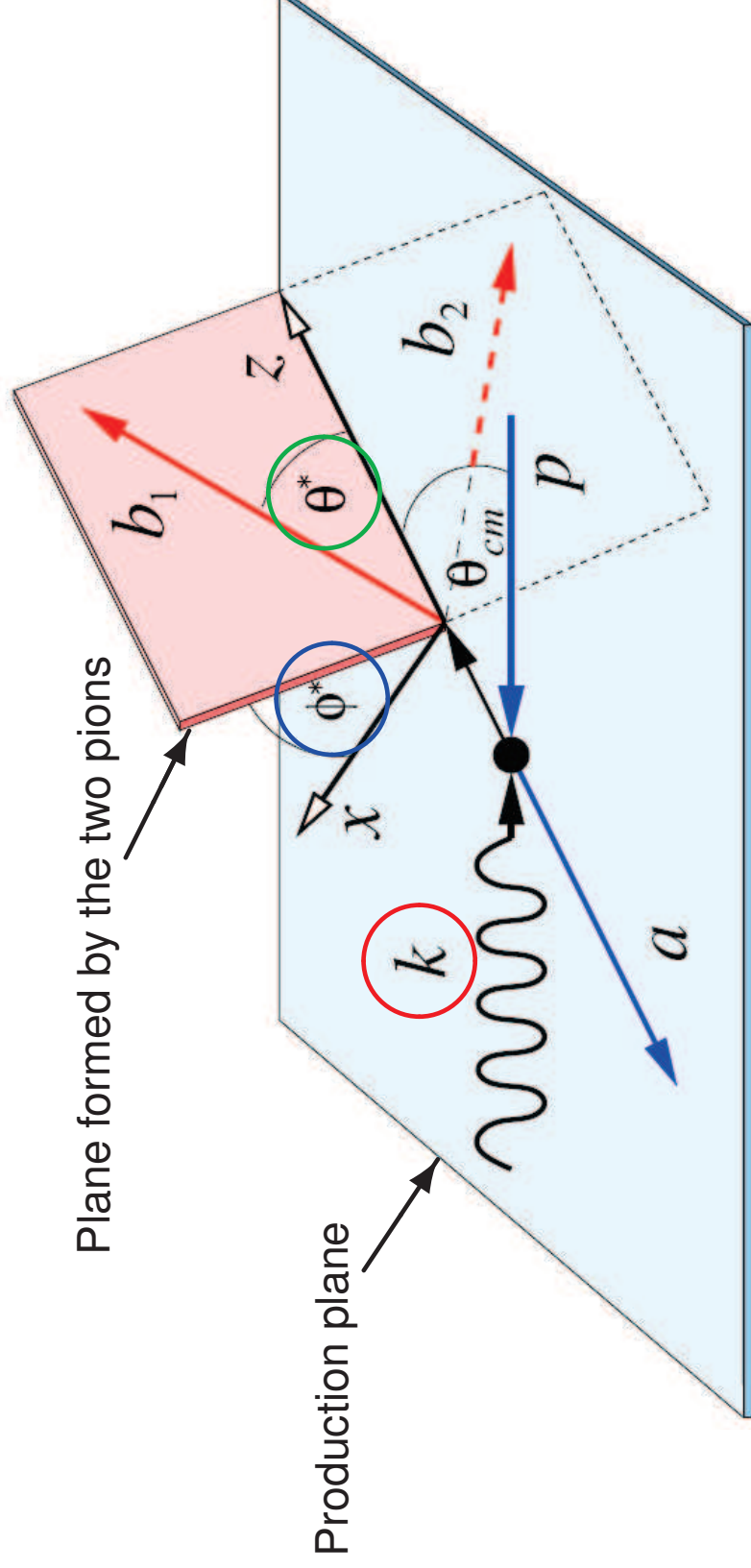
$$\vec{\gamma} p \rightarrow p \pi^+ \pi^-$$

(Run #048326)



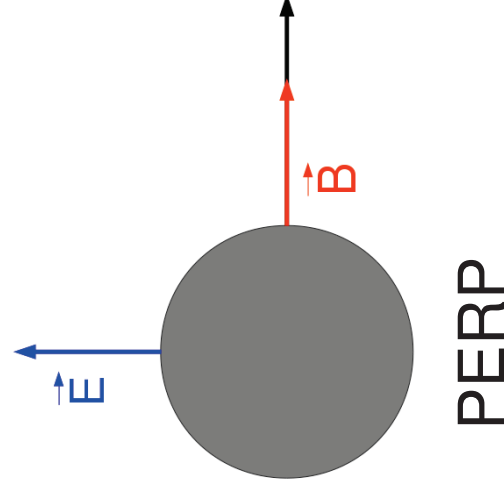
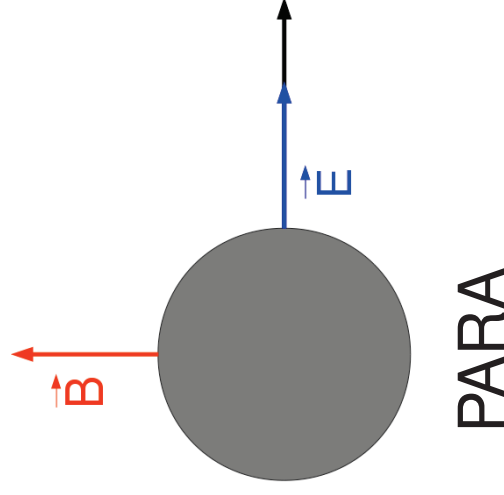
Kinematics

- The analysis of the (double-meson) $\bar{\gamma}p \rightarrow p \pi^+ \pi^-$ channel requires the use of 5 independent kinematic variables: $(m_{p\pi} \text{ or } m_{\pi\pi}), \cos\theta_p^{CM}, \mathbf{k}, \cos\theta_{\pi^+}^*, \phi_{\pi^+}^*$.



Extracting I^s and I^c : ϕ -distributions

- Three types of (photon) polarization settings were used:
 - PARA
 - PERP
 - AMO

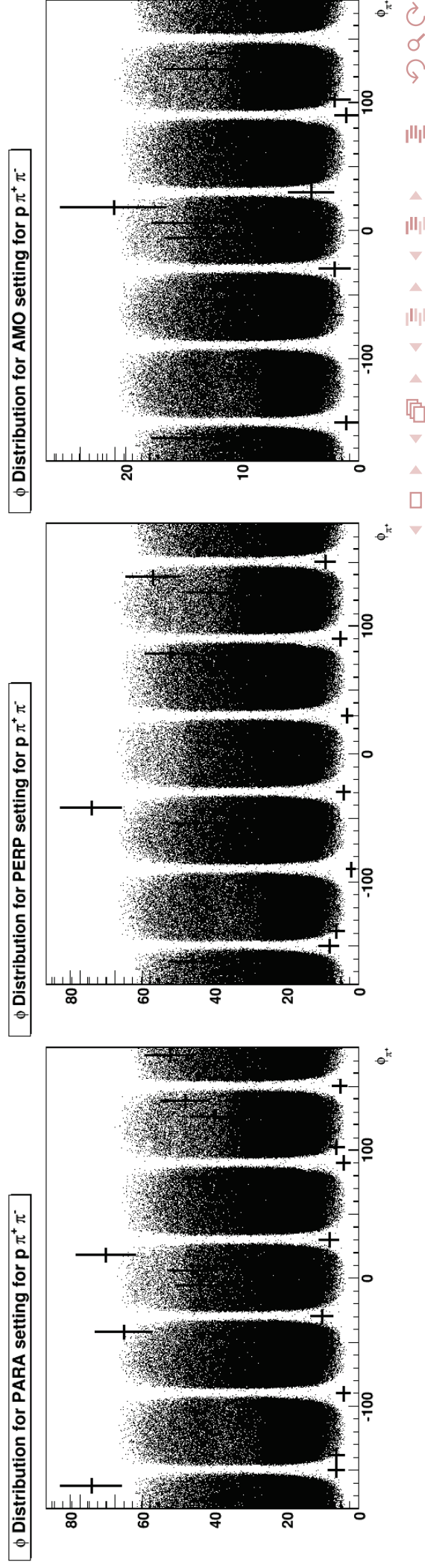


- Events are plotted as a function of lab angle ϕ_{π^+} for each polarization setting.



Extracting I^s and I^c : ϕ -distributions

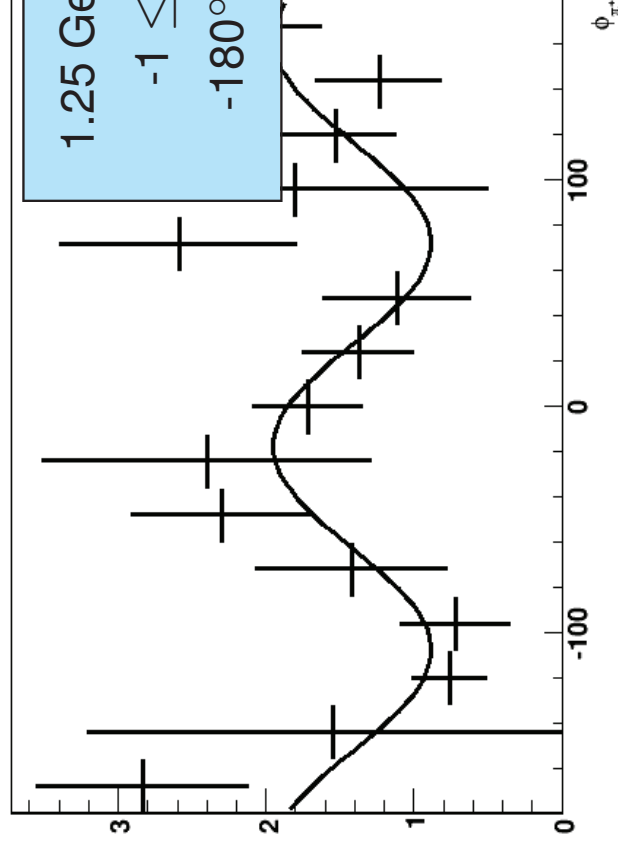
- Three types of (photon) polarization settings were used:
 - PARA
 - PERP
 - AMO
- The polarization of the photon beam breaks the usual ϕ symmetry.
- Events are plotted as a function of lab angle ϕ_{π^+} for each polarization setting.



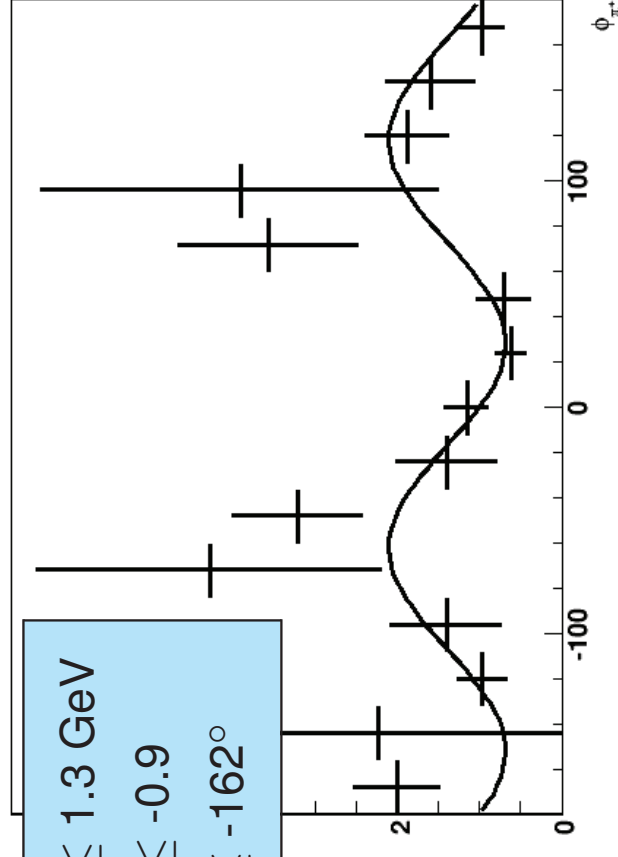
Extracting I^s and I^c : Method 1

- Characteristic: use of AMO ϕ -distributions.
- To remove (divide out) the effects of the experimental setup and acceptance, the ϕ -distributions for PARA and PERP are divided by the AMO ϕ -distributions (for matching bin combinations).
- Fit to $I = I_0 \{ 1 + \delta_I [I^s \sin(2\beta) + I^c \cos(2\beta)] \}$
where $\beta = \phi_{lab} + \phi_{polarization}$

ϕ Distribution for PARA/AMO for $p \pi^+ \pi^-$



ϕ Distribution PERP/AMO for $p \pi^+ \pi^-$



Extracting I^s and I^c : Method 2

- Characteristic: No use of AMO data $\Rightarrow$ uses asymmetry between PARA and PERP.
- Asymmetry between PARA and PERP is formed for matching bin combinations.
- Fit to:

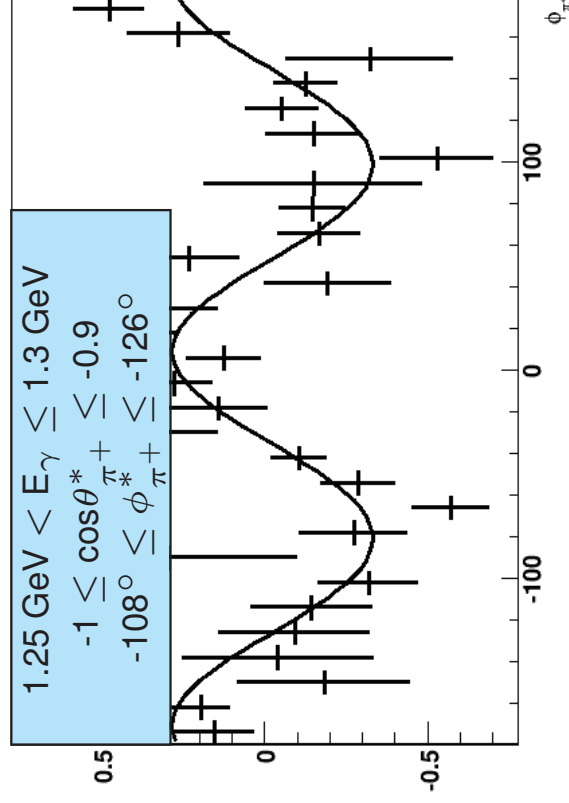
$$A(k, \cos\theta_{\pi^+}^*, \phi_{\pi^+}^*) = \frac{PARA - PERP}{PARA + PERP} = \frac{I_{PARA} - I_{PERP}}{I_{PARA} + I_{PERP}}$$

Extracting I^s and I^c : Method 2

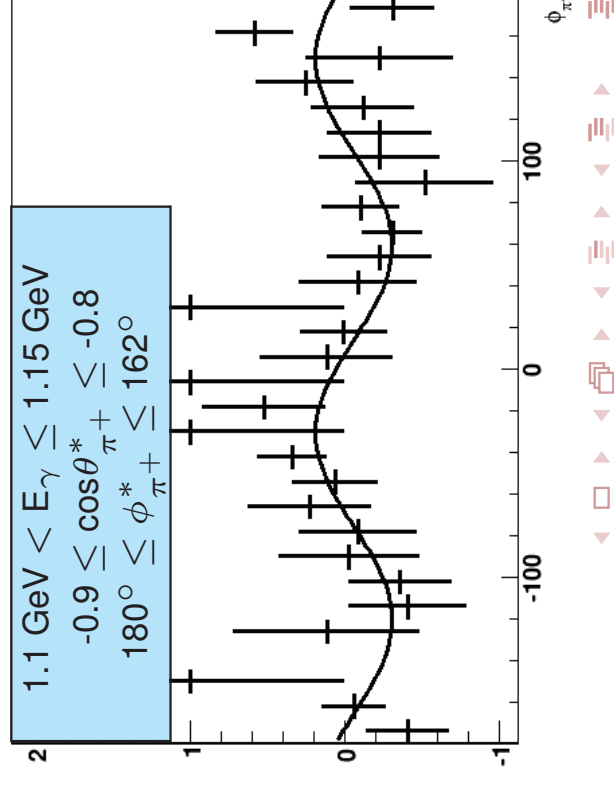
- Characteristic: No use of AMO data $\Rightarrow$ uses asymmetry between PARA and PERP.
- Asymmetry between PARA and PERP is formed for matching bin combinations.
- Fit to:

$$A(k, \cos\theta_{\pi^+}^*, \phi_{\pi^+}^*) = \frac{(\delta_I^{PARA} + \delta_I^{PERP}) I^s \sin(2\beta) + (\delta_I^{PARA} + \delta_I^{PERP}) I^c \cos(2\beta)}{2 + (\delta_I^{PARA} - \delta_I^{PERP}) I^s \sin(2\beta) + (\delta_I^{PARA} - \delta_I^{PERP}) I^c \cos(2\beta)}$$

ϕ Distribution for $p \pi^+ \pi^-$

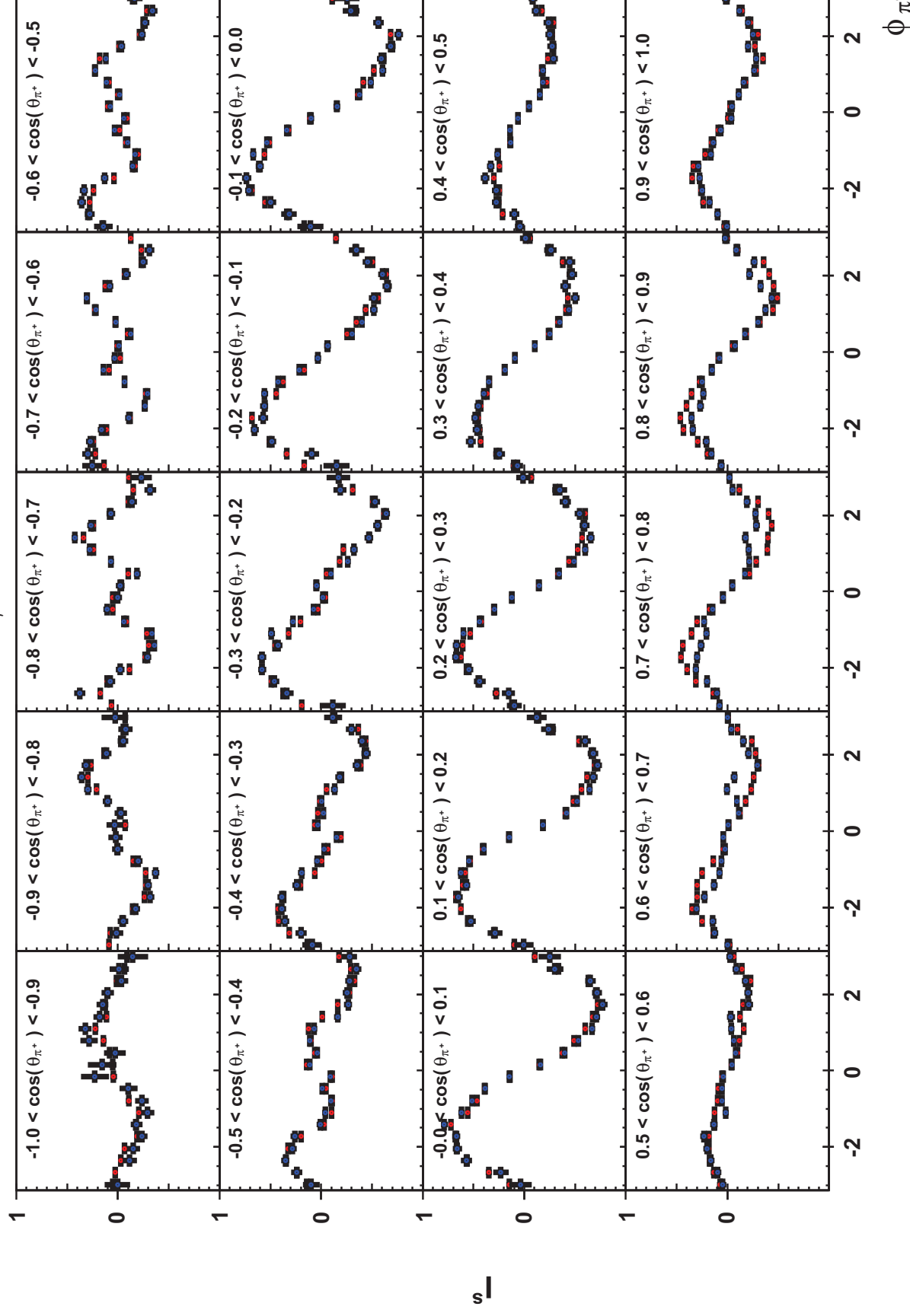


ϕ Distribution for $p \pi^+ \pi^-$



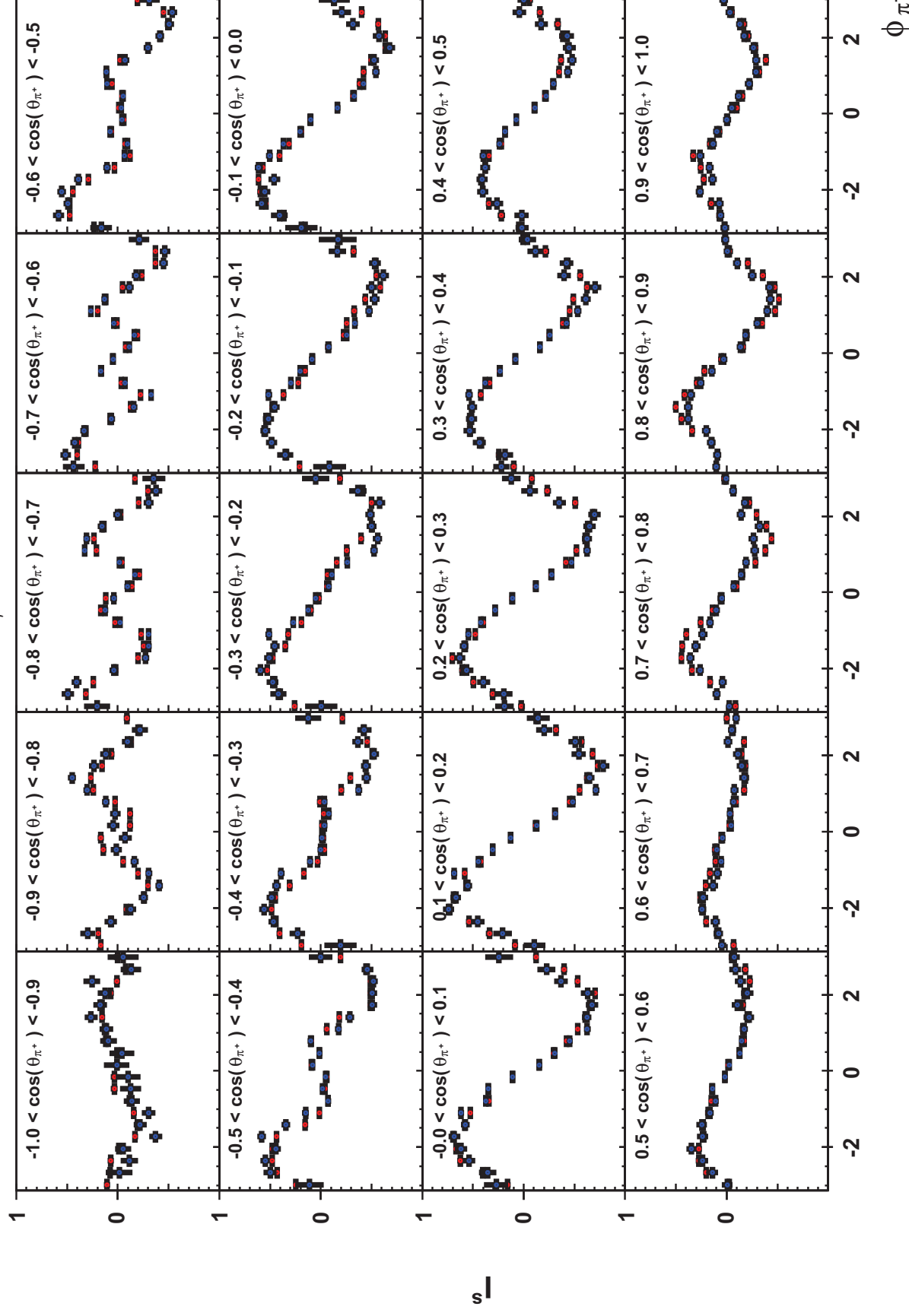
Preliminary Measurement: $\mathbb{I}^s$

$1.2 \text{ GeV} < E_\gamma < 1.25 \text{ GeV}$



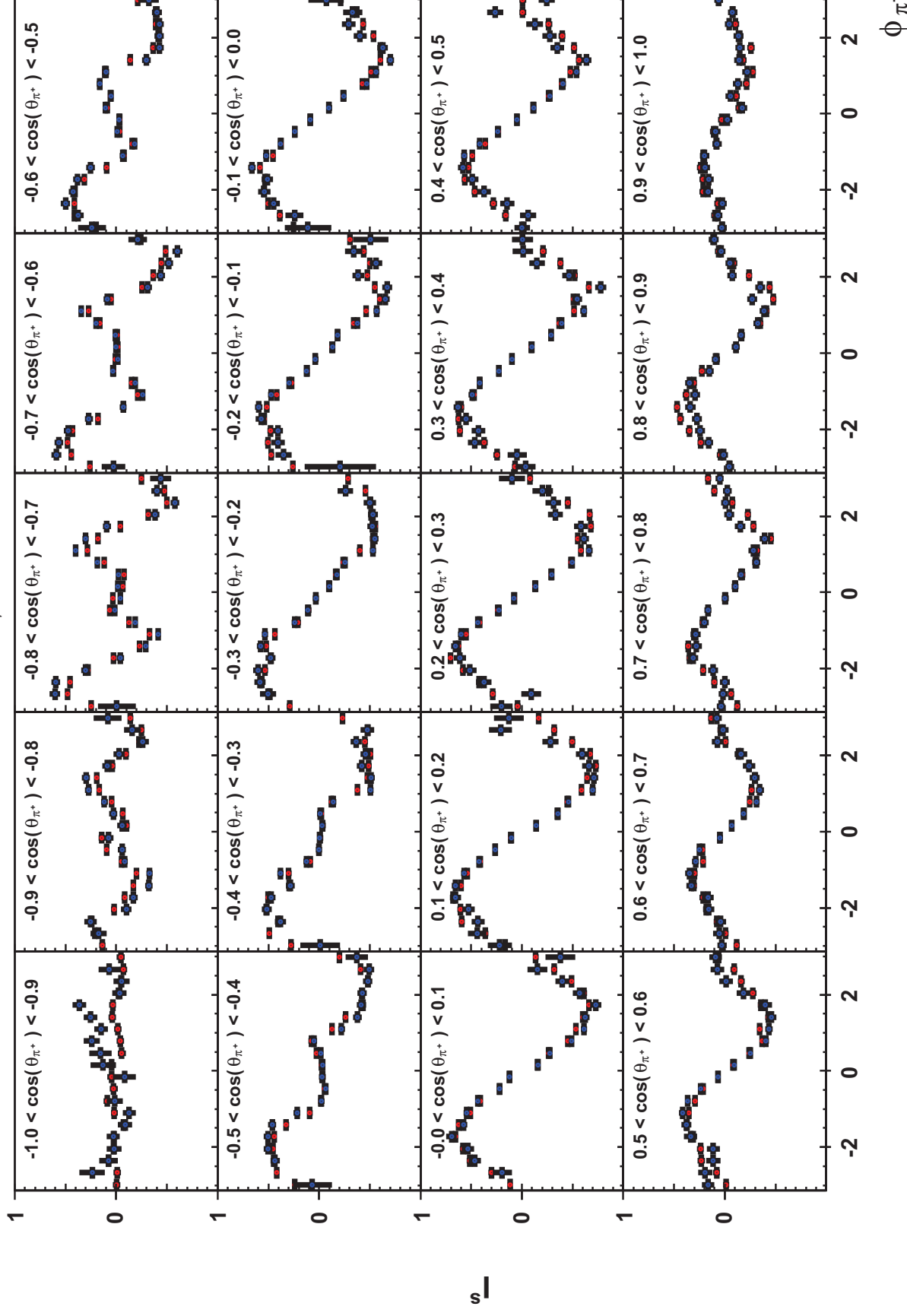
Preliminary Measurement: $\mathbb{I}^s$

$1.4 \text{ GeV} < E_\gamma < 1.45 \text{ GeV}$



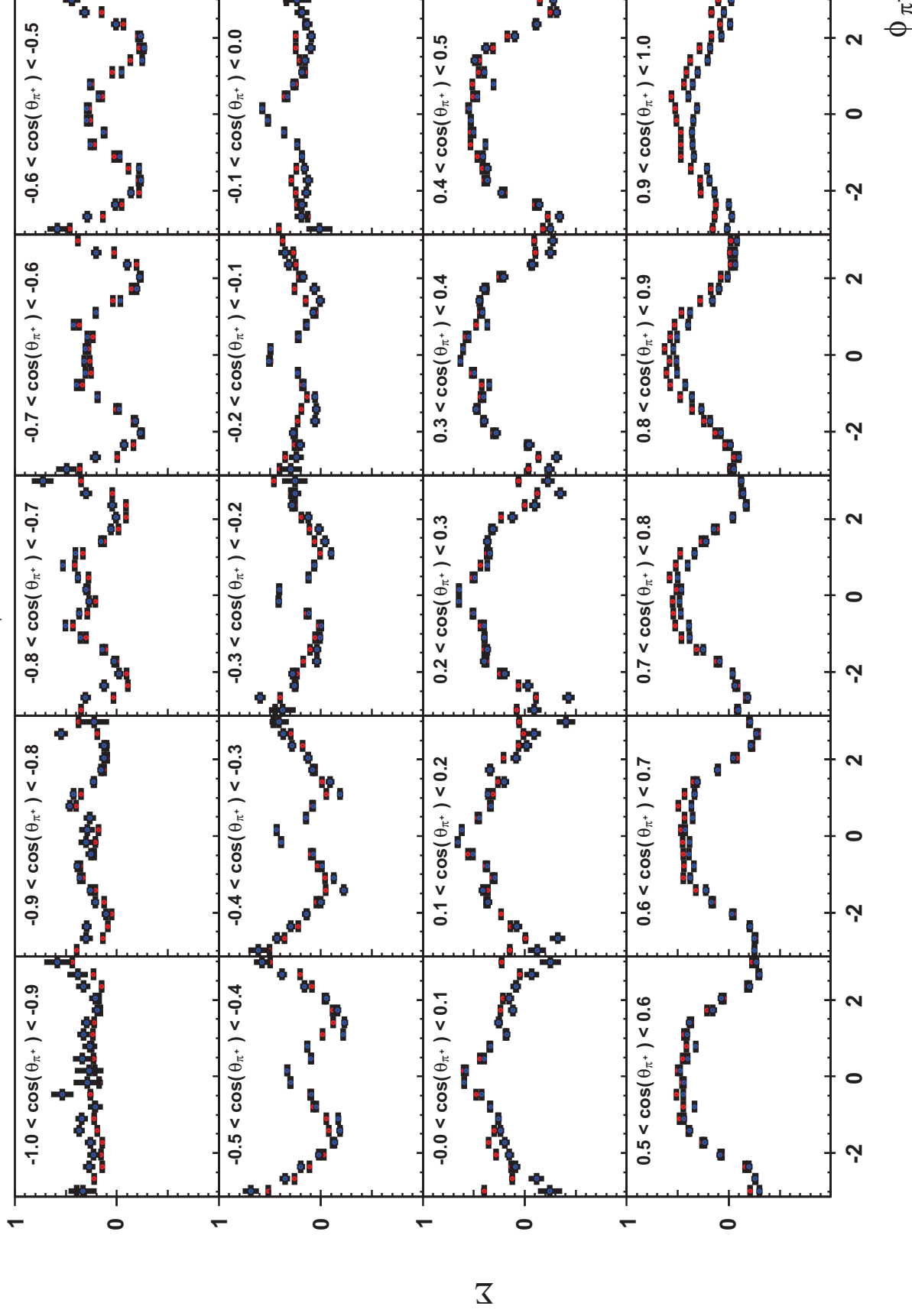
Preliminary Measurement: I^s

$1.6 \text{ GeV} < E_\gamma < 1.65 \text{ GeV}$



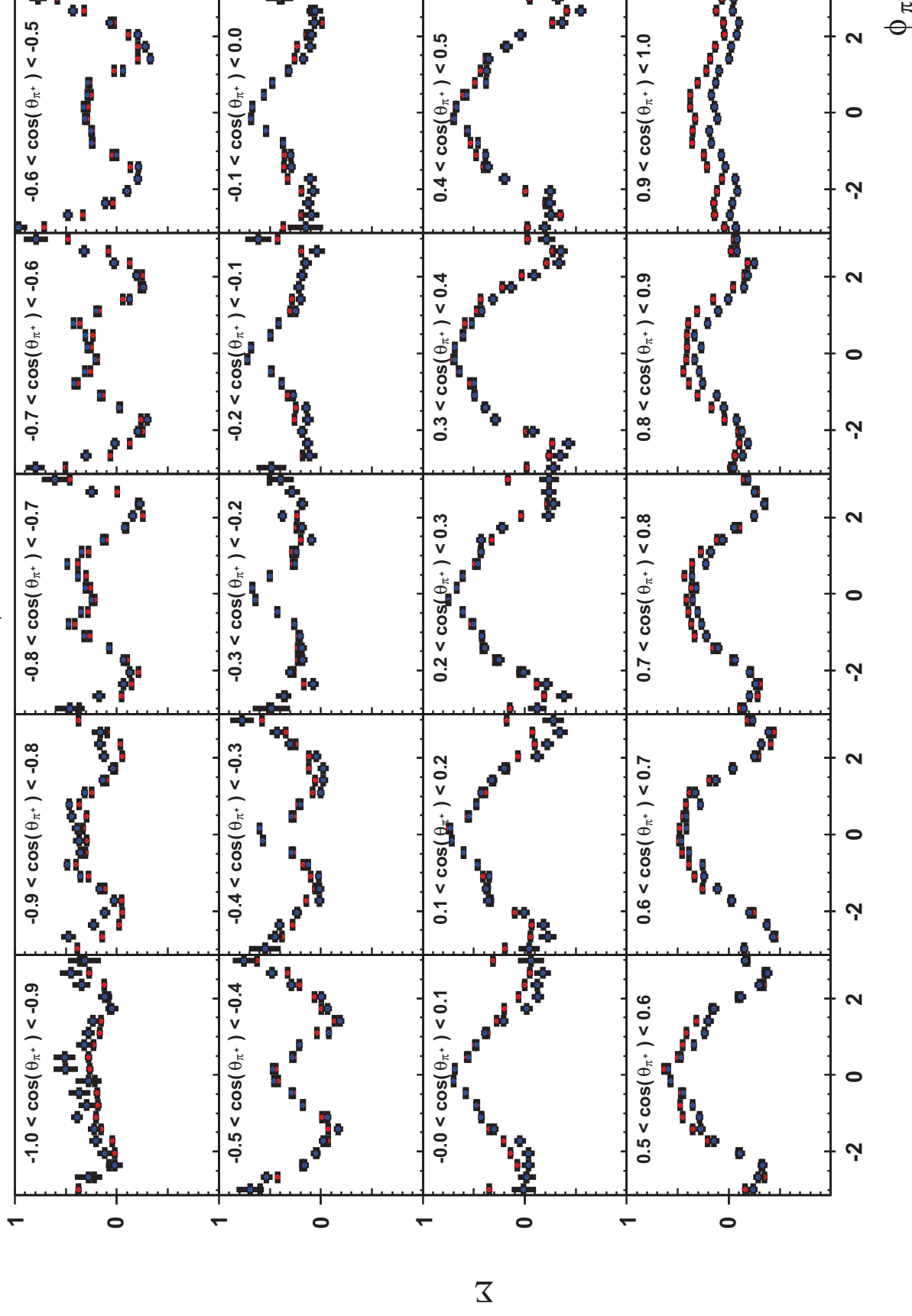
Preliminary Measurement: $\mathbb{I}^c$

$1.2 \text{ GeV} < E_\gamma < 1.25 \text{ GeV}$



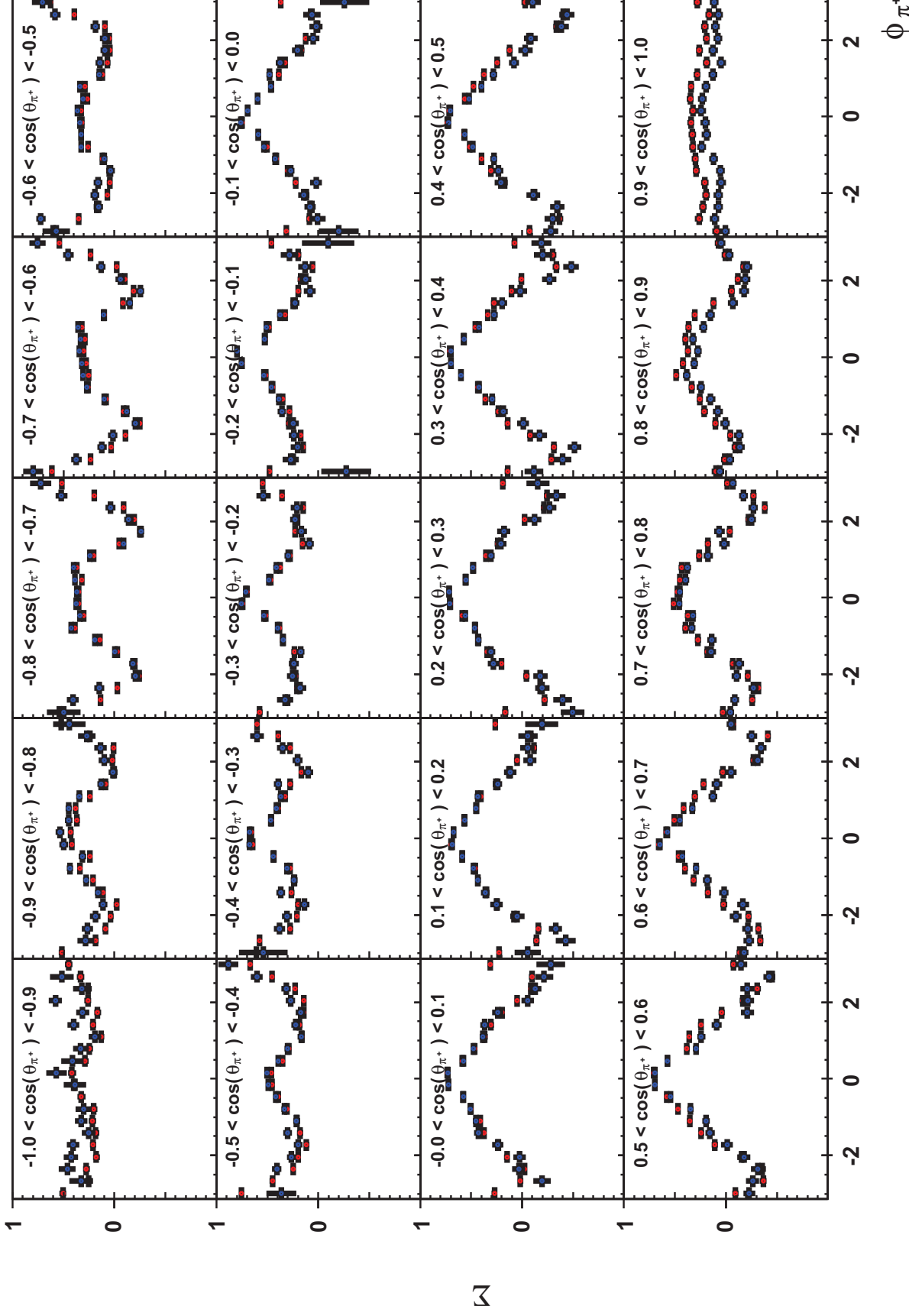
Preliminary Measurement: $\mathbb{I}^c$

$1.4 \text{ GeV} < E_\gamma < 1.45 \text{ GeV}$



Preliminary Measurement: $\mathbb{I}^c$

$1.6 \text{ GeV} < E_\gamma < 1.65 \text{ GeV}$



Summary

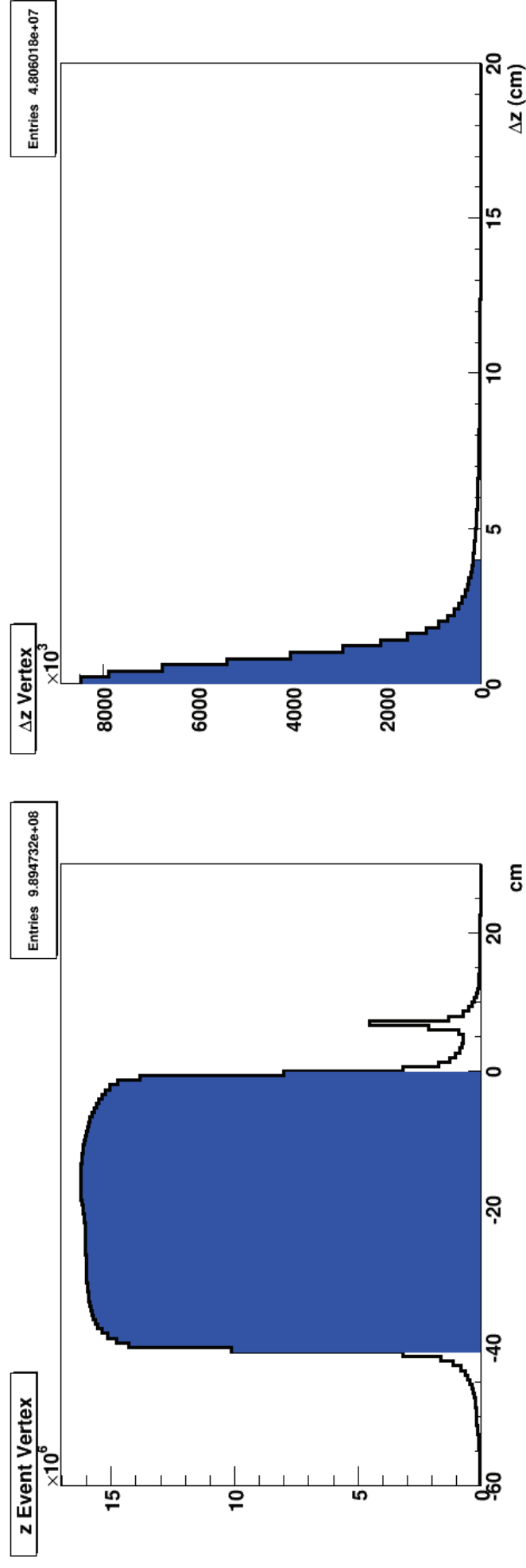
- Hadrons form the basis of most matter.
- These hadrons are systems containing quarks and gluons.
- Through the study of these systems when excited, the internal dynamics of the hadron may be demystified.
- Polarization observables are sensitive to these dynamics.
- Measured for $\vec{\gamma} p \rightarrow p \pi^+ \pi^-$:
 - I^C (first for a two-charged-pion final state)
 - I^S (first for a two-pion final state)
- These observable measurements as well as future measurements will aid in the refinement of CQMs as well aid in the revelation of the internal dynamics of the baryon.
- The FROST experiment will provide access to all 15 observables for a double-meson final state.

END

Backup Slides

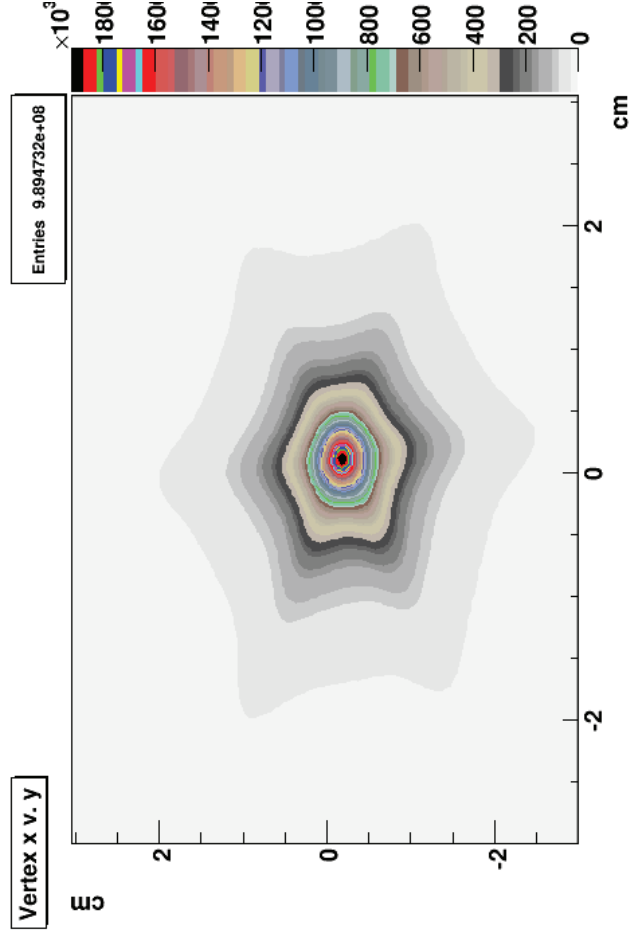
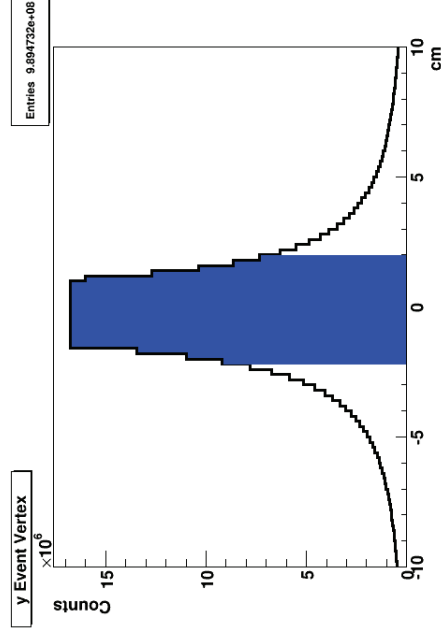
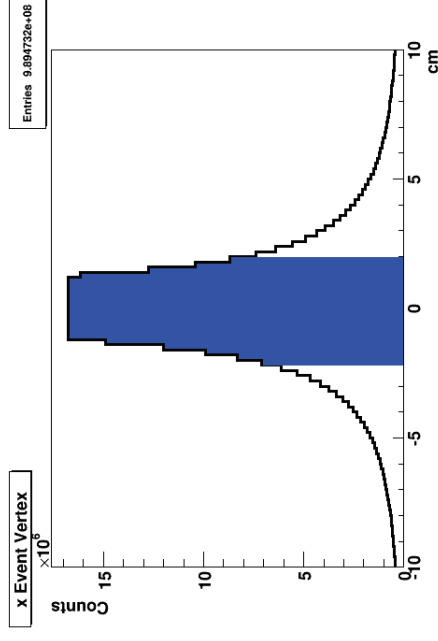
Cuts

- Vertex
 - $-40 \text{ cm} < z_{\text{vertex}} < 0 \text{ cm}$ (for all detected particles)
 - All z-vertices must be within 4 cm of each other.
 - Cut on (event) x,y-vertex requiring it originates in the target cell.
 - $0 \text{ ns} < t_{\text{vertex}} < 20 \text{ ns}$ (for all detected particles)



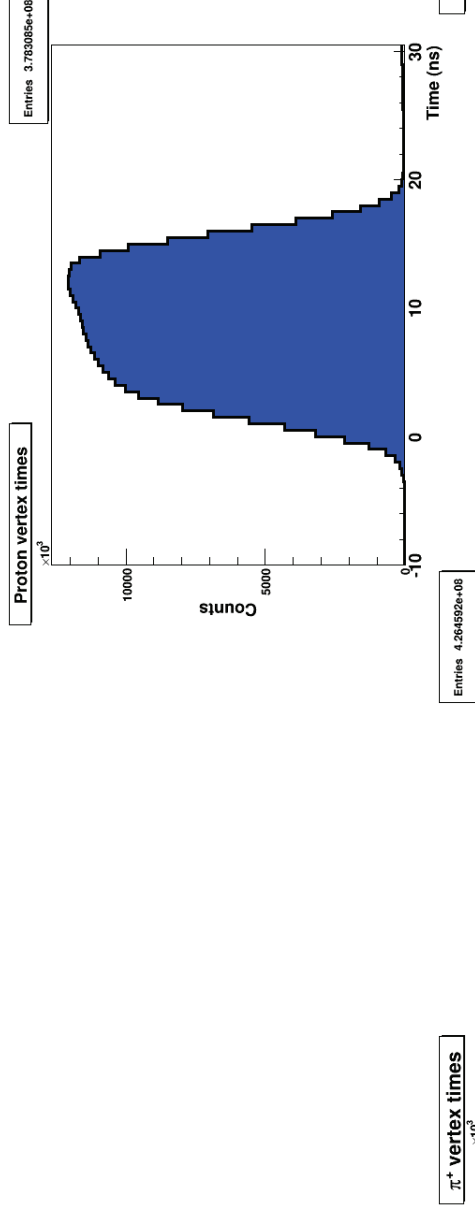
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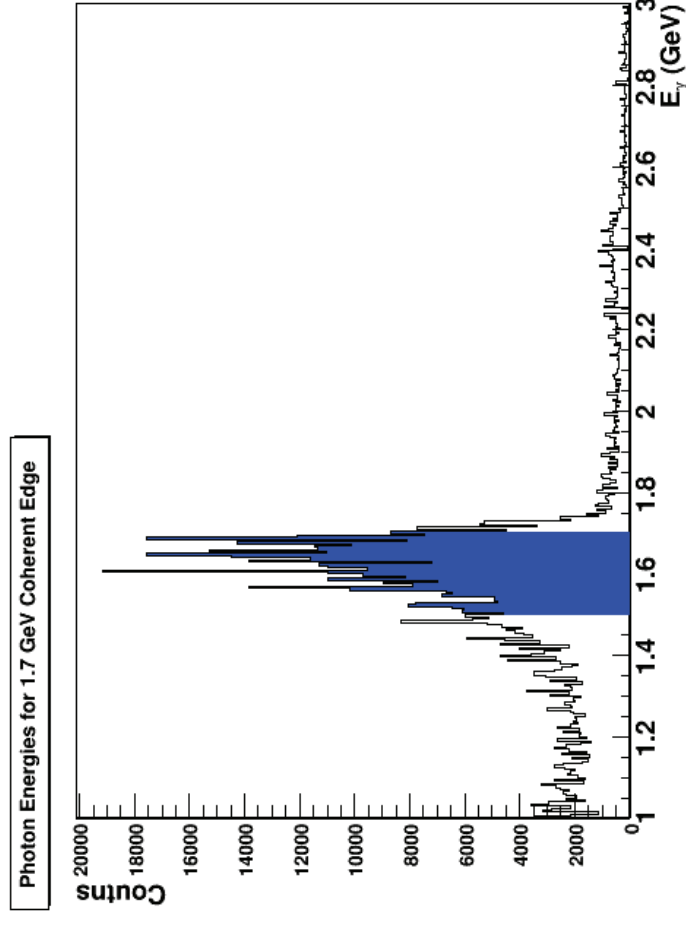
Cuts

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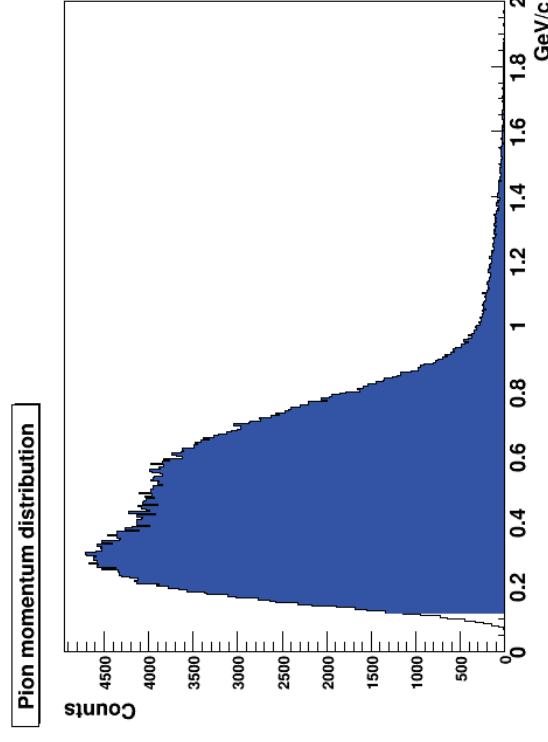
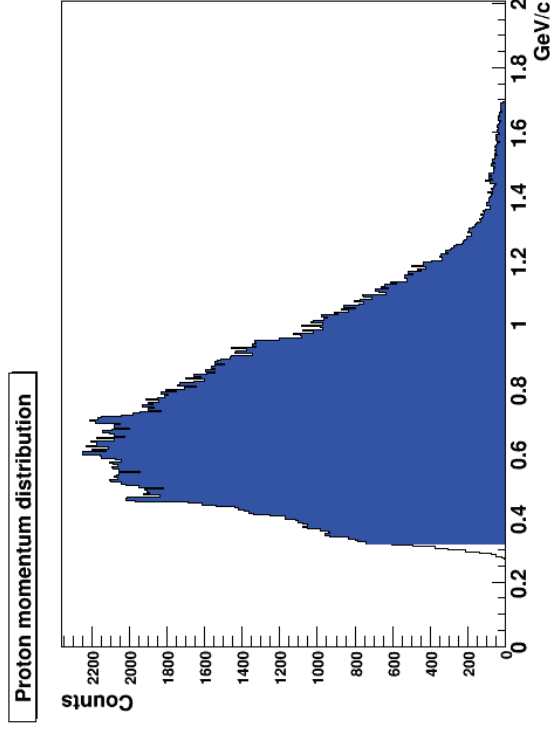
Cuts

- Photon Energy
 - $1.1 \text{ GeV} < E_\gamma < 1.3 \text{ GeV}$
 - $1.3 \text{ GeV} < E_\gamma < 1.5 \text{ GeV}$
 - $1.5 \text{ GeV} < E_\gamma < 1.7 \text{ GeV}$
 - $1.7 \text{ GeV} < E_\gamma < 1.9 \text{ GeV}$
 - $1.9 \text{ GeV} < E_\gamma < 2.1 \text{ GeV}$
- Momentum and Angle
 - $p_{\text{proton}} > 320 \text{ MeV}/c$
 - $p_{\pi^\pm} > 125 \text{ MeV}/c$
 - $\theta > 10^\circ$
 - $\theta_{\pi^\pm} < 120^\circ$
- Other Cuts
 - Confidence level cut of 5%.
 - ‘ngrf’ = 1 for all particles.
 - ‘tagrid’ same for all particles.



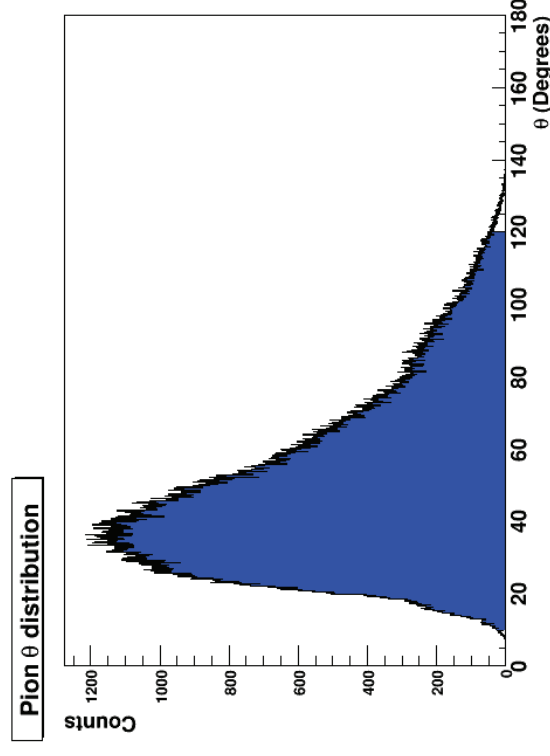
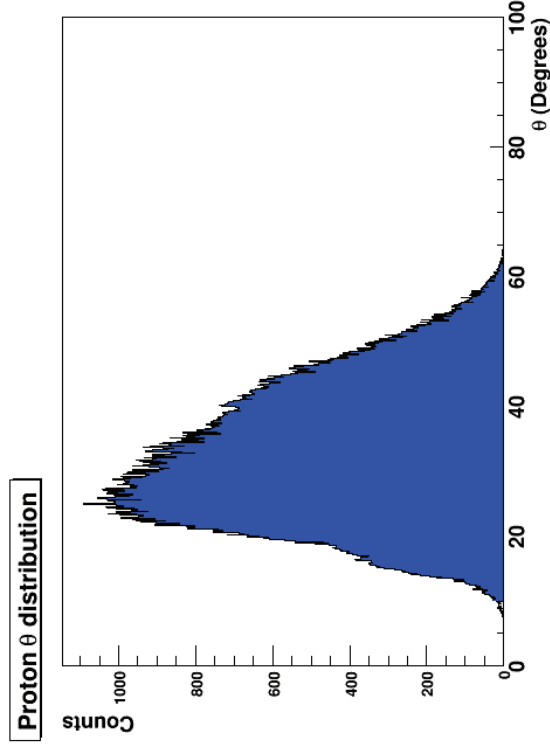
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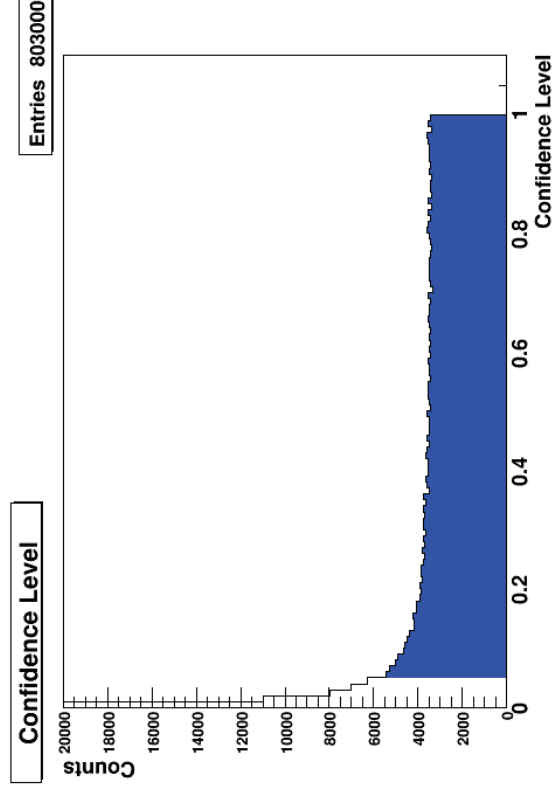
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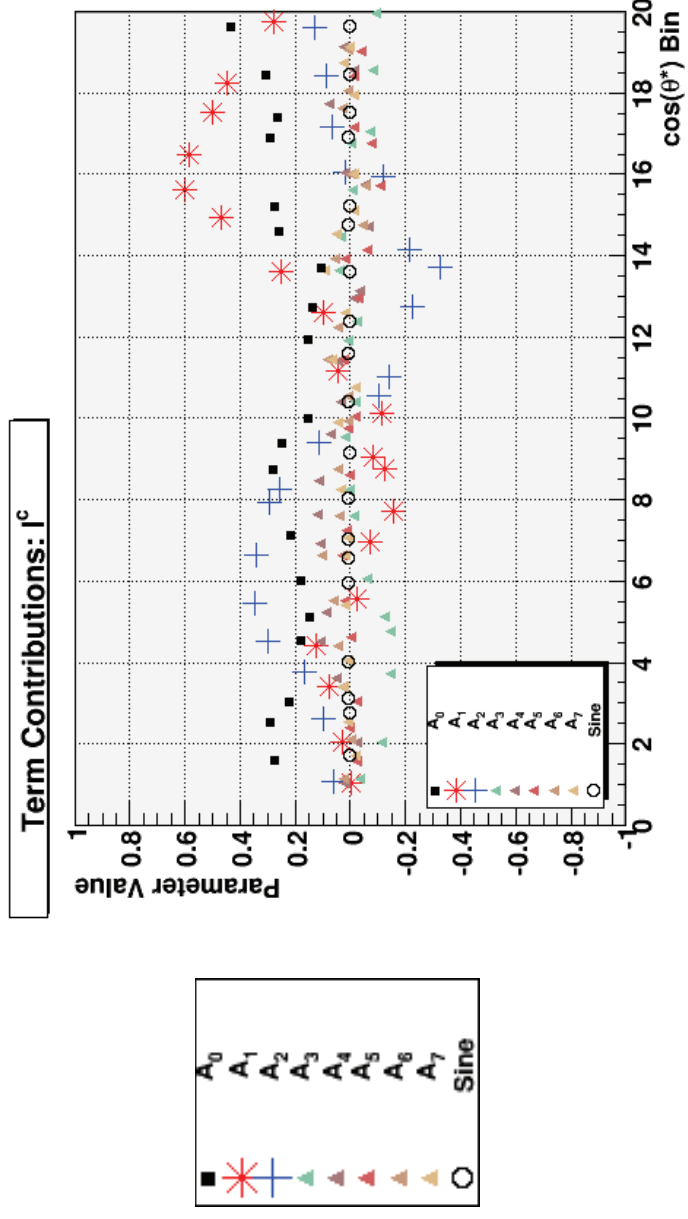
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Examination of the behavior of I^s and I^c

- The behavior of the measured observables is predicted to be either even (I^c) or odd (I^s) as a function of ϕ^* .
- To check the behavior, measurements were fit with expansions of sine and cosine for each $\cos\theta^*$ bin.
- Contributions from the different terms were examined.

$$f(\phi^*) = A_0 + A_1 \cos(\phi^*) + A_2 \cos(2\phi^*) + A_3 \cos(3\phi^*) + A_4 \cos(4\phi^*) \\ + A_5 \cos(5\phi^*) + A_6 \cos(6\phi^*) + A_7 \cos(7\phi^*) + A_8 \sin(\phi^*)$$



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