

# Magnetic structure and excitations in vanadates, $\text{BaV}_{10}\text{O}_{15}$ and $\text{CoV}_2\text{O}_4$

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Sachith Dissanayake  
(University of Virginia)

4<sup>th</sup> year research seminar

# Outline

- **Introduction**

- Frustrated magnetism, , spin and orbital degrees of freedom
- neutron scattering basics

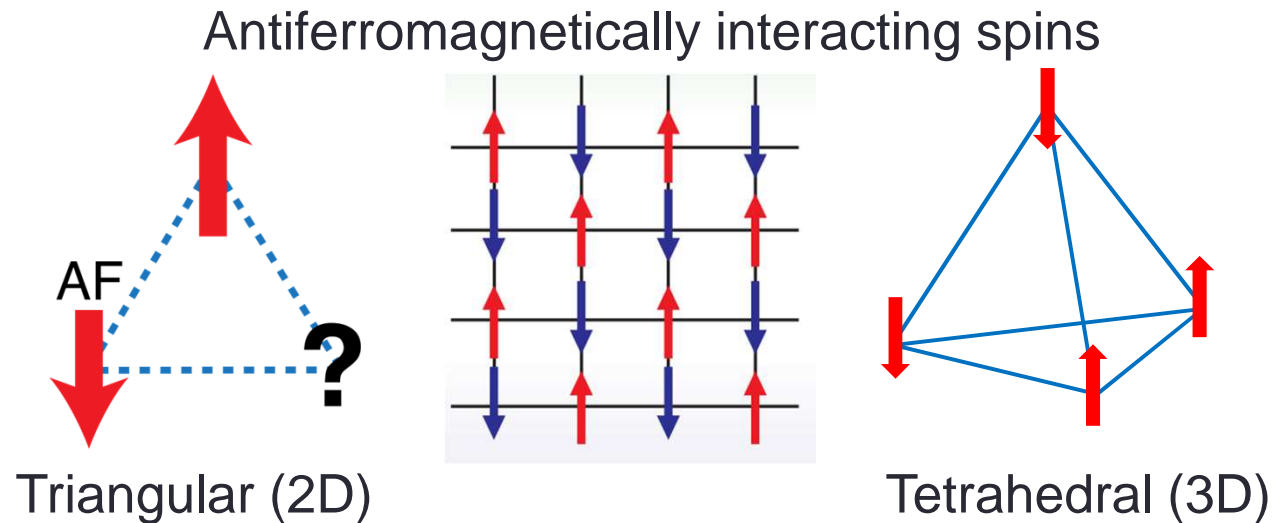
- **BaV<sub>10</sub>O<sub>15</sub>**

- Background
- Experimental data – Diffraction and Inelastic
- Results – Magnetic ground state and excitations
- Conclusions

- **CoV<sub>2</sub>O<sub>4</sub>**

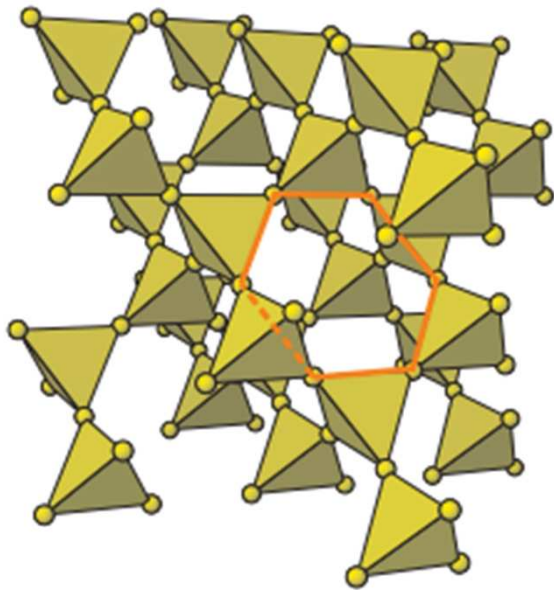
- Background
- Experimental data – Diffraction and Inelastic
- Results – Magnetic Structure and linear spin wave calculation
- Conclusions

# What is Frustration?

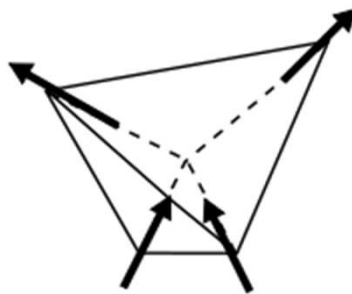


- When a system is geometrically frustrated, absence of a unique ground state gives a finite entropy at absolute zero temperature.
- Some materials may have many nearly-degenerate ground states (a spin glass), or may retain dynamic disorder (a quantum spin liquid).
- Orbital degree of freedom can lift the degeneracy and drive the system into a particular ground state or reduce the frustration.

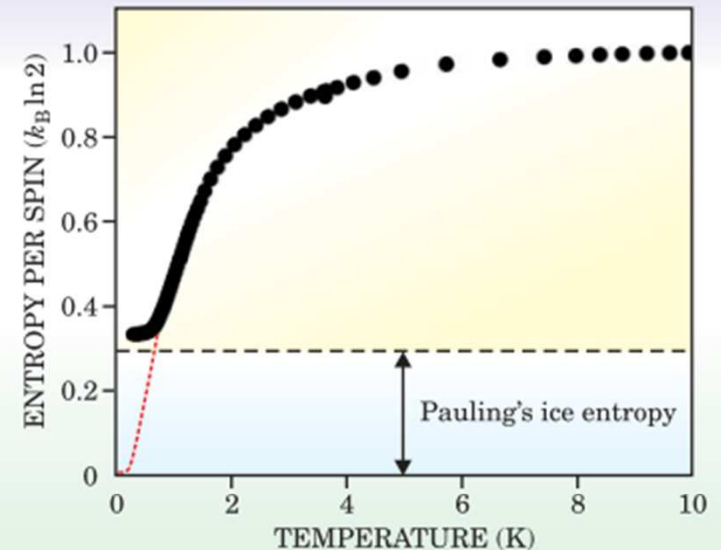
# What is Frustration?



Most degenerate- and thus most frustrated-lattice  
Readily realizable in three dimensions or less is  
one made up with corner sharing tetrahedra, the  
pyrochlore lattice.



Roderich Moessner and Arthur P. Ramirez  
Physics Today / Volume 59 / Issue 2 /



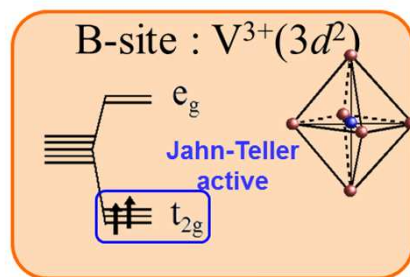
**Figure 3.** The entropy  $S$  of the spin-ice compound dysprosium titanate  $\text{Dy}_2\text{Ti}_2\text{O}_7$  as a function of temperature  $T$ . At high  $T$ ,  $S = k_B \ln 2$  per spin, the same as for free spins. Cooling the system causes the entropy to drop as correlations develop between spins. An unfrustrated magnet follows the plot's schematic red line down to zero entropy because the system assumes a unique ground state. In spin ice, geometrical frustration creates an exponentially large number of degenerate ground states. The large degeneracy manifests itself in a non-vanishing entropy, which is close to the value that Linus Pauling predicted for ordinary water ice. (Adapted from ref. 5.)

# Frustrated magnets

- Lies in two fundamental enterprises in condensed matter physics.
- On the applied side, the instabilities exhibited by frustrated magnets open a window on richness of nature realized in different materials
- On the fundamental side is the search for principles that help organize the variety of behavior we observed around us.

# Jahn–Teller effect

- Any nonlinear molecule with a spatially degenerate electronic ground state will undergo a geometrical distortion that removes that degeneracy, because the distortion lowers the overall energy of the species.
- The Jahn–Teller effect is most often encountered in octahedral complexes of the transition metals.



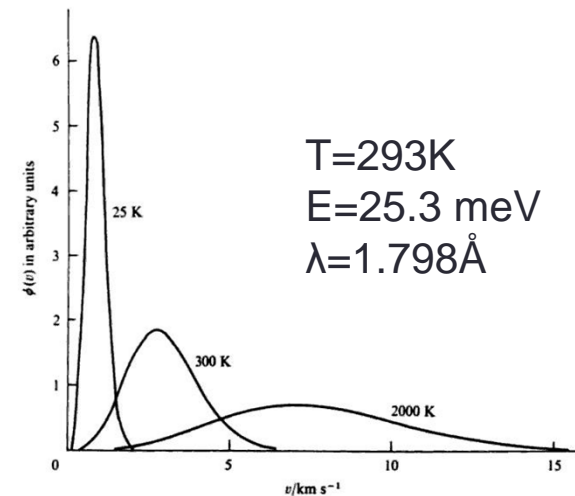
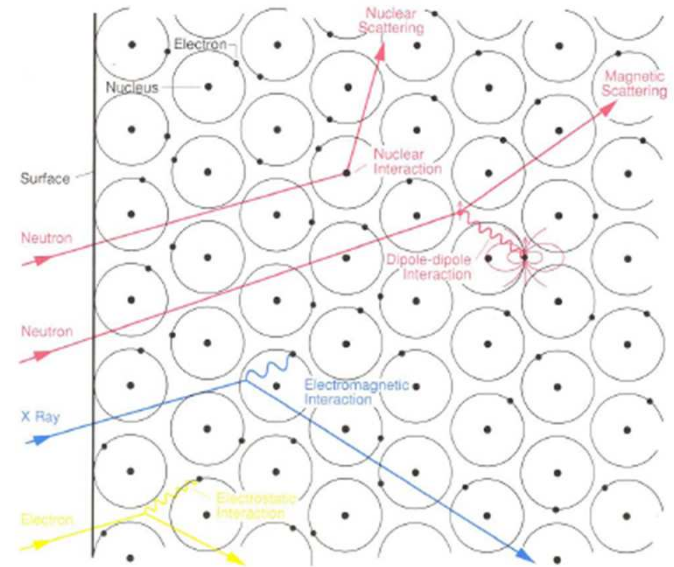
Number of d electrons	1	2	3	4	5	6	7	8	9	10
High spin	w	w		s		w	w		s	
Low spin	w	w		w	w		s		s	

**w**: weak Jahn–Teller effect ( $t_{2g}$  orbitals unevenly occupied),  
**s**: strong Jahn–Teller effect expected ( $e_g$  orbitals unevenly occupied), **blank**: no Jahn–Teller

# Neutron Scattering

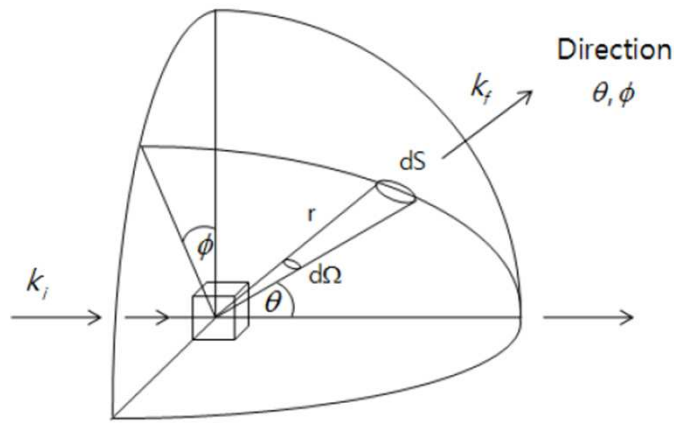
Why neutron?

1. **Zero net charge** makes neutron to penetrate deep inside the sample without being scattered by the electrons like X-ray, hence measuring bulk properties.
2. Its internal **magnetic moment** interact with the unpaired electrons spins in an atom via dipole-dipole interaction enabling direct measurement of magnetic moments.
3. Due to **large neutron mass**, when moderated with water or liquid hydrogen, the kinetic energy becomes comparable to that of many elementary excitations in a solid.



Maxwell distribution of thermal neutron

# What we measure



The scattering cross section, the actual number measured by the detector, is proportional to the scattering function,  $S(\vec{Q}, \omega)$  which is the Fourier transform of time dependent pair correlation function of either atomic density or spin component.

nuclear scattering

$$\frac{d^2\sigma}{d\Omega_f dE_f} = N \frac{k_f}{k_i} b^2 S(\vec{Q}, \omega)$$

$$\begin{aligned} S(\vec{Q}, \omega) &= \frac{1}{2\pi\hbar N} \sum_{ll'} \int_{-\infty}^{\infty} dt \langle e^{-i\vec{Q}\cdot\vec{r}_{l'}(0)} e^{i\vec{Q}\cdot\vec{r}_l(t)} \rangle e^{-i\omega t} \\ &= \frac{1}{2\pi\hbar N} \int_{-\infty}^{\infty} dt \langle \rho_{\vec{Q}}(0) \rho_{-\vec{Q}}(t) \rangle e^{-i\omega t}. \end{aligned}$$

$$\rho_{\vec{Q}}(t) = \sum_l e^{i\vec{Q}\cdot\vec{r}_l(t)}$$

magnetic scattering

$$\frac{d^2\sigma}{d\Omega_f dE_f} = \frac{N}{\hbar} \frac{k_f}{k_i} p^2 e^{-2W} \sum_{\alpha, \beta} (\delta_{\alpha, \beta} - \hat{Q}_{\alpha} \hat{Q}_{\beta}) S^{\alpha\beta}(\vec{Q}, \omega)$$

$$S^{\alpha, \beta}(\vec{Q}, \omega) = \frac{1}{2\pi} \int_{-\infty}^{\infty} dt e^{-i\omega t} \sum_l e^{i\vec{Q}\cdot\vec{r}_l} \langle S_0^{\alpha}(0) S_l^{\beta}(t) \rangle$$

$$pS = \left(\frac{\gamma r_0}{2}\right) g f(\vec{Q}) S$$



# Neutron scattering

$$\vec{Q} = \vec{k}_i - \vec{k}_f$$

$$\hbar\omega = \Delta E = \hbar/(2m_n)(k_i^2 - k_f^2)$$

- **Elastic Neutron Scattering (Neutron diffraction)**

- Reveals the crystal structure.
- Reveals the microscopic magnetic structure

Instruments – Powder and single crystal diffractometer

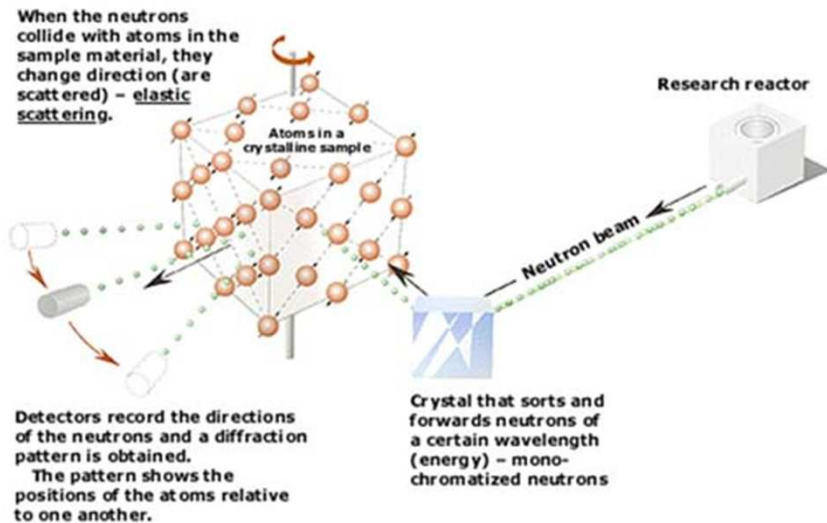
- **Inelastic Neutron Scattering**

- To study atomic and molecular motion as well as magnetic and crystal field excitations.

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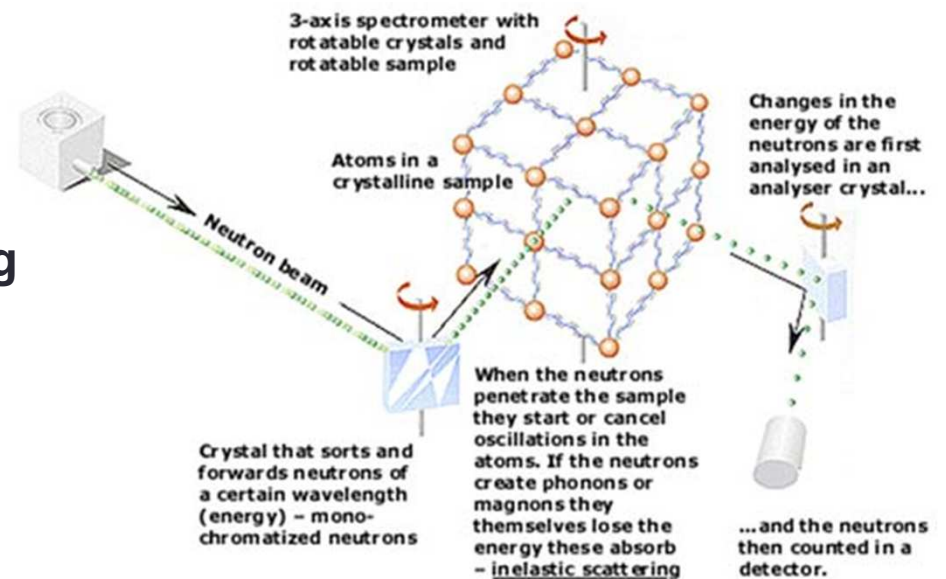
Instruments – Triple axis spectrometer

Time of flight spectrometer



## Elastic Neutron Scattering

## Inelastic Neutron Scattering



# Neutron diffraction

**Powder Diffractometer**



**HB2A, HFIR**

**Single Crystal Diffractometer**

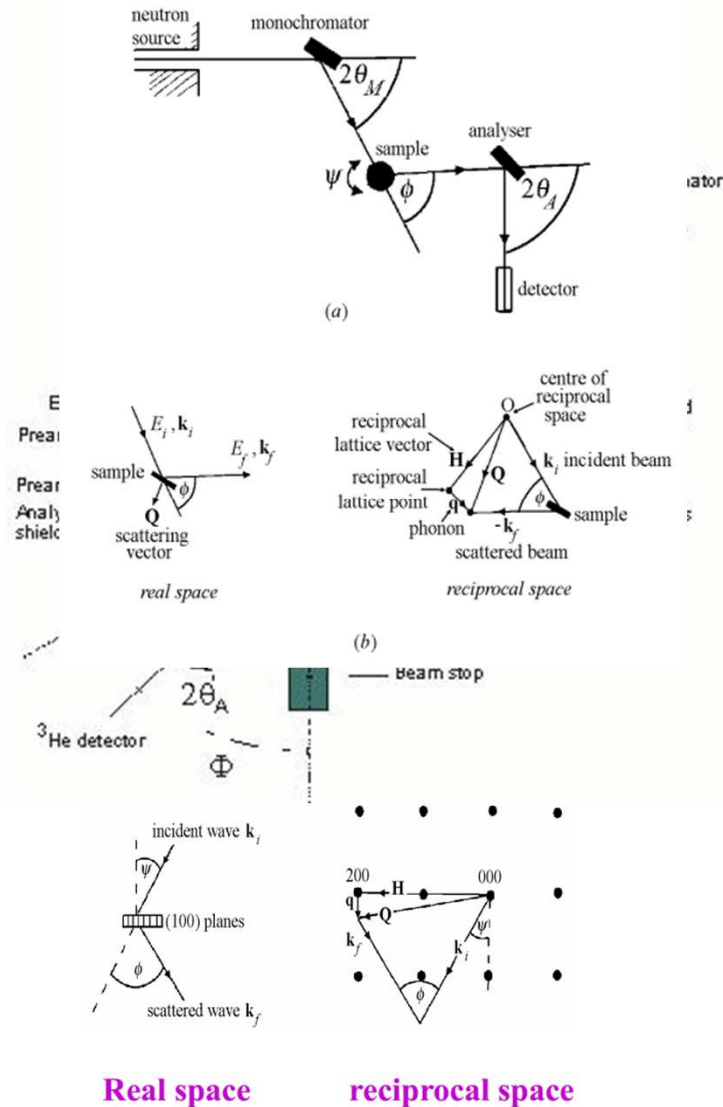


**RESI, FRM II**

Bragg's law  $n\lambda = 2d \sin \theta$

# Triple axis spectrometer

### HB-1 Schematic Drawing



It allows measurement of the scattering function at any point in energy and momentum space physically accessible by the spectrometer.

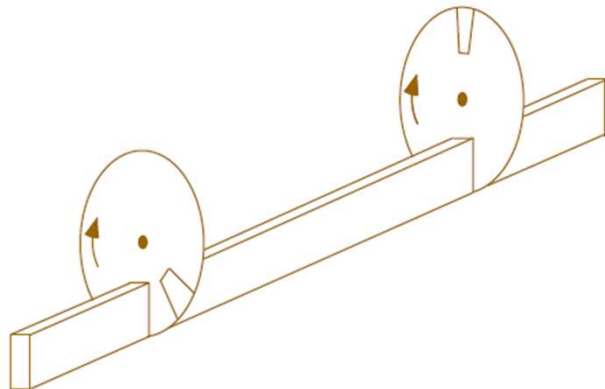
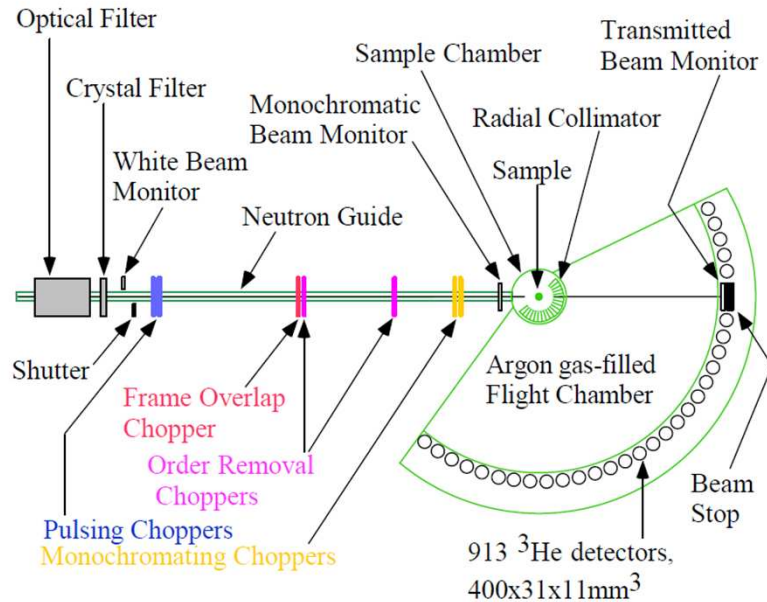
$$Q = k_i - k_f$$

$$\hbar\omega = E_i - E_f = \frac{\hbar^2}{2m} (k_i^2 - k_f^2)$$

$$|\mathbf{Q}|^2 = |\mathbf{k}_i|^2 + |\mathbf{k}_f|^2 - 2 |\mathbf{k}_i| |\mathbf{k}_f| \cos (2\theta)$$



# Time of Flight measurement



The time of flight spectrometer, as can be expected from its name, determines neutron energy by measuring its time of flight from one point to another.

The beams are monochromated using several choppers rotating at different frequencies allowing only neutrons at certain velocity can pass through.

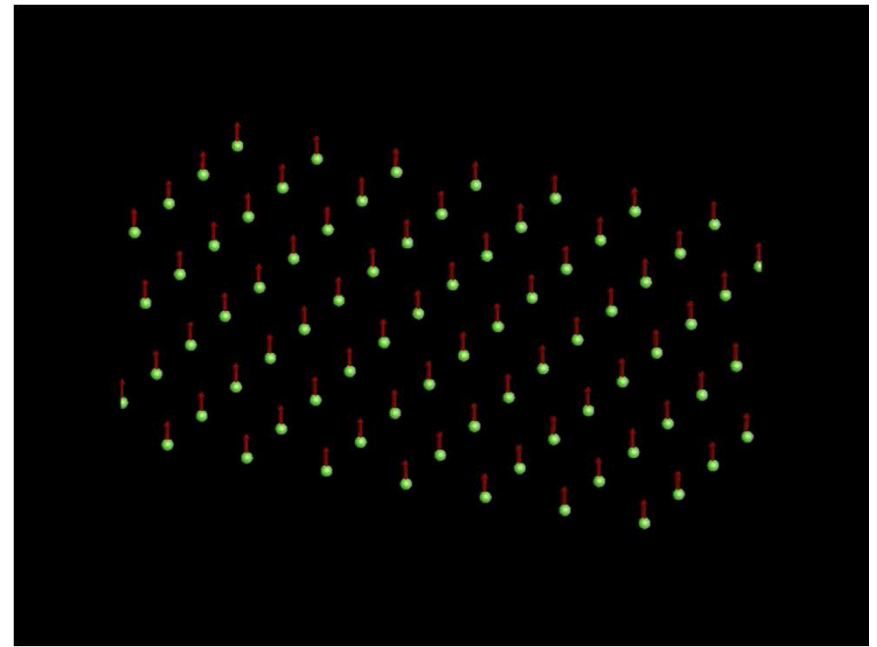
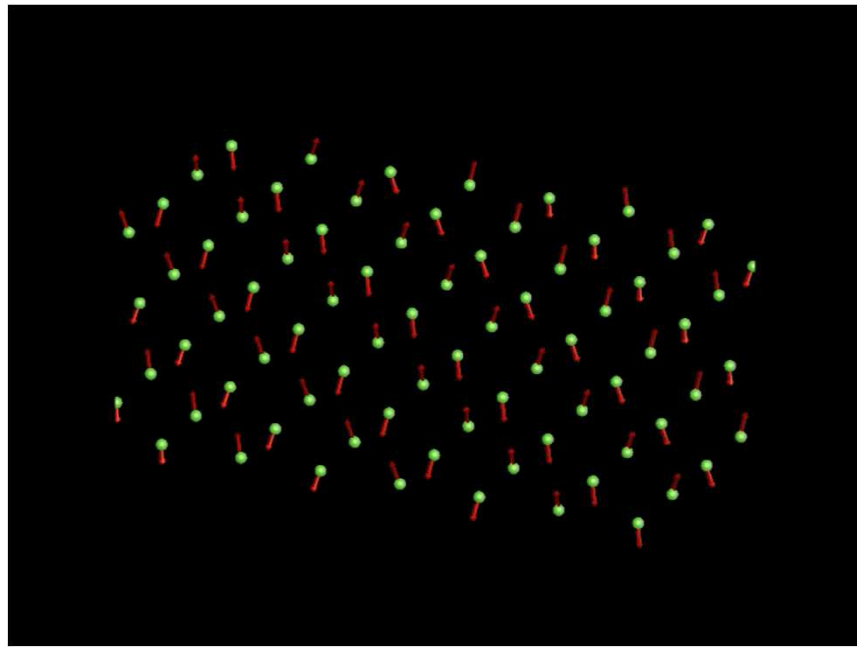
Since we know the distance between the sample and detectors and can measure the time for neutron to fly from the sample to the detector, we then can calculate the energy of scattered neutrons.

The direction of momentum can be figured out from the detector angle.

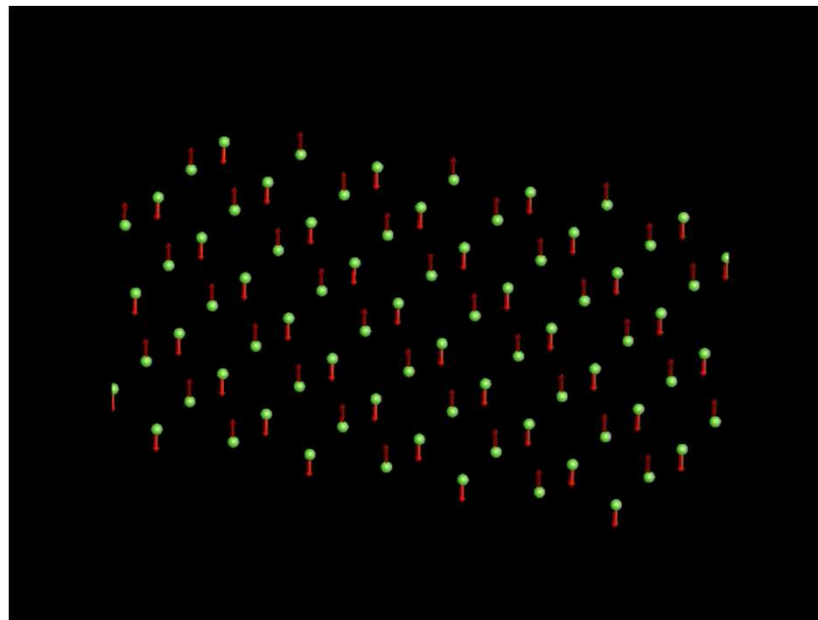
Time-of-flight method is powerful in that it can map out the huge Q-E space at a time.

# Spin waves

- Propagating disturbances in the ordering of magnetic materials. These low-lying collective excitations occur in magnetic lattices with continuous symmetry.
- Spin waves are the analog for magnetically ordered systems of lattice waves in solid systems; and just as a quantized lattice wave is called a “phonon”, a quantized spinwave is called a “magnon”.



**Animations are  
Created by IDL  
code written by  
Prof. Seunghun Lee**



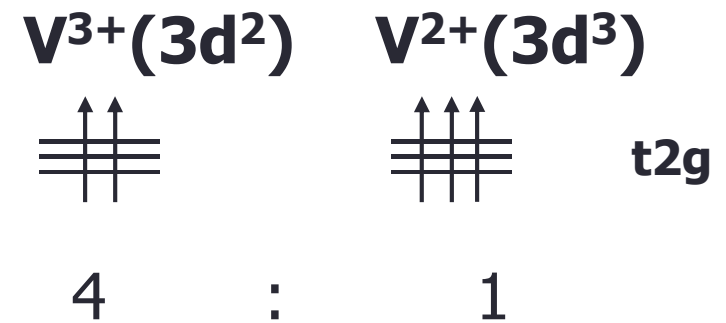
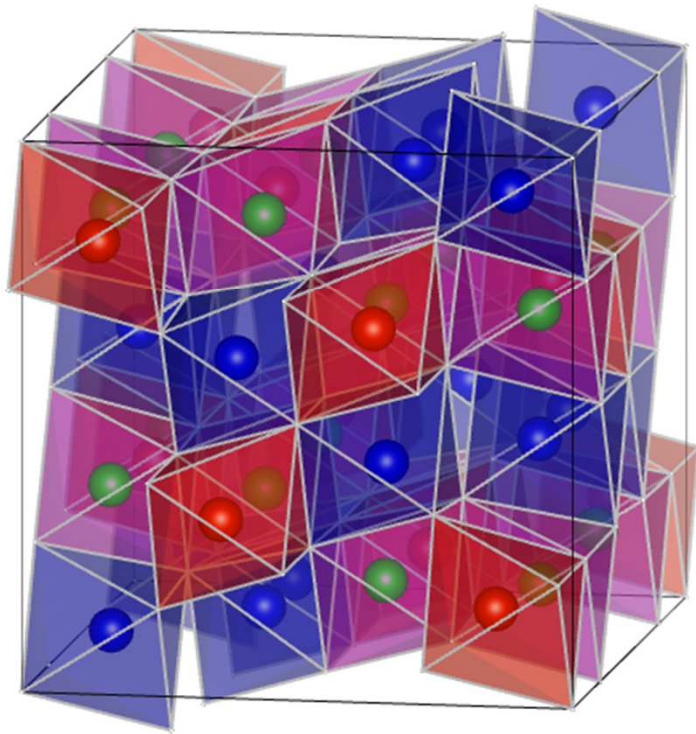


# Frustrated magnet $\text{BaV}_{10}\text{O}_{15}$

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# Introduction – $\text{BaV}_{10}\text{O}_{15}$



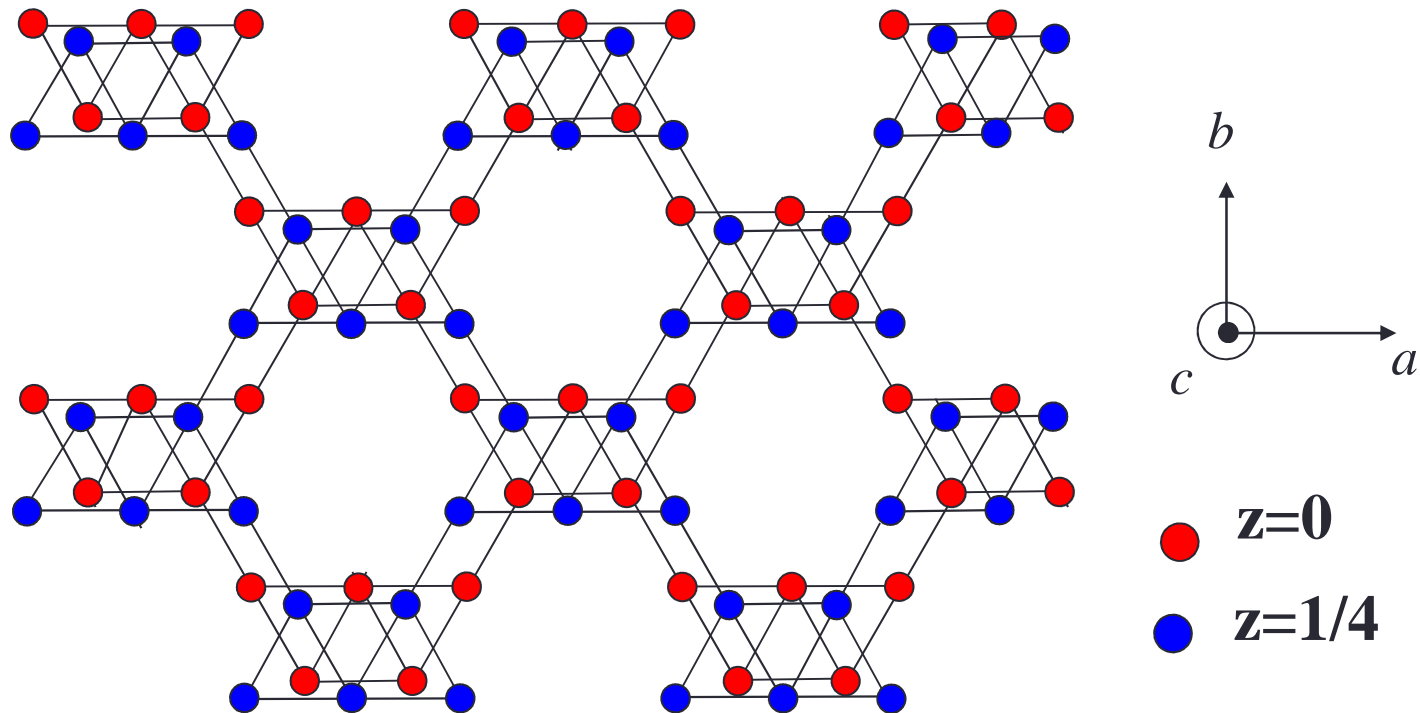
V ions surrounded by oxygen in octahedral fashion

face-, edge-, corner- sharing

mixed valence

t2g orbital ground state

# Background : V structure



$z=0$ : V5 boat

$z=1/4$ :  $180^\circ$  rotation

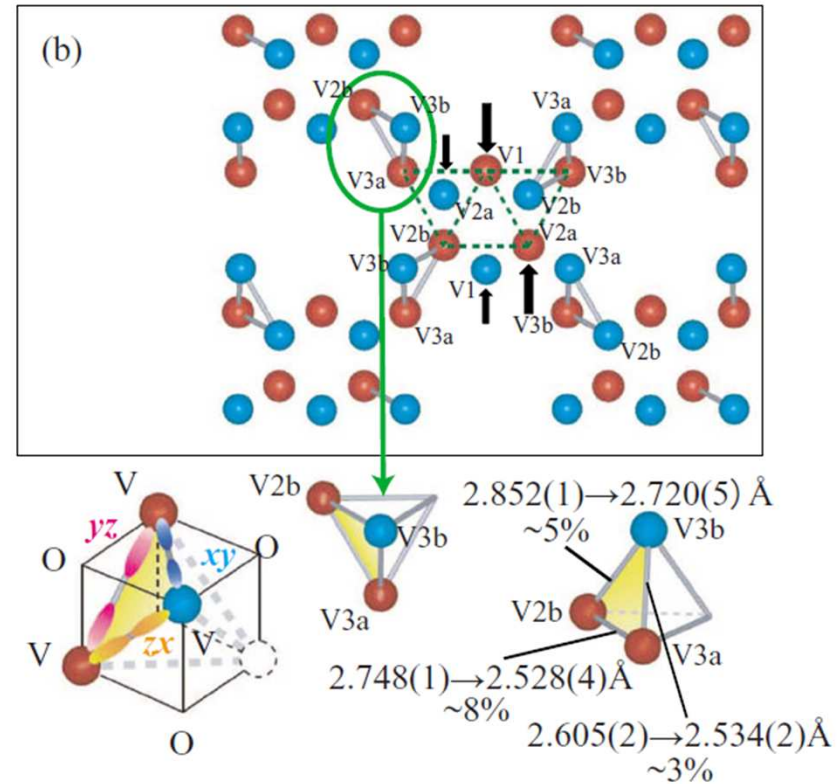
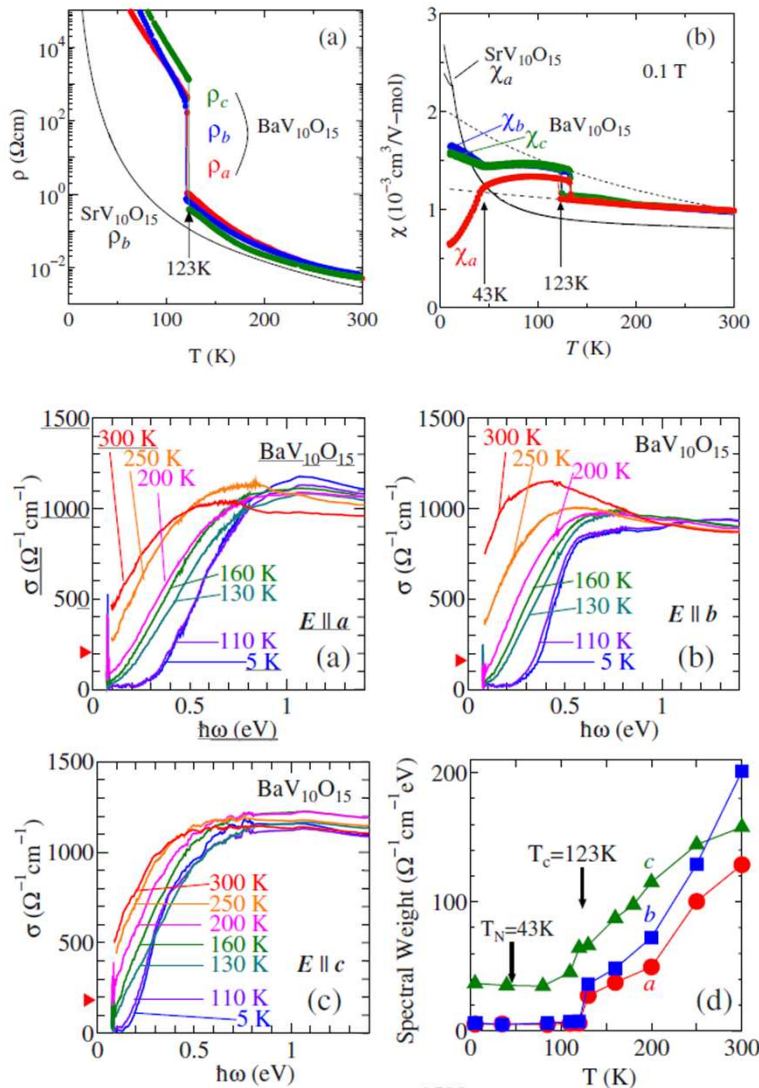
C.Bridges et. al., PRB **74**, 024426 (2006)

J.Miyazaki et. al., PRB **79** 180410(R)

(2009)

due to small overlap between  $z=0$  and  $z=-1/4$ ,  
considered as bi-layer structure stacked in ABAB fashion with B shifted along  $a$

# Bulk Properties



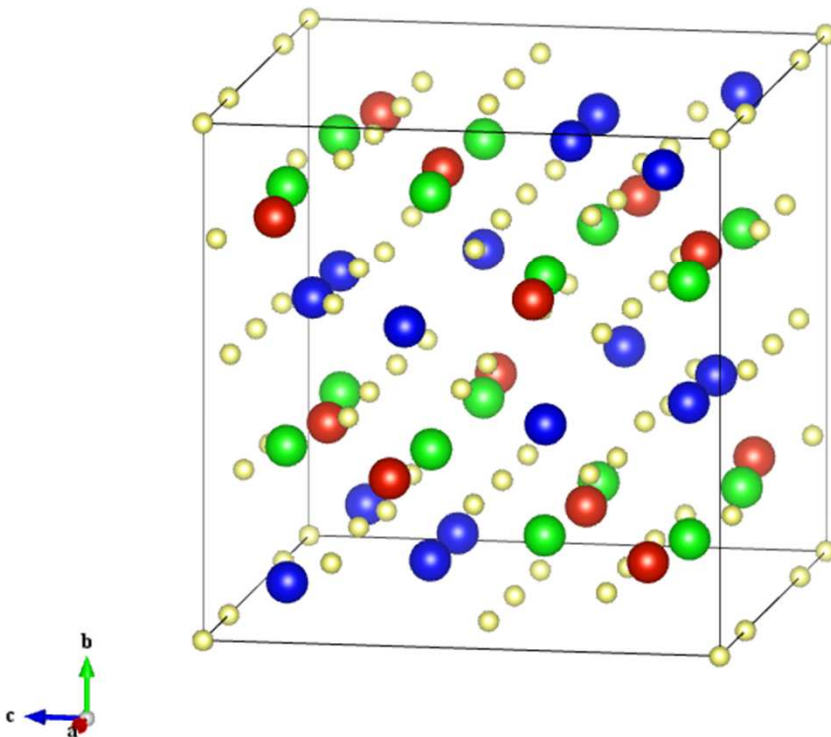
- structural phase transition at 123 K ( $\text{Cmce} \rightarrow \text{Pbca}$ )
- opening of charge gap
- magnetic phase is unknown

# Strength of interaction

- **which connection is strong?**
  - depends on the amount of orbital overlap
- **In octahedrally coordinated 3d ions case**
  - corner < edge < face

cf.) SCGO, S.-H. Lee et. al., PRL (1996)

- \*Green ions: V3a(light), V3b(dark)
- \*Yellow ions: Oxygen



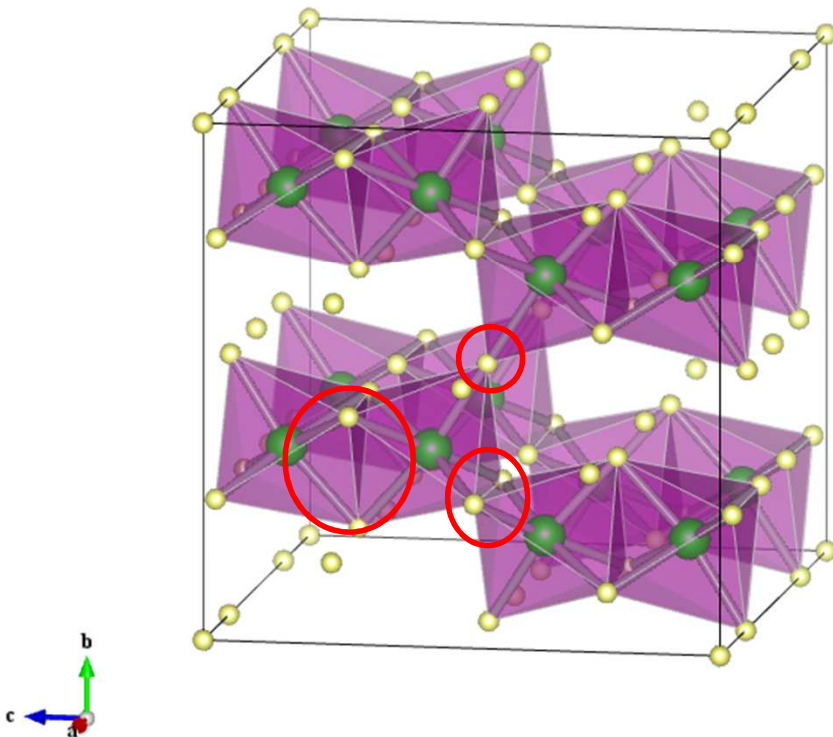
example of V3 ions interaction

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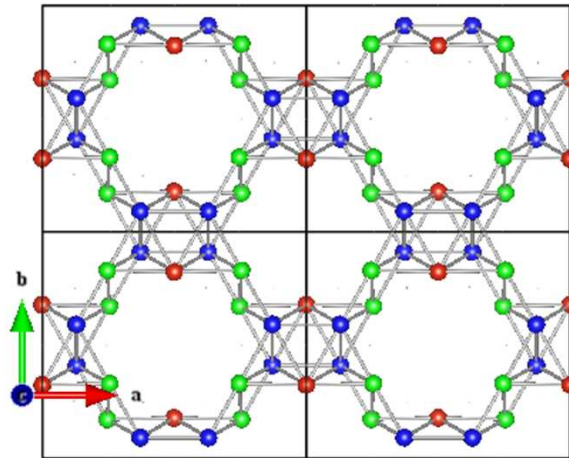


example of V3 ions interaction

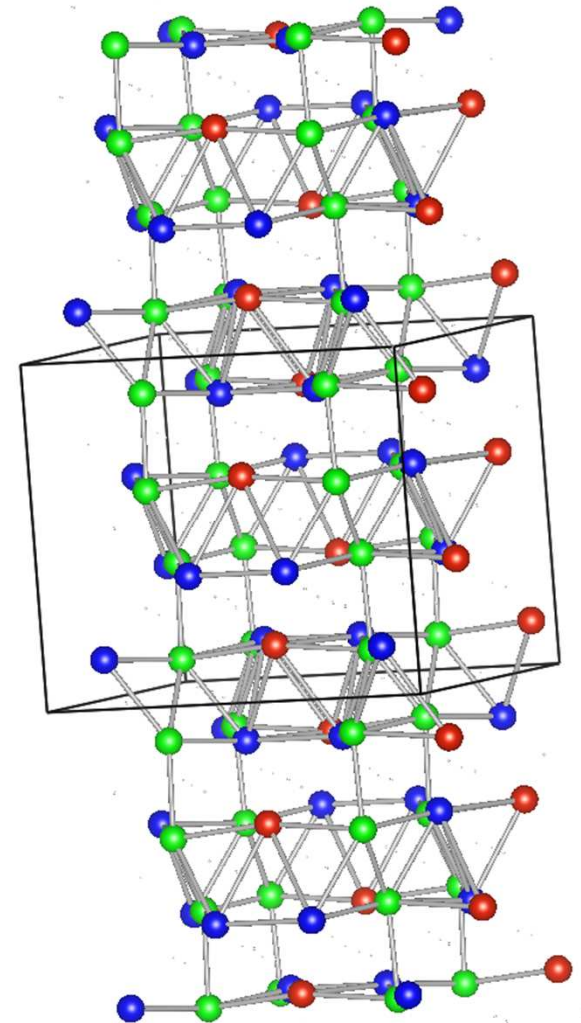
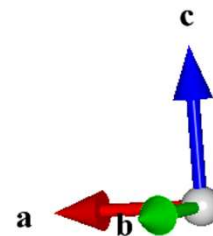
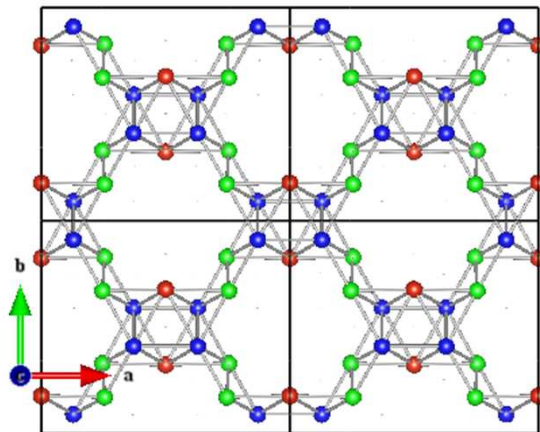


# Interactions (face and edge sharing bonds)

$0.3 < z < 0.7$

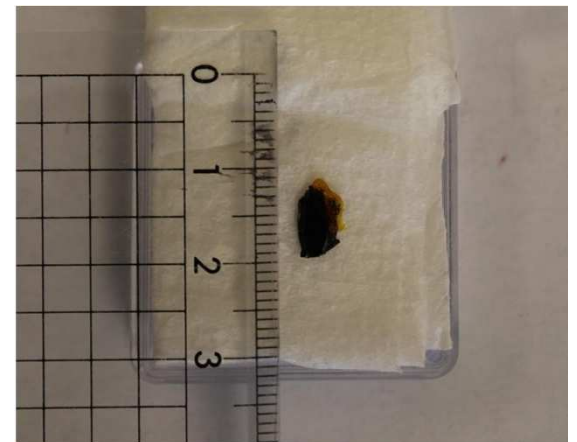
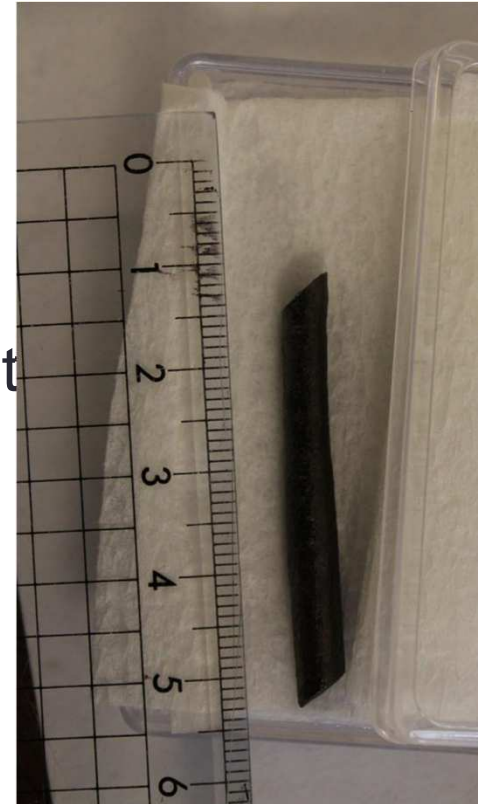


$-0.2 < z < 0.3$



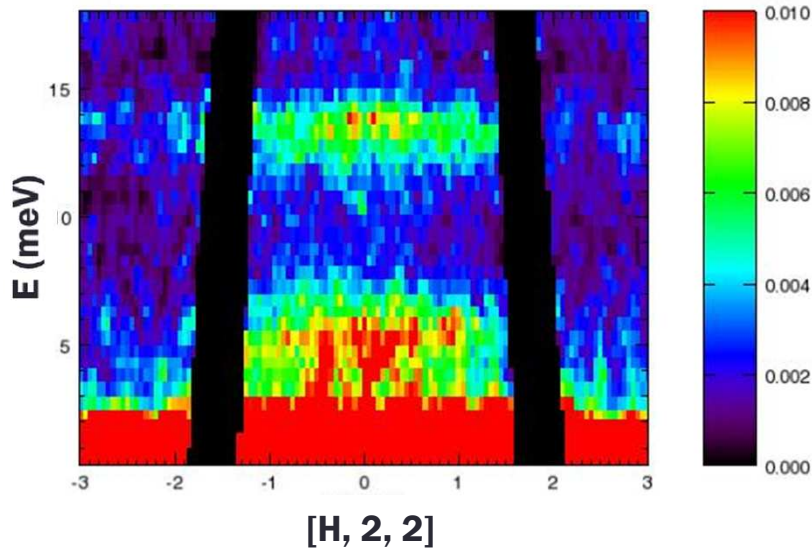
# Experiments

- ARCS, SNS
  - Time of Flight neutron scattering measurement
  - $T=4\text{K}$  with  $E_i=65\text{meV}$
  - $k_i/c$  with rotation angle  $-10$  to  $40$  deg.
- HB2A, HFIR
  - Neutron powder diffractometer
  - $T=4\text{K}$  with Wavelength  $1.538 \text{ \AA}$
- RESI, FRM II
  - Single Crystal Diffractometer
  - $T=4\text{K}$  with wavelength  $1 \text{ \AA}$

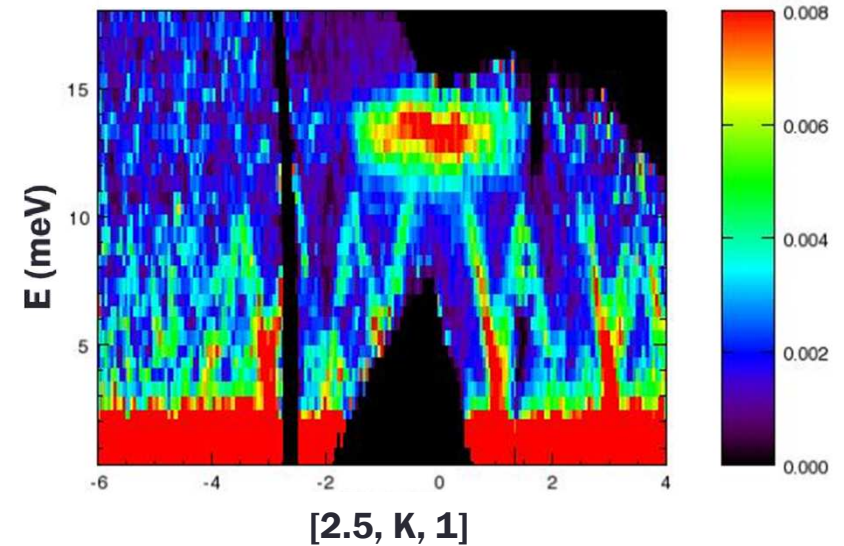


# Magnetic excitations

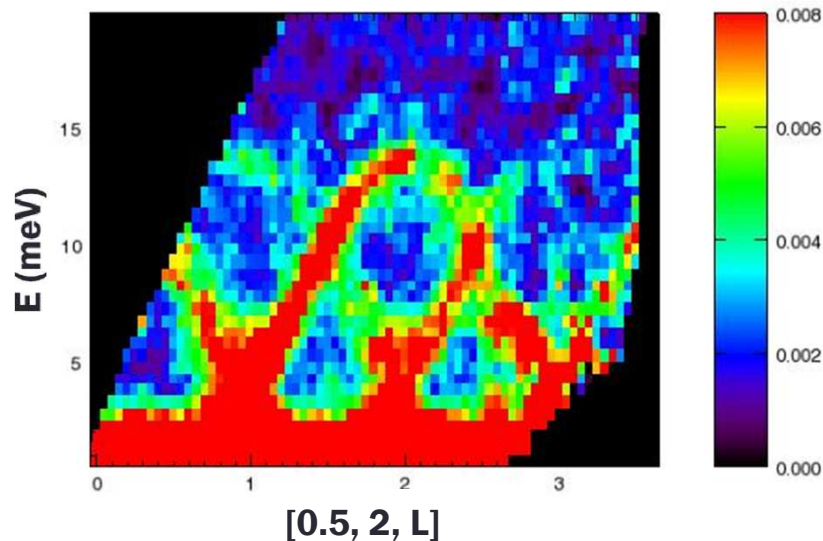
EH cut ( $K=[1.8, 2.2]$ ,  $L=[1.8, 2.2]$ )



EK cut ( $H=[2.3, 2.7]$ ,  $L=[0.8, 1.2]$ )



EL cut ( $H=[0.3, 0.7]$ ,  $K=[1.8, 2.2]$ )



**H:** Band 1: almost dispersionless, 5~6 meV  
Band 2: flat, ~13 meV

**K:** Band 1: highly dispersive, 0-10 meV  
Band 2: flat, ~13 meV

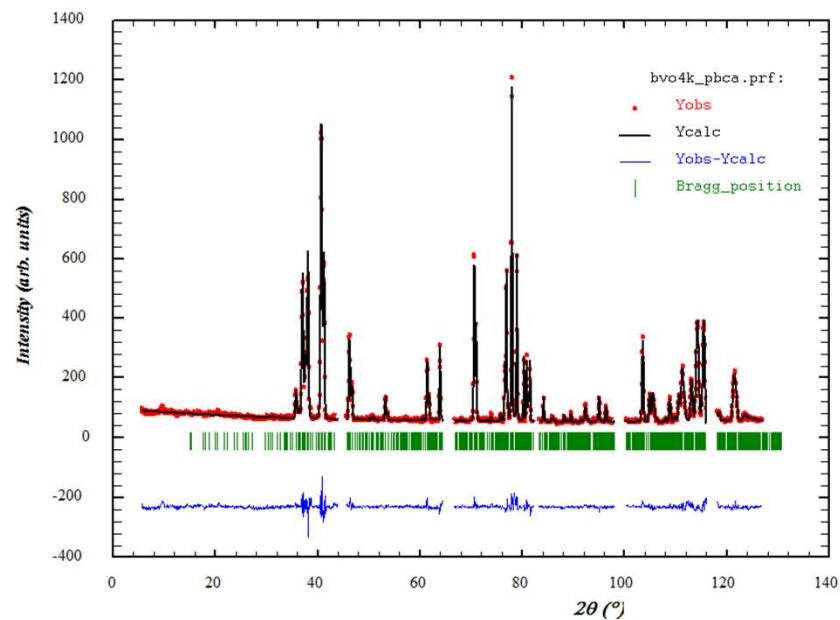
**L:** Band 1: 0-13 meV, highly dispersive

**-not consistent with bi-layer stacked along c**



# Powder Diffraction

BVO\_4K\_Pbca CELL: 11.61890 9.85910 9.42221 90.0000 90.0000 90.0000 SPGR: Pb



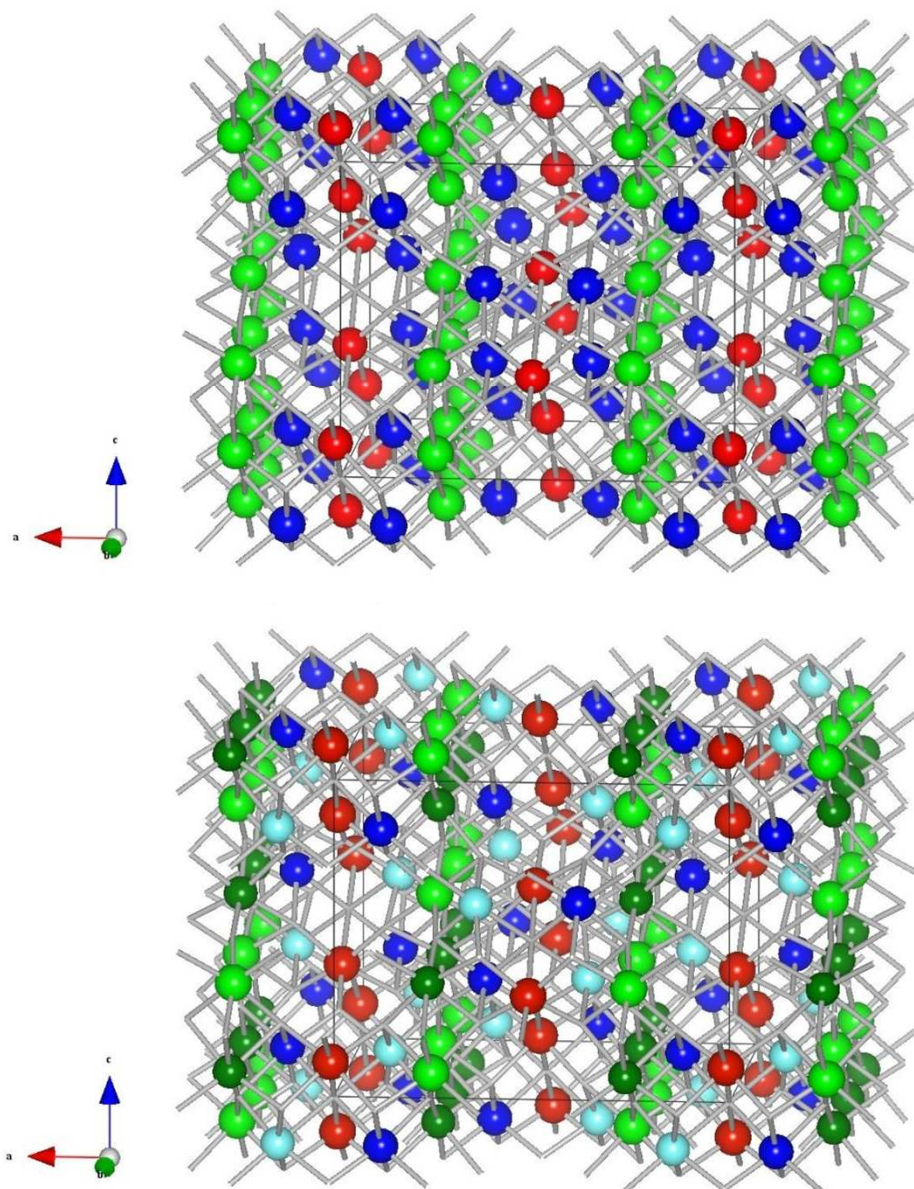
Cmca  $\rightarrow$  Pbca

V1(Red) $\rightarrow$ V1 (Red)

V2(Blue) $\rightarrow$ V2a, V2b(Blue, Sky)

V3(Green) $\rightarrow$ V3a, V3b(Green, Dark Green)

:mirror symmetry is lost



# Single Crystal Data

Magnetic peaks					
H	K	L	Q	Observed	Sigma
0.5	-1	0	0.692294	10098.89	252.9
0.5	0	1	0.719572	-177.64	282.46
-0.5	1	1	0.9612	503.91	233.4
0.5	1	1	0.9612	406.75	195.22
1.5	1	0	1.031545	7367.78	137.95
-0.5	0	2	1.360842	2930.91	207.17
-0.5	-2	1	1.463726	15696.28	286.76
0.5	-2	1	1.463726	17409.14	284.03
-0.5	1	2	1.502641	456.01	189.4
0.5	1	2	1.502641	83.16	213.6
1.5	-2	0	1.51084	647.57	205.6
1.5	0	2	1.560982	694.32	226.38
2.5	1	1	1.636639	1084.05	235.15
-1.5	2	1	1.651446	2442.92	206.92
1.5	2	1	1.651446	1989.52	127.95
-0.5	2	2	1.864512	4621.35	189.72
0.5	2	2	1.864512	3919.91	189.21
-1.5	3	1	2.181297	2053.82	242.04
1.5	-3	1	2.181297	2055.34	208.58
-1.5	1	3	2.250818	446.08	205.26
1.5	1	3	2.250818	573.47	234.38
0.5	3	2	2.346745	-406.49	187.29

RESI, FRM II

-Single Crystal Diffractometer

--T= 4K with wavelength 1 Å



# Representational Analysis

$$\Gamma_{mag} = \sum_{\Gamma_i} n_i \Gamma_i \quad \Gamma_i - \text{Irreducible representation}$$

$$\mathbf{s}^{kv} = \sum_L \sum_{\lambda} C_{\lambda}^{kLv} \psi_{\lambda}^{kLv}$$

Spin of the  $n^{\text{th}}$  cell  $S_{ni} = \exp(ikt_n) S_{oi}$

# Representational Analysis

$$\Gamma_{mag} = \sum_{\Gamma_i} n_i \Gamma_i$$

$\Gamma_i$  – Irreducible representation

Magnetic structure

$$m_j = \sum_L \sum_{\lambda} C_{\lambda}^{kLv} \psi_{\lambda}^{kLv} e^{-2\pi i \mathbf{k} \cdot \mathbf{r}}$$

**The propagation vector**

$\vec{V}_j = \vec{V}_i \exp(-2\pi i \vec{k} \cdot \vec{r})$

$\vec{V}_j = \vec{V}_i \exp \left[ -2\pi i \begin{pmatrix} 0 \\ 0 \\ 0.5 \end{pmatrix} \cdot \begin{pmatrix} 0 \\ 0 \\ 3 \end{pmatrix} \right] = \vec{V}_i \exp[-3\pi i] = -\vec{V}_i$

$\vec{V}_j = \vec{V}_i \exp \left[ -2\pi i \begin{pmatrix} 0 \\ 0 \\ 0.5 \end{pmatrix} \cdot \begin{pmatrix} 0 \\ 0 \\ 2 \end{pmatrix} \right] = \vec{V}_i \exp[-2\pi i] = \vec{V}_i$

$\vec{V}_j = \vec{V}_i \exp \left[ -2\pi i \begin{pmatrix} 0 \\ 0 \\ 0.5 \end{pmatrix} \cdot \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} \right] = \vec{V}_i \exp[-\pi i] = -\vec{V}_i$

$\vec{V}_j = \vec{V}_i \exp \left[ -2\pi i \begin{pmatrix} 0 \\ 0 \\ 0.5 \end{pmatrix} \cdot \begin{pmatrix} 0 \\ 0 \\ 0 \end{pmatrix} \right] = \vec{V}_i \exp[-2\pi i \cdot 0] = \vec{V}_i$



# Refining the magnetic structure - Basis functions

80 Magnetic V atoms

$$\Gamma_{mag} = 6\Gamma_1 + 6\Gamma_2$$

$\Gamma_1, \Gamma_2$ - 2 Dimensional

Table B.3: The basis functions of the irreducible representations for the space group Pbcu with  $\mathbf{k} = (1/2, 0, 0)$

BV	Atom 1			Atom 2			Atom 3			Atom 4			Atom 5			Atom 6			Atom 7			Atom 8		
	$m_x$	$m_y$	$m_z$	$m_x$	$m_y$	$m_z$	$m_x$	$m_y$	$m_z$	$m_x$	$m_y$	$m_z$	$m_x$	$m_y$	$m_z$	$m_x$	$m_y$	$m_z$	$m_x$	$m_y$	$m_z$	$m_x$	$m_y$	$m_z$
$\psi_1$	1	0	0	0	0	0	1	0	0	0	0	0	-1	0	0	0	0	0	-1	0	0	0	0	0
$\psi_2$	0	1	0	0	0	0	0	-1	0	0	0	0	0	-1	0	0	0	0	0	1	0	0	0	0
$\psi_3$	0	0	1	0	0	0	0	0	1	0	0	0	0	0	-1	0	0	0	0	0	-1	0	0	0
$\psi_4$	0	0	0	1	0	0	0	0	0	1	0	0	0	0	0	-1	0	0	0	0	0	-1	0	0
$\psi_5$	0	0	0	0	1	0	0	0	0	0	-1	0	0	0	0	0	-1	0	0	0	0	0	1	0
$\psi_6$	0	0	0	0	0	-1	0	0	0	0	0	-1	0	0	0	0	0	1	0	0	0	0	0	1
$\psi_7$	0	0	0	1	0	0	0	0	0	-1	0	0	0	0	0	1	0	0	0	0	0	-1	0	0
$\psi_8$	0	0	0	0	1	0	0	0	0	0	1	0	0	0	0	0	1	0	0	0	0	0	1	0
$\psi_9$	0	0	0	0	0	-1	0	0	0	0	0	1	0	0	0	0	0	-1	0	0	0	0	0	1
$\psi_{10}$	1	0	0	0	0	0	-1	0	0	0	0	0	1	0	0	0	0	0	-1	0	0	0	0	0
$\psi_{11}$	0	1	0	0	0	0	0	1	0	0	0	0	0	1	0	0	0	0	0	1	0	0	0	0
$\psi_{12}$	0	0	1	0	0	0	0	0	-1	0	0	0	0	0	1	0	0	0	0	0	-1	0	0	0
$\psi_{13}$	1	0	0	0	0	0	1	0	0	0	0	0	1	0	0	0	0	0	1	0	0	0	0	0
$\psi_{14}$	0	1	0	0	0	0	0	-1	0	0	0	0	0	1	0	0	0	0	0	-1	0	0	0	0
$\psi_{15}$	0	0	1	0	0	0	0	0	1	0	0	0	0	0	1	0	0	0	0	0	1	0	0	0
$\psi_{16}$	0	0	0	1	0	0	0	0	0	1	0	0	0	0	0	1	0	0	0	0	0	1	0	0
$\psi_{17}$	0	0	0	0	1	0	0	0	0	0	-1	0	0	0	0	0	1	0	0	0	0	0	-1	0
$\psi_{18}$	0	0	0	0	0	-1	0	0	0	0	0	-1	0	0	0	0	0	-1	0	0	0	0	0	-1
$\psi_{19}$	0	0	0	1	0	0	0	0	0	-1	0	0	0	0	0	-1	0	0	0	0	0	1	0	0
$\psi_{20}$	0	0	0	0	1	0	0	0	0	0	1	0	0	0	0	0	-1	0	0	0	0	0	-1	0
$\psi_{21}$	0	0	0	0	0	-1	0	0	0	0	0	1	0	0	0	0	0	1	0	0	0	0	0	-1
$\psi_{22}$	1	0	0	0	0	0	-1	0	0	0	0	0	-1	0	0	0	0	0	1	0	0	0	0	0
$\psi_{23}$	0	1	0	0	0	0	0	1	0	0	0	0	0	-1	0	0	0	0	0	-1	0	0	0	0
$\psi_{24}$	0	0	1	0	0	0	0	0	-1	0	0	0	0	0	-1	0	0	0	0	0	1	0	0	0

- C. Bridges Thesis 2002

## Refining the magnetic structure

- Refining the Magnetic structure by carrying out all the possibilities is very hard, due to enormous possibilities
- We approached the refinement in two ways.

# Refining the magnetic structure

## 1. Trimer with tetramer

There are 4 basis vectors that is pointing in a-direction. BV1, BV4, BV7, BV10

a.) All bonds with in all tetramers in unit cell are satisfied

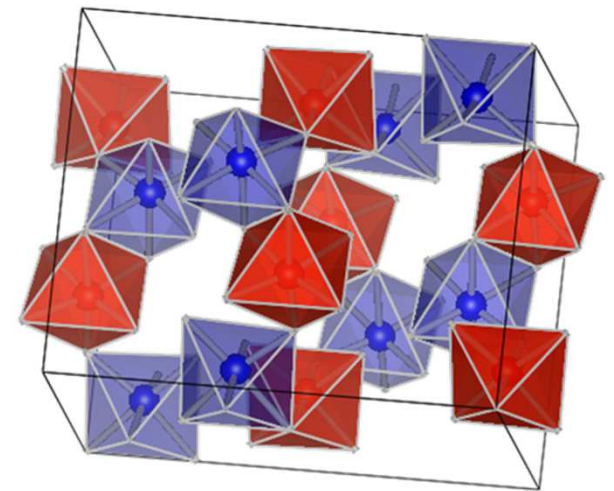
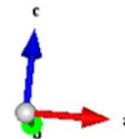
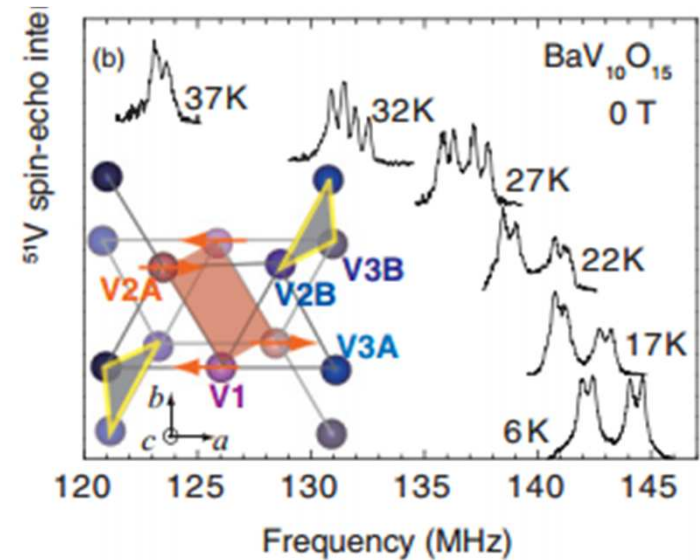
$\Gamma_1$  - BV1 and BV 10

$\Gamma_2$  - BV4 and BV 7

b.) Some bonds within tetramers does not satisfy

Other combinations of  
BV1, BV4, BV7, BV10

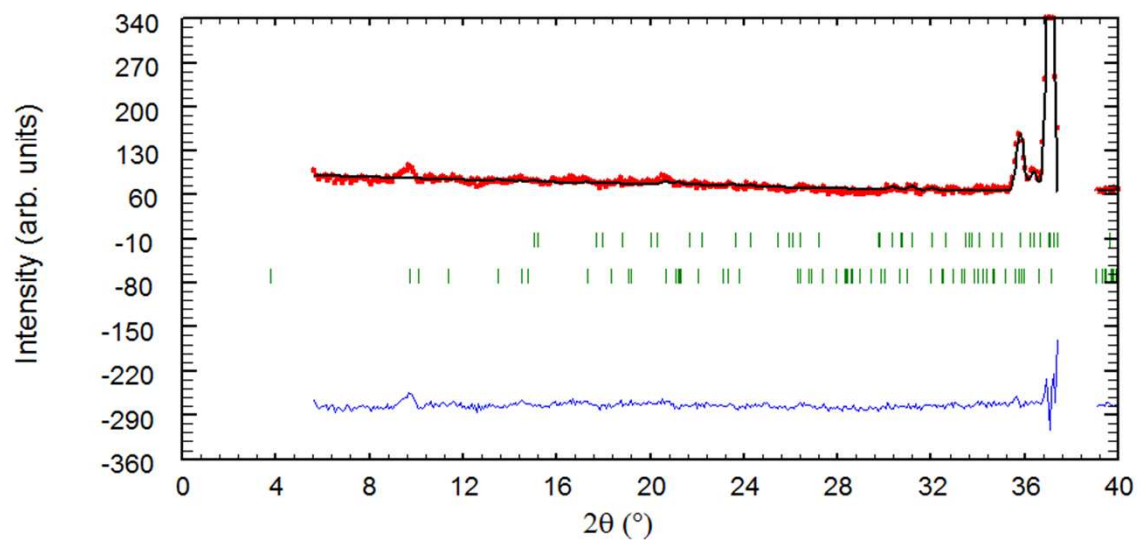
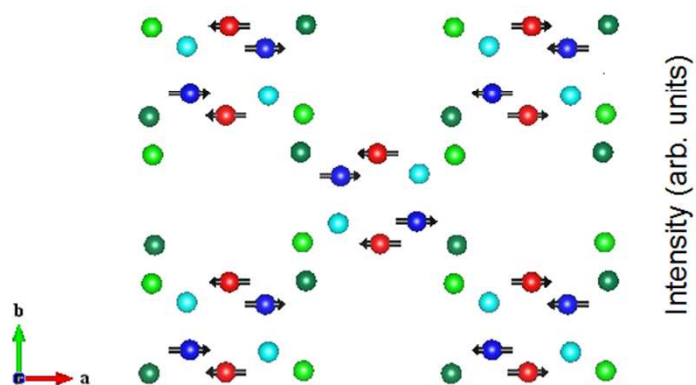
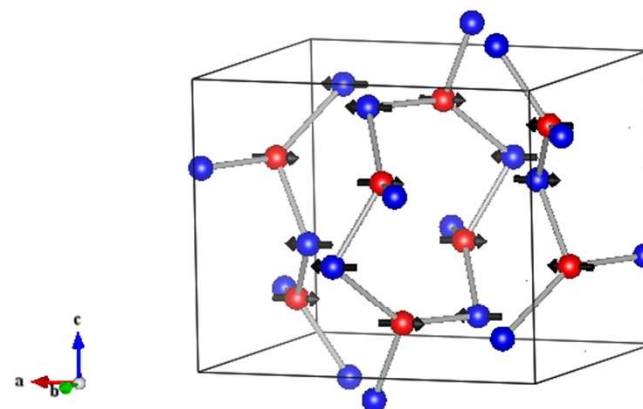
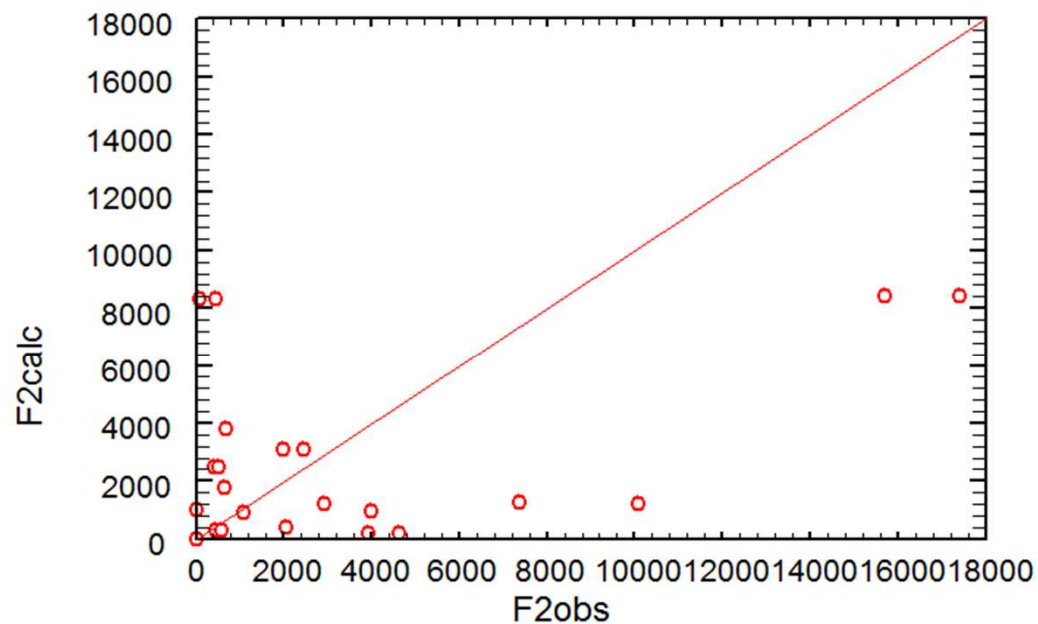
Shimizu et. al. **PRB** 84, 064421 (2011)



# Trimer with tetramer - All bonds within all tetramers in unit cell are satisfied

$\Gamma_2$

BaV10O15, lambda 1.00 RESI

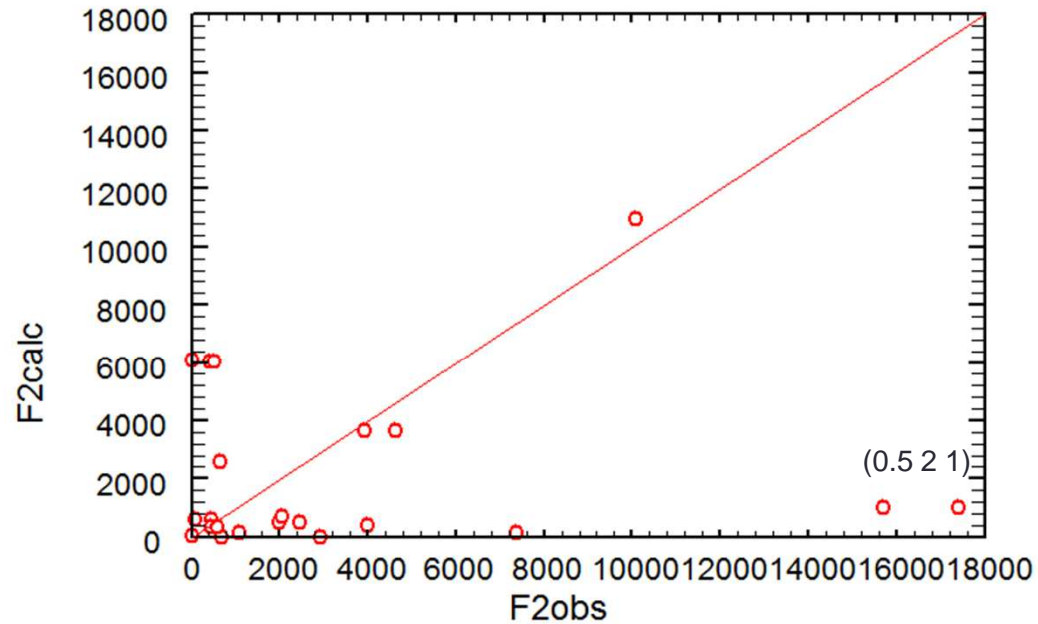




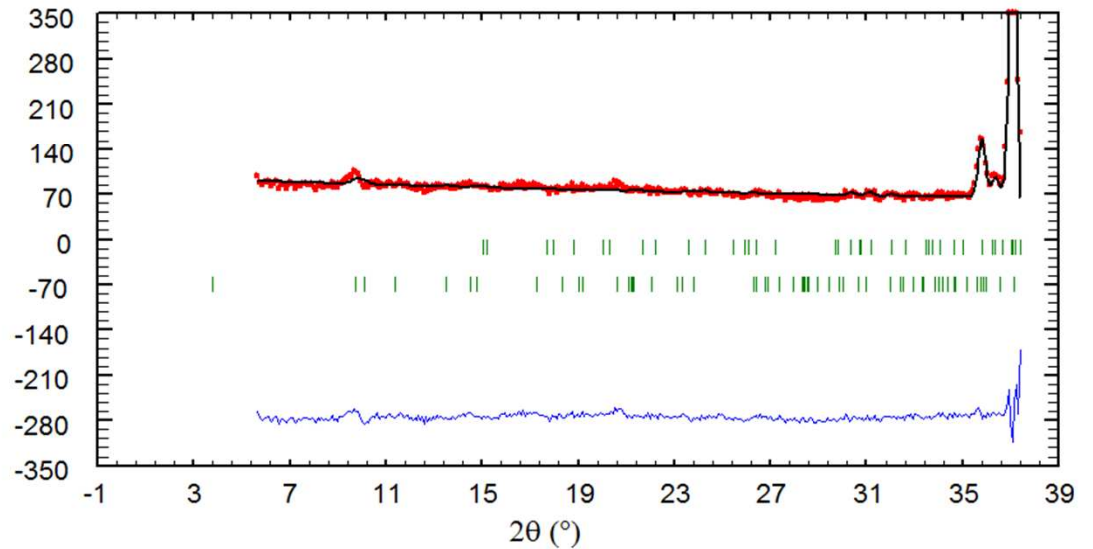
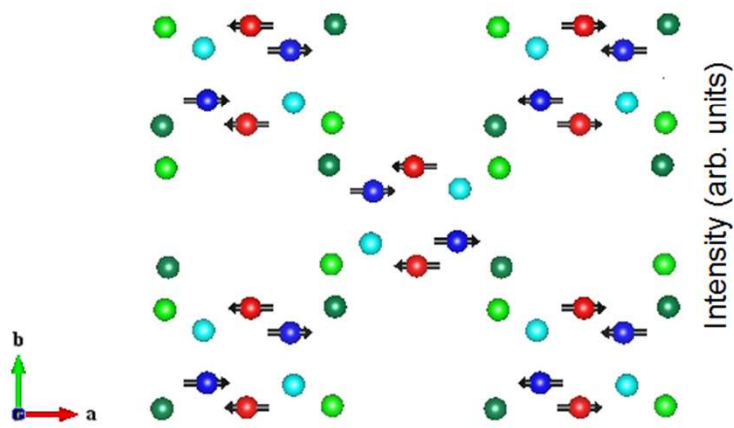
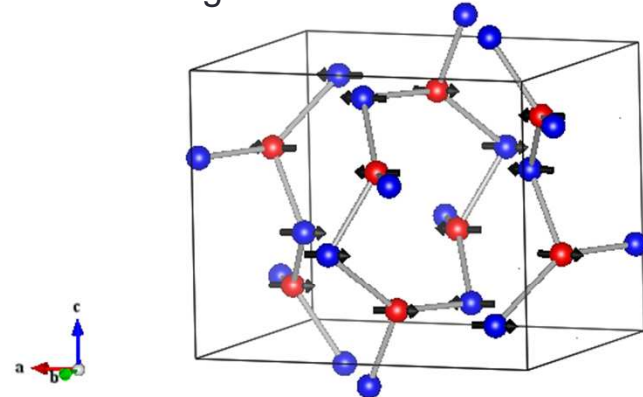
# Trimer with tetramer - All bonds within all tetramers in unit cell are satisfied

$\Gamma_1$

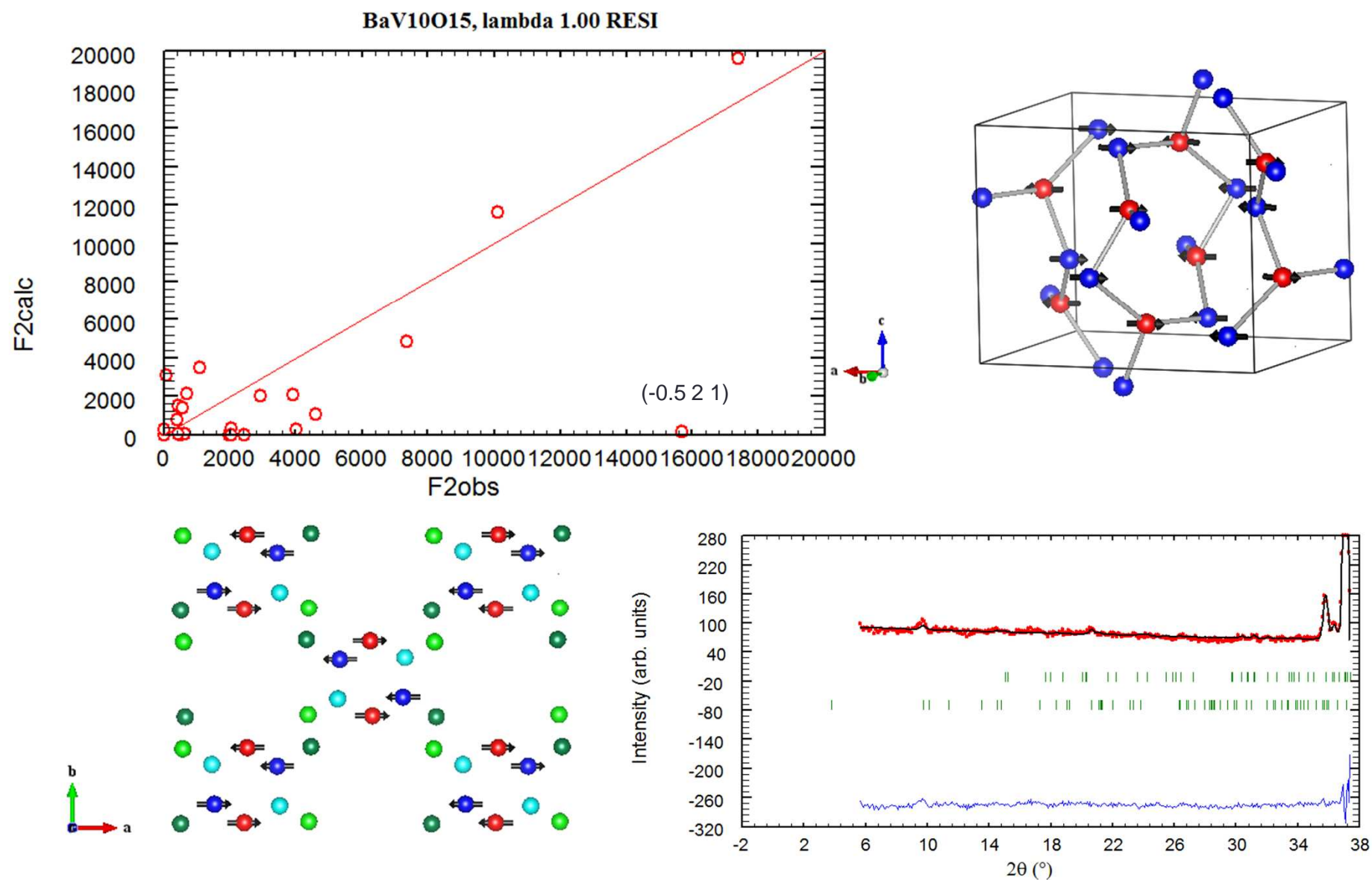
BaV10O15, lambda 1.00 RESI



Corner sharing bond along C, not satisfied or ferromagnetic



## Trimer with tetramer - Some bonds with in tetramers does not satisfy



# Refining the magnetic structure

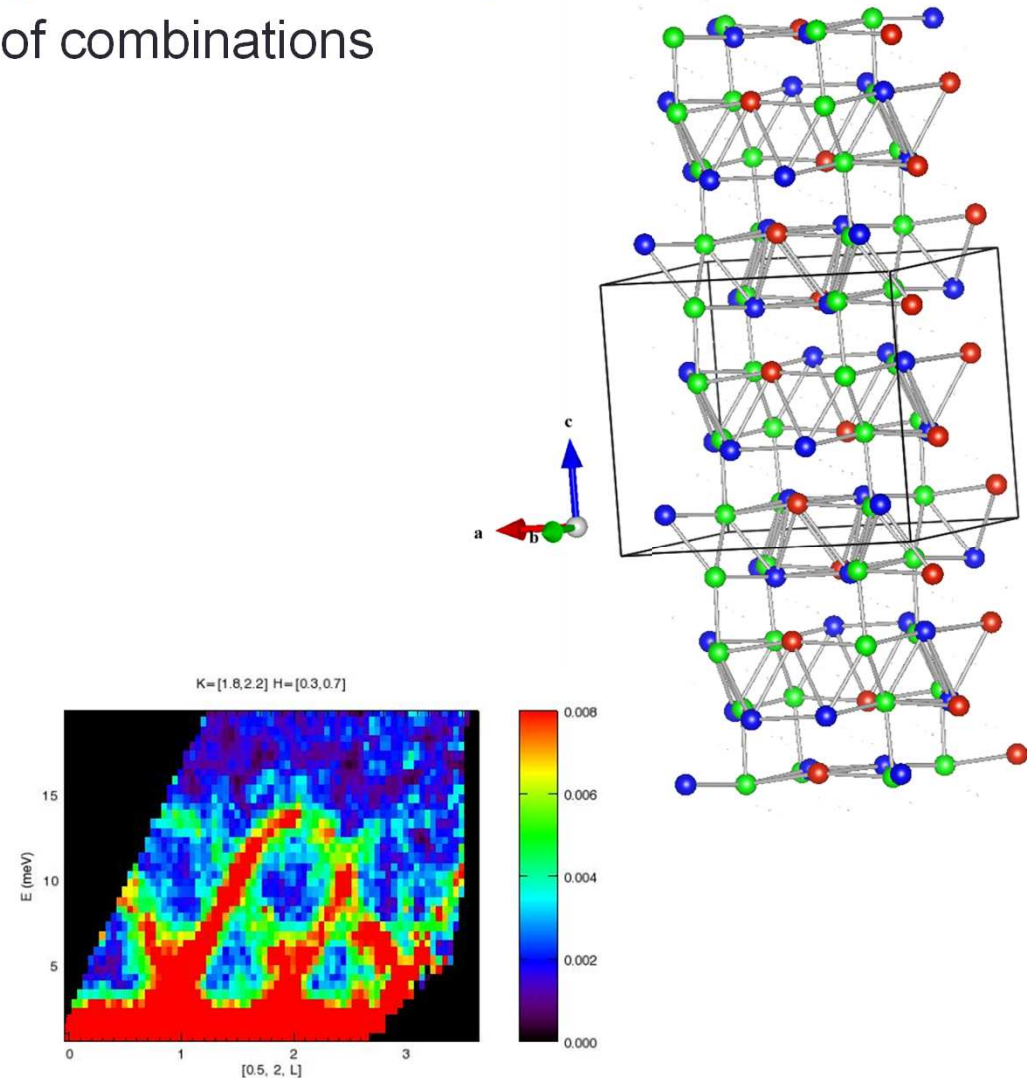
2. Non trimer model – All five V atoms are ordered

Here we assumed that chain along c-direction is always satisfied.

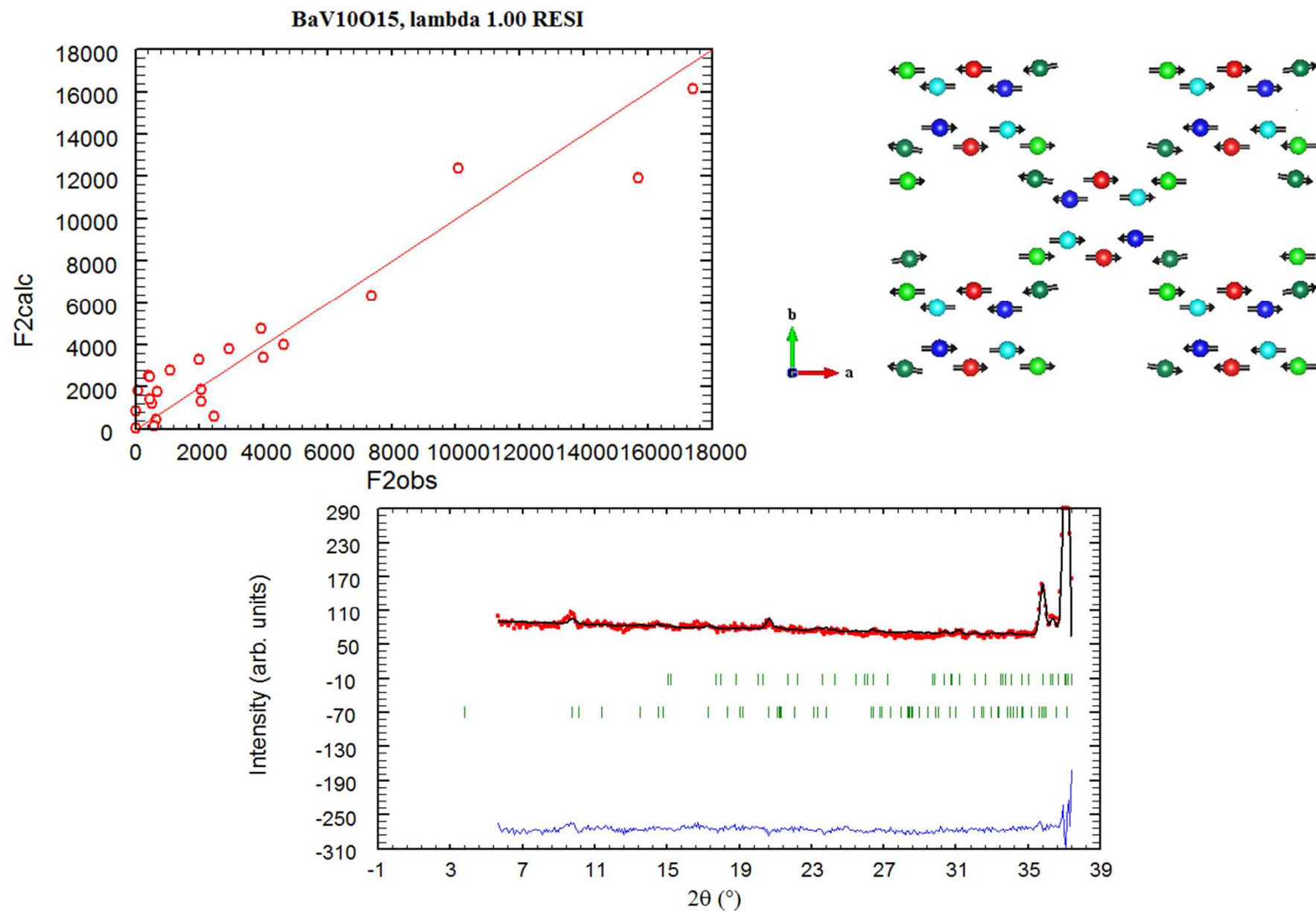
This allows us to restrict number of combinations in the basis vectors.

For eg. :

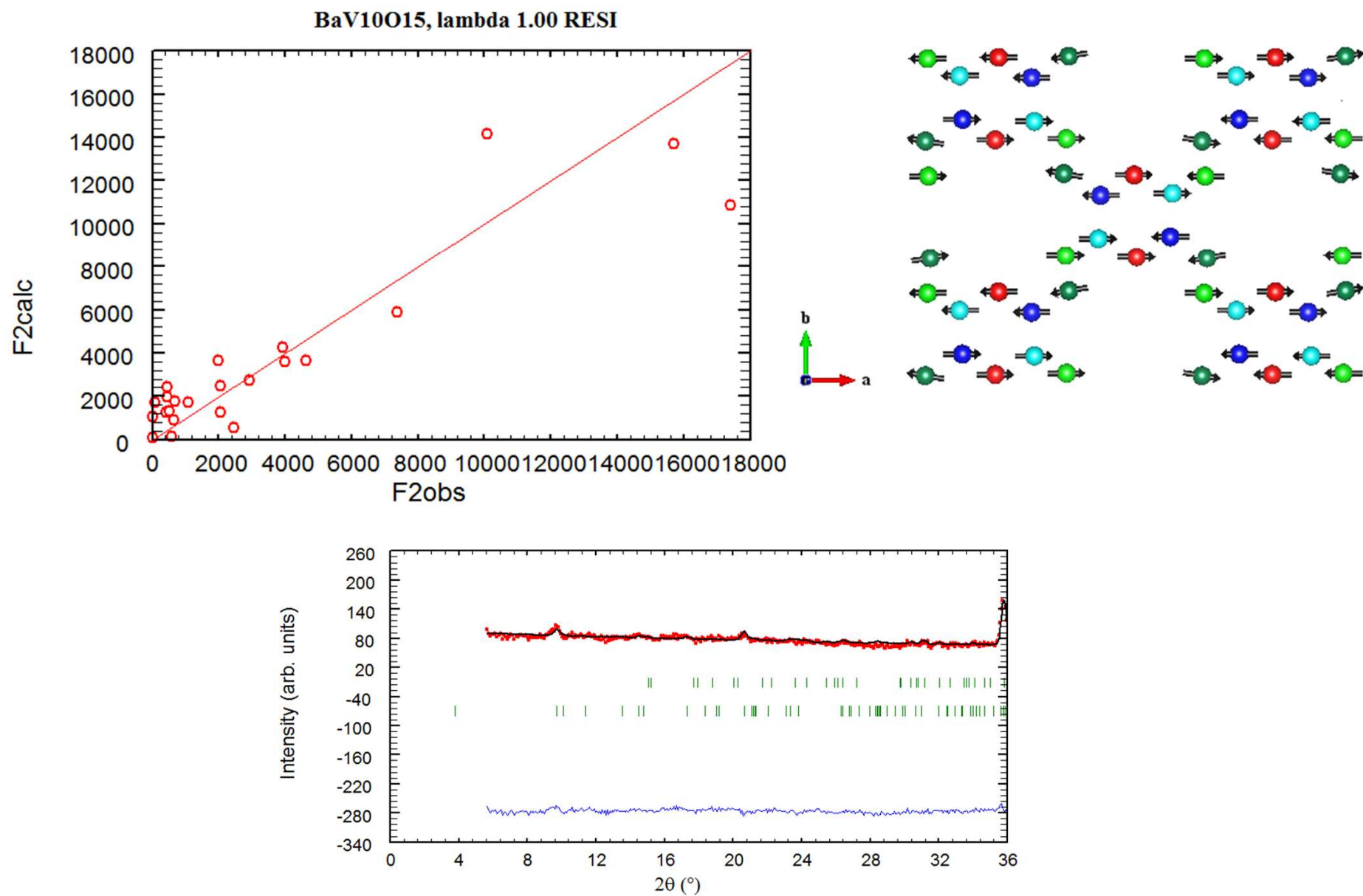
	$\Gamma_1, \Gamma_2$
a direction	V3 (BV1) + V3B(BV4) + V3 (BV7) + V3B (BV10)
b direction	V3 (BV5) + V3B (BV2) + V3 (BV11) + V3B (BV8)
c direction	V3 (BV6) + V3B (BV3) + V3 (BV12) + V3B (BV9)



## Non Trimer models – All five V atoms are ordered



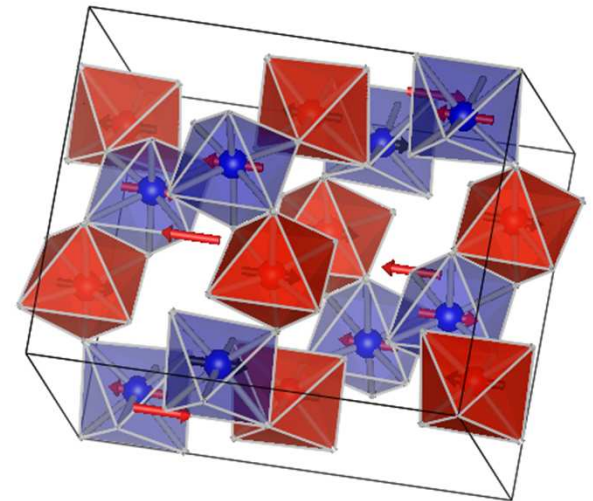
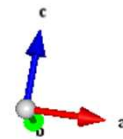
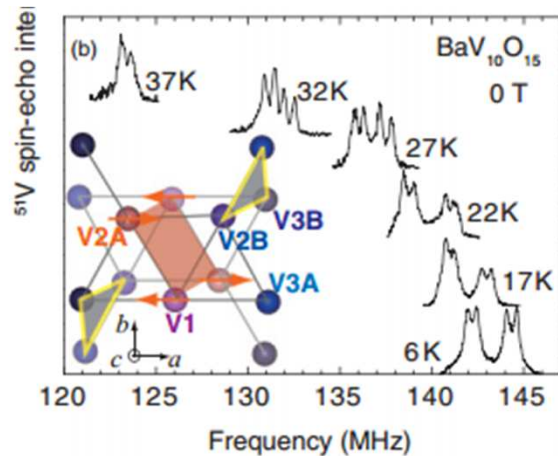
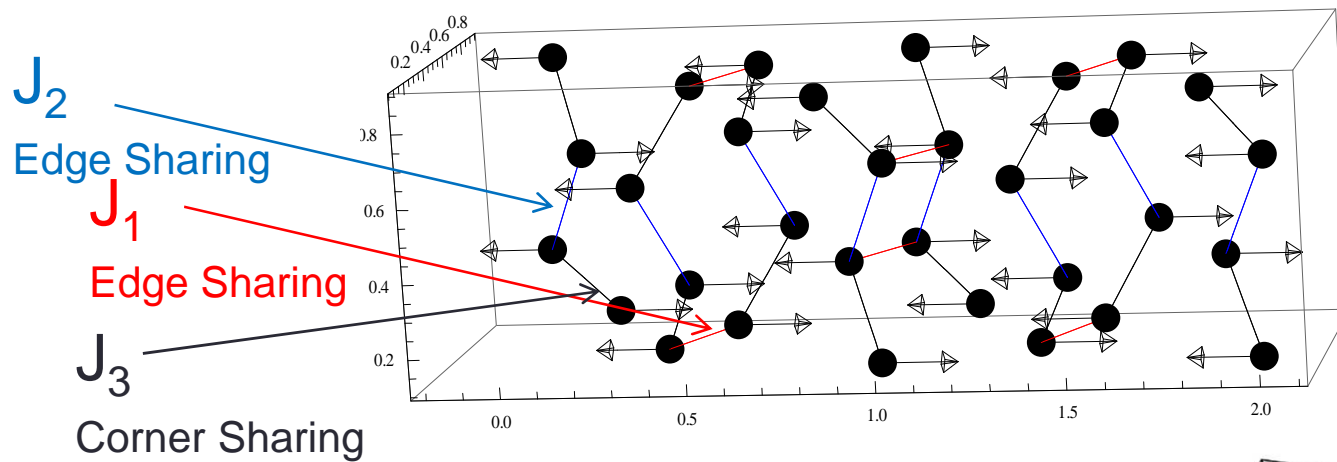
## Non Trimer models – All five V atoms are ordered



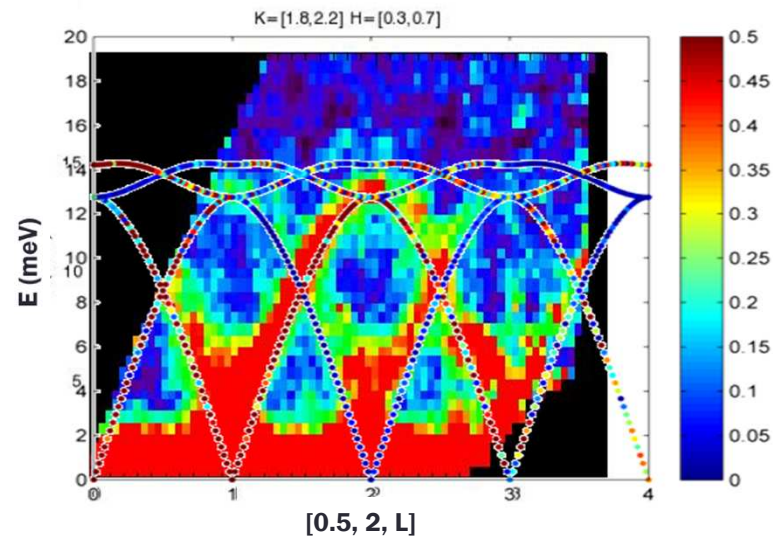
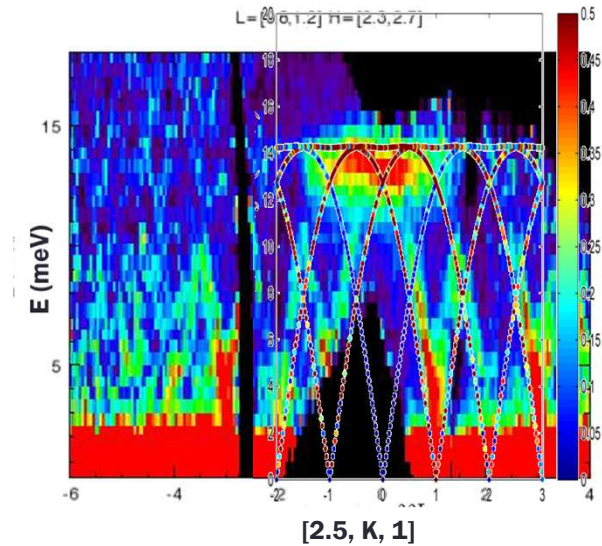
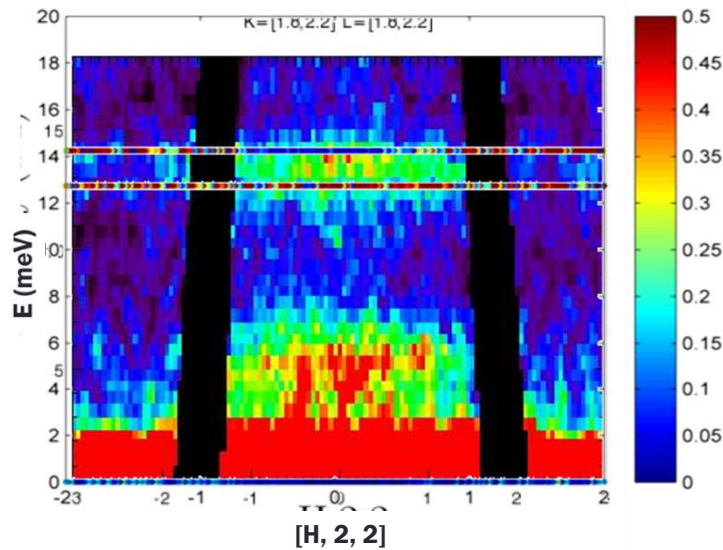


# Linear spinwave calculation - Tetramer

$$H = \sum_{i,j} J_{ij} \mathbf{S}_i \cdot \mathbf{S}_j$$

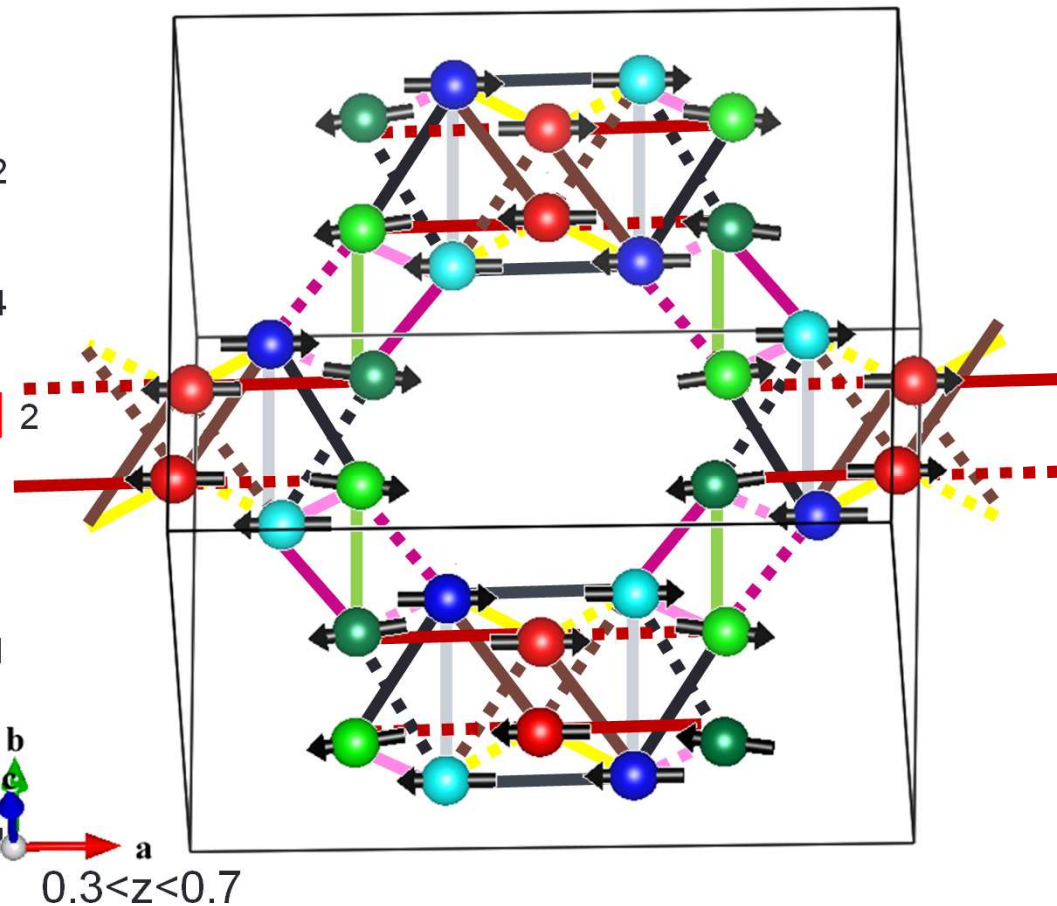
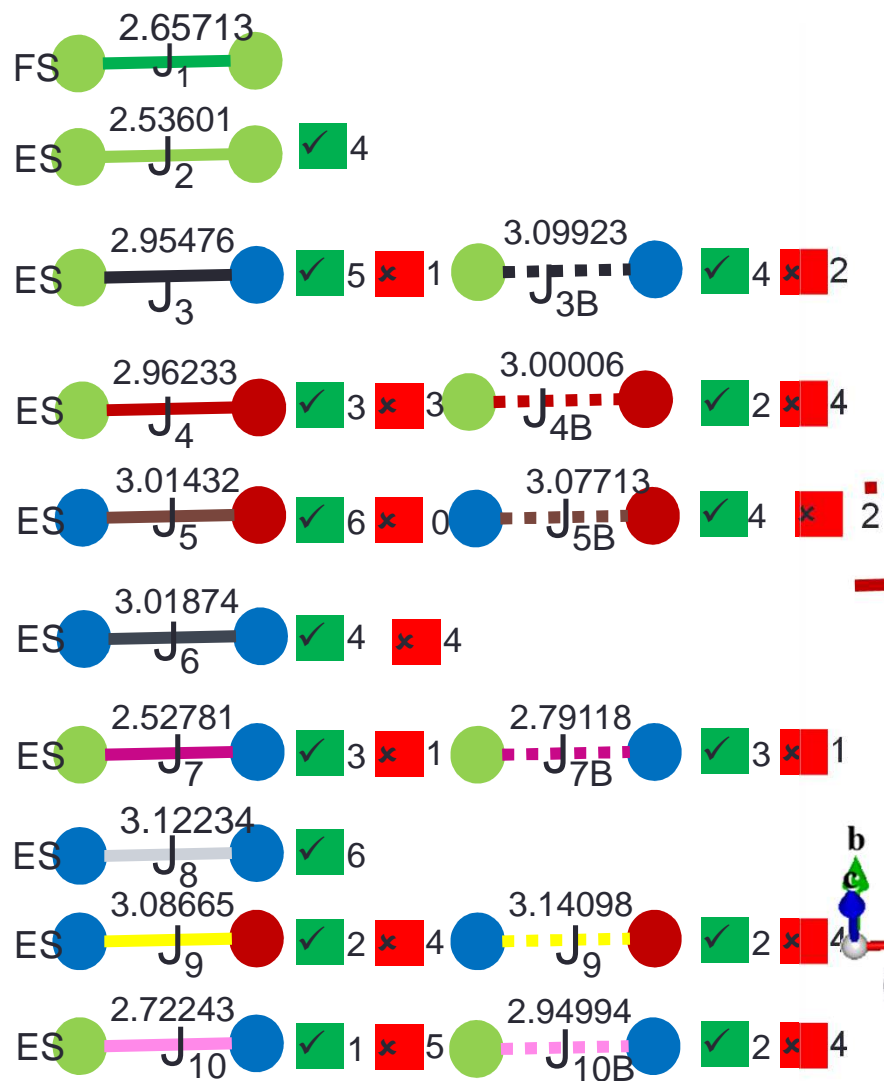


# Linear spinwave calculation - Tetramer



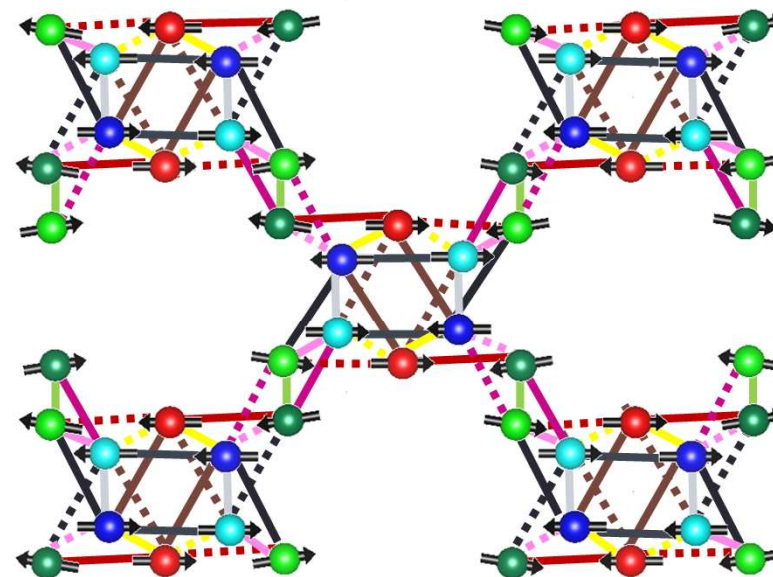
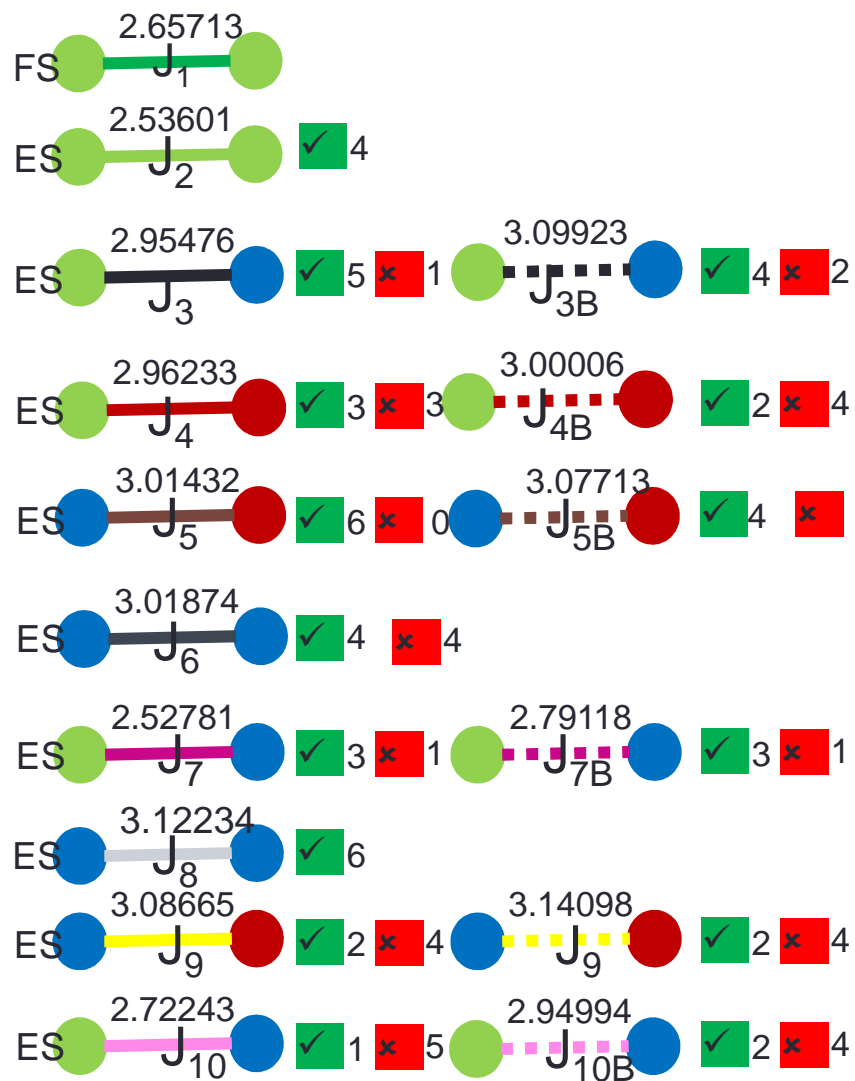
$J_2 = 7.5$  meV;  
 $J_1 = 0.3 \cdot J_2$ ;  
 $J_3 = 0.6 \cdot J_2$ ;

# Interactions – Non Trimer





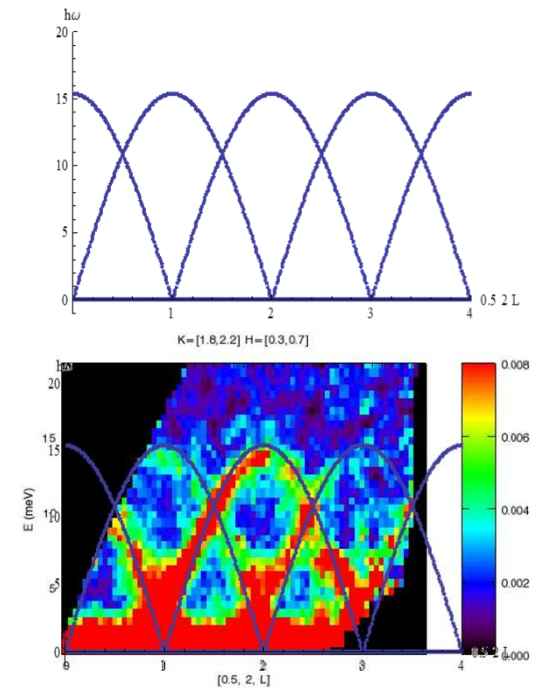
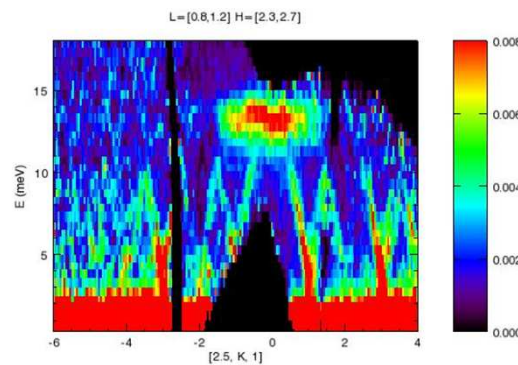
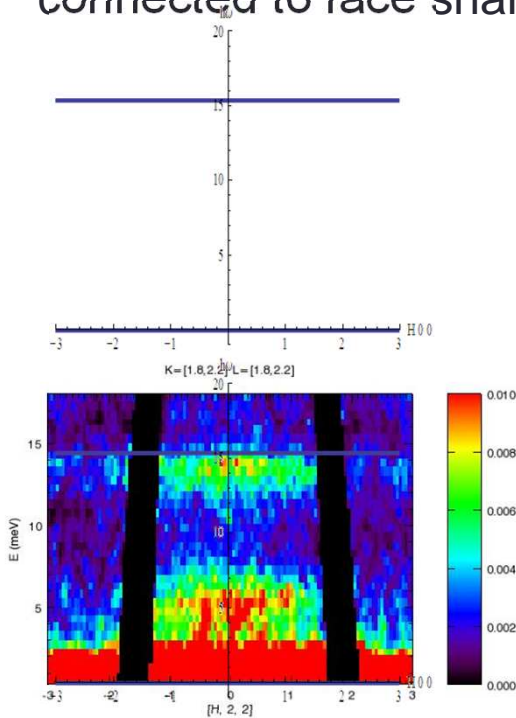
# Interactions – Non Trimer



$$-0.2 < z < 0.3$$

## Non Trimer models – All five V atoms are ordered – Spin waves

- Chain along c-direction with one face sharing and one edge sharing bonds can satisfy the L dependence of the excitations.
- Also, along a-direction it is mostly disconnected and thus can produce the flat modes observed along H.
- The reason for highly dispersive mode along K, is still a question and our hypothesis is frustrated bonds some how create a similar chain along b direction which might be connected to face sharing bond along c direction.



# Summary and future work

- We have examined the magnetic structure of BaV<sub>10</sub>O<sub>15</sub> using Single Crystal and powder diffraction data.
- Our analysis shows that Trimer model with only V1 and V2a atoms ordered, cannot reproduce our diffraction data.
- Frustrated model with all five atoms ordered can reasonably reproduce our both single crystal and powder diffraction data.
- We will continue our analysis to find the best Hamiltonian for this system.
- Our future experiments consists of,
  1. Field dependence of the excitations
  2. Inelastic experiment to see whether there is any high energy modes.



Spinel  $\text{CoV}_2\text{O}_4$

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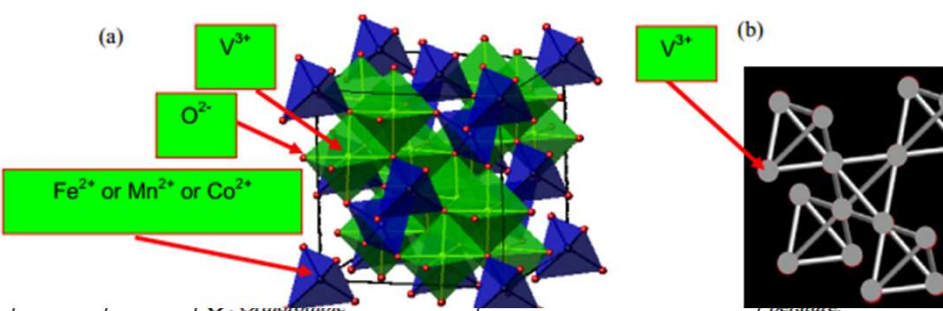
# Spinel Family

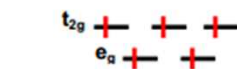
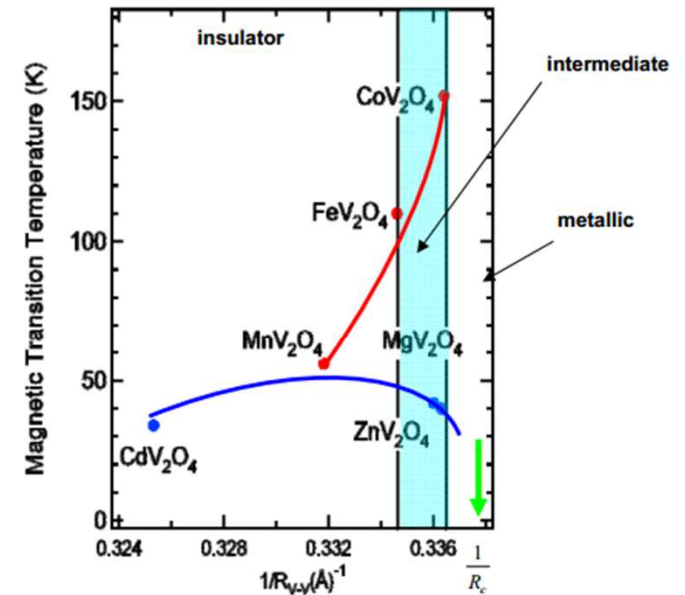
- Spinel is an important class of mixed-metal oxides, which has the general chemical composition of  $AB_2O_4$ .
- When the shape of the orbital is changed by an external stimulus, the magnetic, electric, elastic, and optical properties may also be altered.
- Spinel oxides with the general formula  $AB_2O_4$  provide a fertile playground for studying the interplay between these degrees of freedom.

# Spinel Family

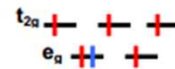
Kismarahardja, Dissertation 2010

**Table 2** The physical properties of transition-metal spinel vanadate.

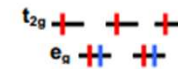
	$\rho$ at RT	structural transition and	magnetic ordering and	comments
				
CoV <sub>2</sub> O <sub>4</sub>	$65 \times 10^{-3}$	<div style="display: flex; align-items: center;"> <div style="background-color: red; width: 100px; height: 20px; margin-right: 5px;"></div> <div style="text-align: center;">C</div> </div> <div style="text-align: center;">no transition</div>	<div style="display: flex; align-items: center;"> <div style="background-color: lightblue; width: 100px; height: 20px; margin-right: 5px;"></div> <div style="text-align: center;">ferrimagnetic</div> </div> <div style="display: flex; align-items: center;"> <div style="background-color: red; width: 100px; height: 20px; margin-right: 5px;"></div> <div style="text-align: center;">para</div> </div> <div style="text-align: center;">152 T(K)</div>	Time decaying behaviour observed at low temperature.



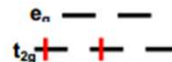
Mn<sup>2+</sup>, tetrahedral, high spin,  $e_g 2$ ,  $t_{2g} 3$   
 Magnetic moment =  $5.91 \mu_B$ ,  
 free electrons = 5



Fe<sup>2+</sup>, tetrahedral, high spin,  $e_g 3$ ,  $t_{2g} 3$   
 Magnetic moment =  $4.89 \mu_B$ ,  
 free electrons = 4



Co<sup>2+</sup>, tetrahedral, high spin,  $e_g 4$ ,  $t_{2g} 3$   
 Magnetic moment =  $3.87 \mu_B$ ,  
 free electrons = 3



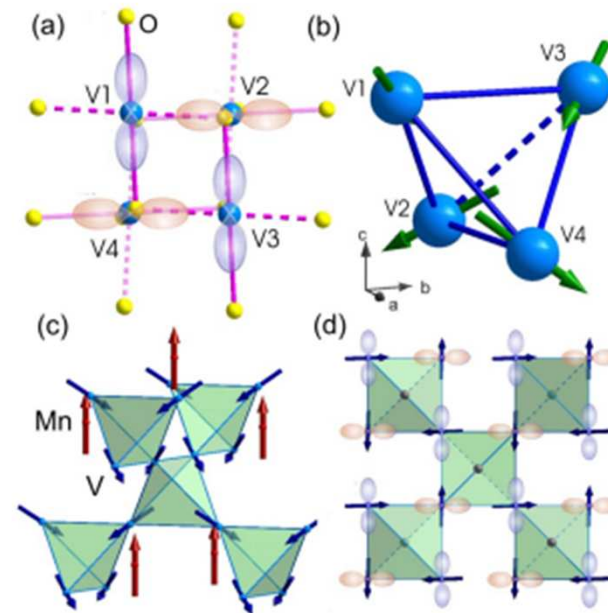
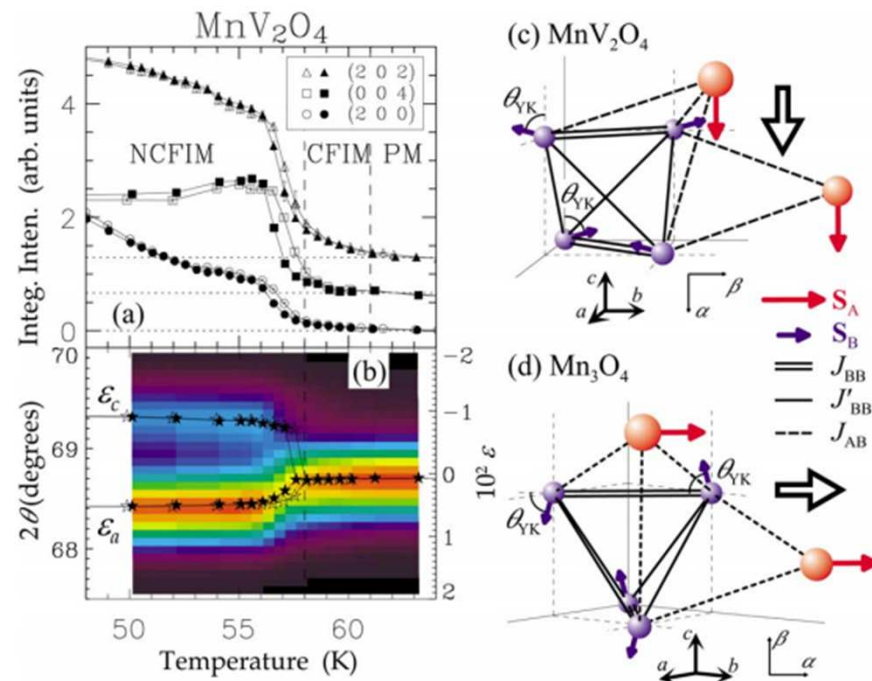
V<sup>3+</sup>, octahedral, high spin,  $e_g 2$ ,  $t_{2g} 3$   
 Magnetic moment =  $2.82 \mu_B$ ,  
 free electrons = 2



# MnV<sub>2</sub>O<sub>4</sub>: Magnetic structure

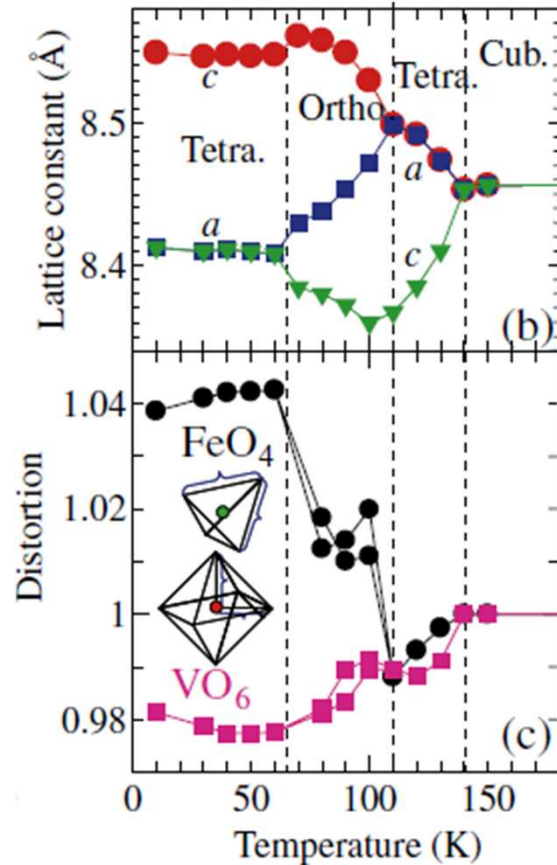
PHYSICAL REVIEW B **77**, 054412 (2008)

PRL **100**, 066404 (2008)



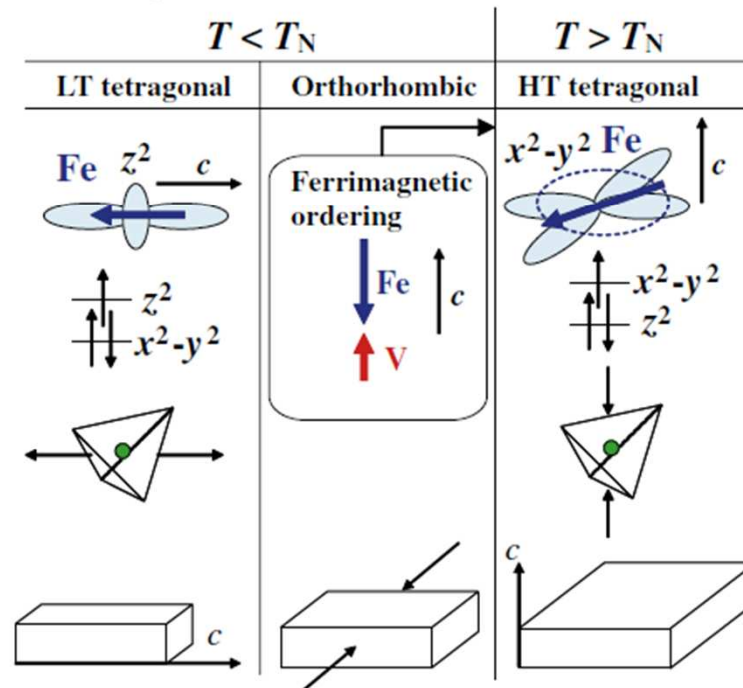
# FeV<sub>2</sub>O<sub>4</sub>: Magnetic structure

Journal of the Physical Society of Japan  
Vol. 77, No. 5, May, 2008, 053708  
©2008 The Physical Society of Japan



## Structural and Magnetic Properties of Spinel FeV<sub>2</sub>O<sub>4</sub> with Two Ions Having Orbital Degrees of Freedom

Takuro KATSUFUJI<sup>1,2,3</sup>, Takehito SUZUKI<sup>1</sup>, Haruki TAKEI<sup>1</sup>, Masao SHINGU<sup>1</sup>,  
Kenichi KATO<sup>4,5,6</sup>, Keiichi OSAKA<sup>5</sup>, Masaki TAKATA<sup>4,5,6</sup>,  
Hajime SAGAYAMA<sup>7</sup>, and Taka-hisa ARIMA<sup>7</sup>



# Magnetic excitations in $\text{MnV}_2\text{O}_4$

PHYSICAL REVIEW B 77, 054412 (2008)

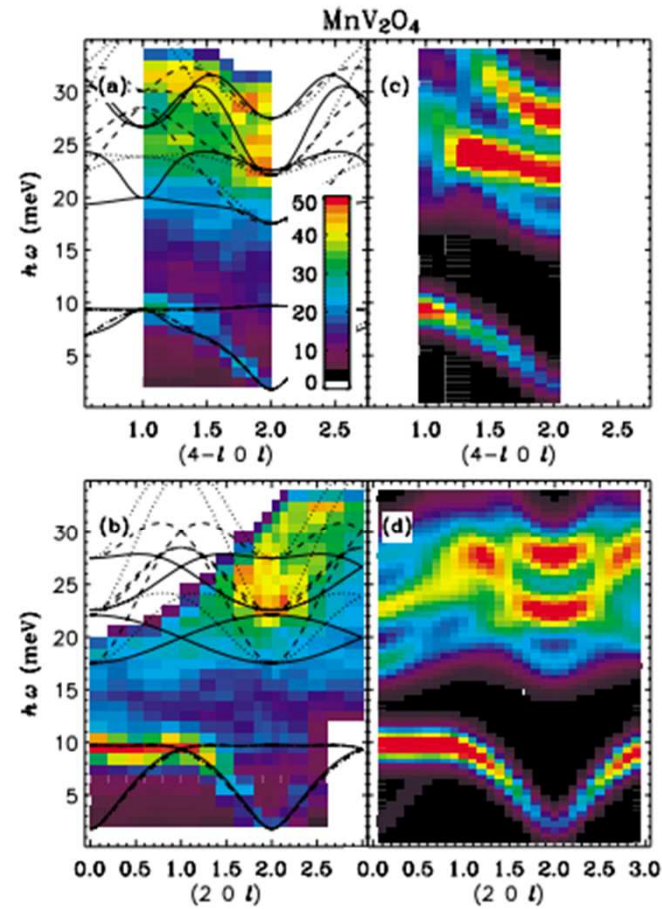
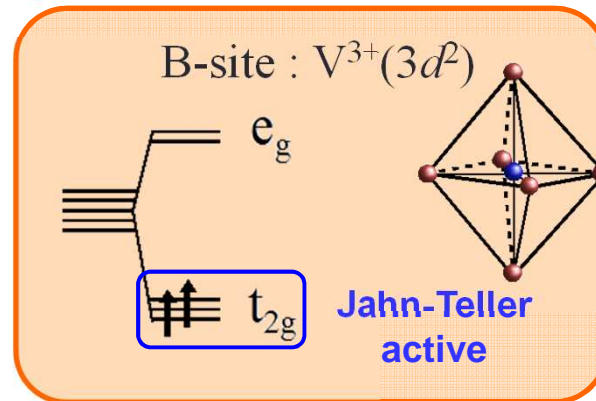
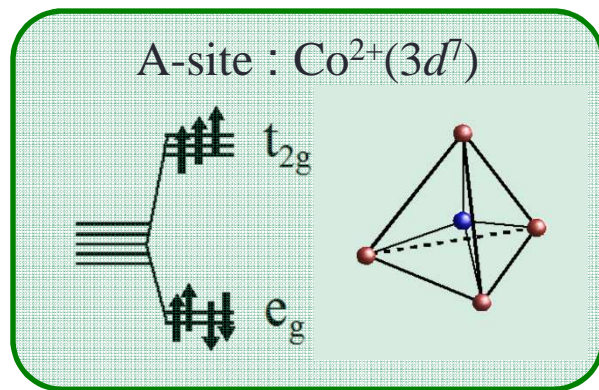
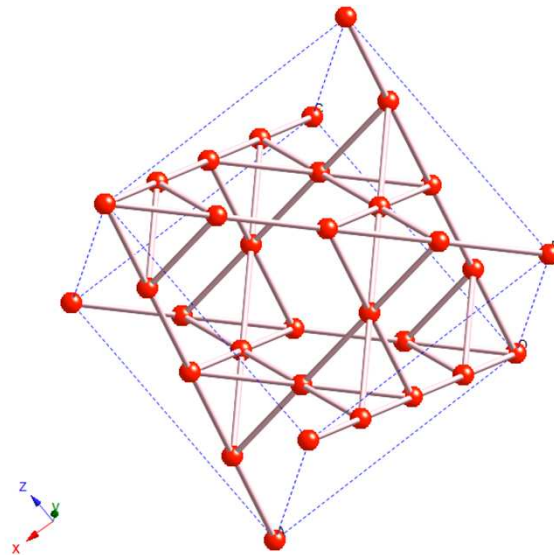
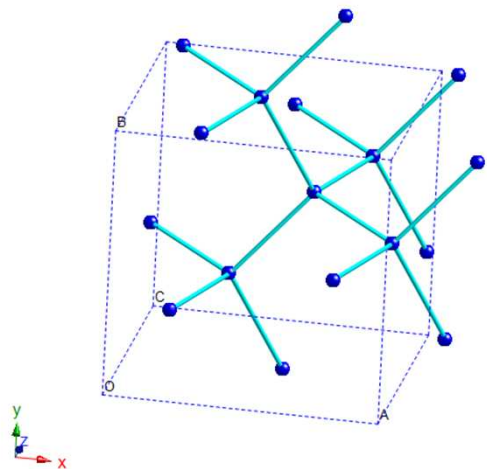


TABLE I. The optimal parameters used to calculate spin wave dispersions of  $\text{Mn}_3\text{O}_4$  [Fig. 2(b)] and  $\text{MnV}_3\text{O}_4$  (Fig. 3) ( $\tilde{J}$  and  $\tilde{D}$  are in meV).

	$\tilde{J}_{AB}$	$\tilde{J}_{BB}$	$\tilde{J}'_{BB}$	$\tilde{D}_A^{\tilde{1}\tilde{1}0}$	$\tilde{D}_B^z$	
Mn <sub>3</sub> O <sub>4</sub>	2.7(1)	19(1)	-1.1(7)	-0.1(3)	-0.28(3)	
	$\tilde{J}_{AB}$	$\tilde{J}_{BB}$	$\tilde{J}'_{BB}$	$\tilde{D}_A^z$	$\tilde{D}_B^x$	$\tilde{D}_B^y$
MnV <sub>2</sub> O <sub>4</sub>	2.8(2)	9.8(9)	3.0(8)	-0.6(4)	-4.0(4)	2.7(9)

# CoV<sub>2</sub>O<sub>4</sub>: Structure (Cubic with $a = 8.41 \text{ \AA}$ )

A-site ( $\text{Co}^{2+}$ ,  $S=3/2$ ) forms a diamond lattice while B-site ( $\text{V}^{3+}$ ,  $S=1$ ) forms a pyrochlore lattice.  $\text{V}^{3+}$  ions are Jahn-Teller active.





# CoV<sub>2</sub>O<sub>4</sub>: Detailed structure

PRL 106, 056602 (2011)

PHYSICAL REVIEW LETTERS

week ending  
4 FEBRUARY 2011

## Co[V<sub>2</sub>]O<sub>4</sub>: A Spinel Approaching the Itinerant Electron Limit

A. Kismarhardja,<sup>1,2</sup> J. S. Brooks,<sup>1,2</sup> A. Kiswandhi,<sup>1,2</sup> K. Matsubayashi,<sup>3</sup> R. Yamanaka,<sup>3</sup> Y. Uwatoko,<sup>3</sup> J. Whalen,<sup>1</sup>  
T. Siegrist,<sup>1,4</sup> and H. D. Zhou<sup>1,\*</sup>

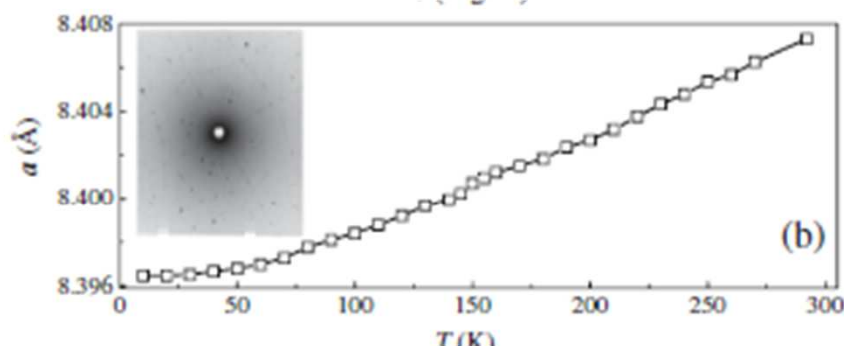


TABLE I. Room temperature crystallographic data for Co[V<sub>2</sub>]O<sub>4</sub>. (a)  $R_1 = \sum ||F_o| - |F_c|| / \sum |F_o|$ , (b)  $wR_2 = [\sum w(F_o^2 - F_c^2)^2 / \sum w(F_o^2)]^{1/2}$ ,  $w = [\sigma^2(F_o)^2 + (A \cdot p)^2 + B \cdot p]^{-1}$ , and  $p = (F_o^2 + 2F_c^2)/3$ ;  $A = 0.0067$ ,  $B = 0$ .

Space Group	$Fd\bar{3}m$ (No. 227)
$a$ (Å)	8.4073(1)
$Z$	8
Atom Positions, $U_{iso}$	Co 0.375, 0.006 70(13)
$(x = y = z)$	V 0, 0.005 68(12)
	O 0.239 79(10), 0.0071(2)
V (Å <sup>3</sup> )	594.251(12)
$\rho_{cal}$ (g/cm <sup>3</sup> )	5.026
$\mu$ (mm <sup>-1</sup> )	11.497
Data Collection Range (deg)	$8.06 < \theta < 61.47$
Reflections Collected	7124
Independent Reflections	260 [ $R_{int} = 0.097$ ]
Parameter Refined	8
$R_1, wR_2$ ( $F_o > 4\sigma F_o$ )	0.0370, 0.1035
$R_1, wR_2$ (All Data)	0.0398, 0.1015
Goodness-of-Fit	1.112

$Fd\bar{3}m$

$O_h^7$

$m\bar{3}m$

Cubic

No. 227

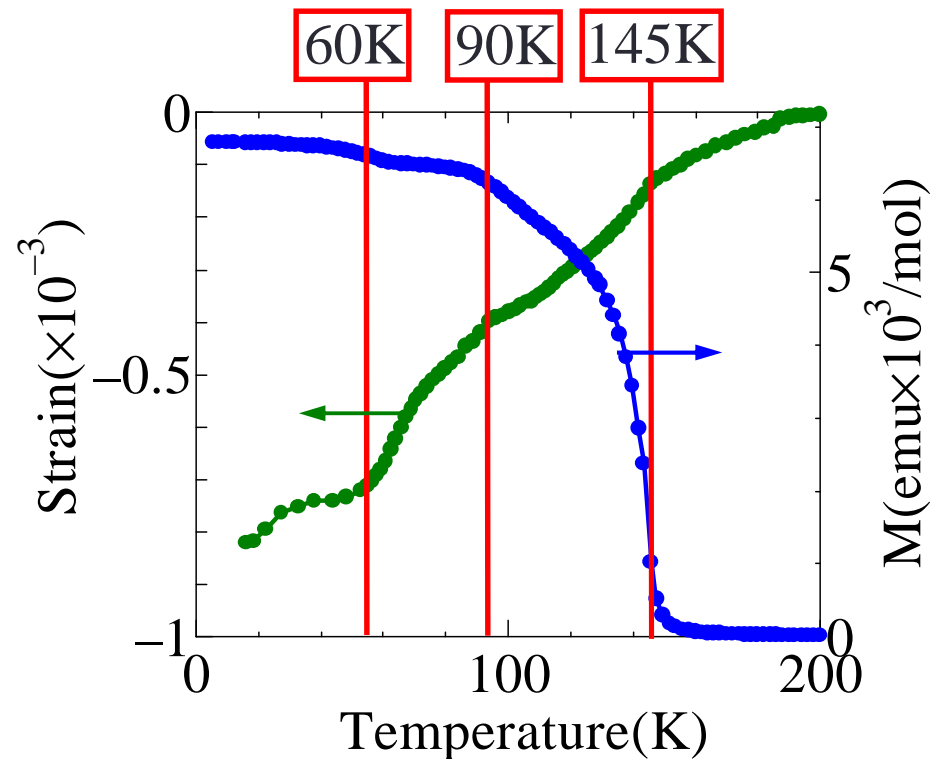
$F 4_1/d \bar{3} 2/m$

Patterson symmetry  $Fm\bar{3}m$

ORIGIN CHOICE 2

# CoV<sub>2</sub>O<sub>4</sub>: Magnetism

CoV<sub>2</sub>O<sub>4</sub> shows a ferrimagnetic transition at  $T = 145$  K. And there may be two additional transitions at  $T = 60$  and  $90$  K (spin re-orientation ?).

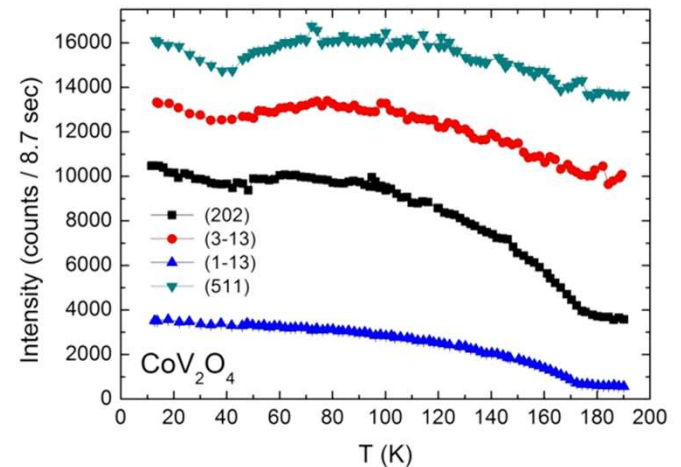
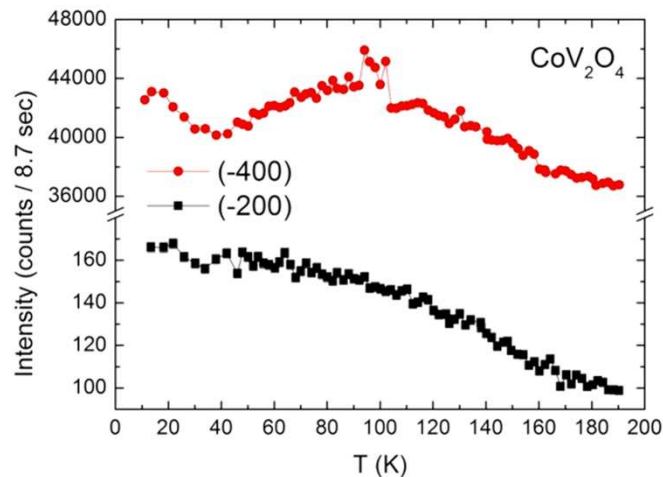
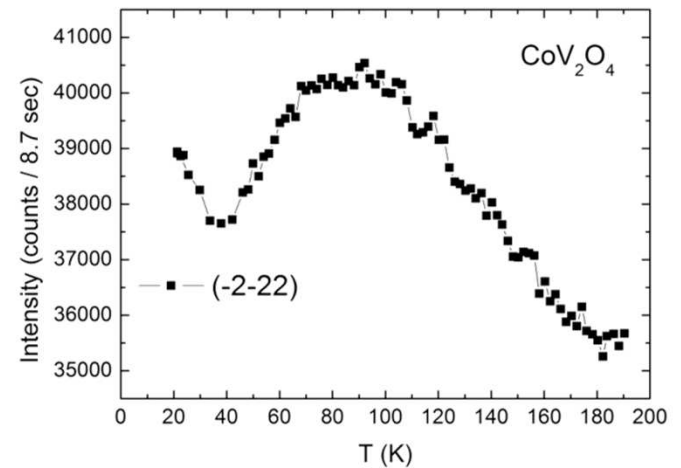
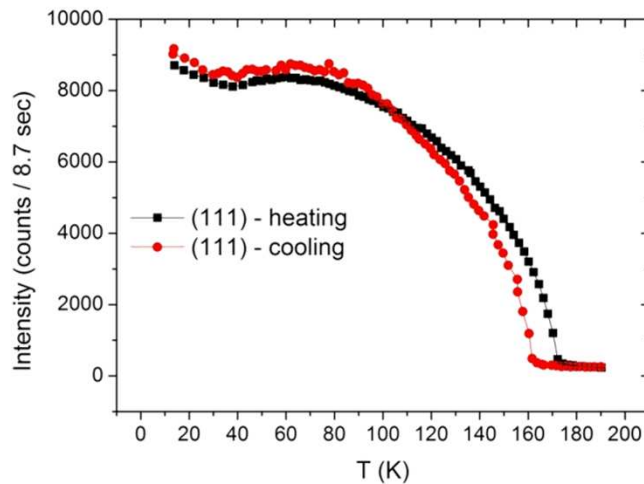




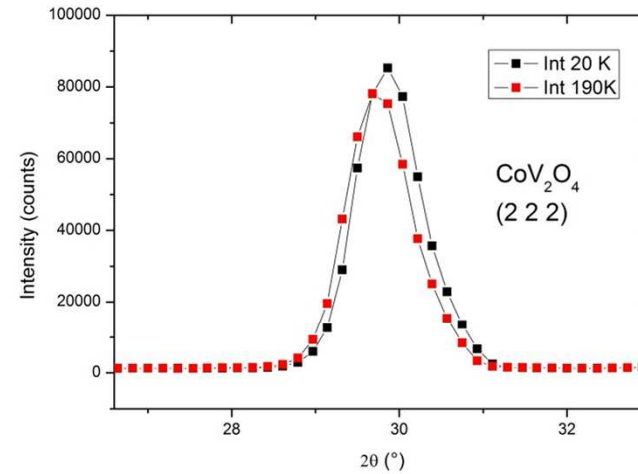
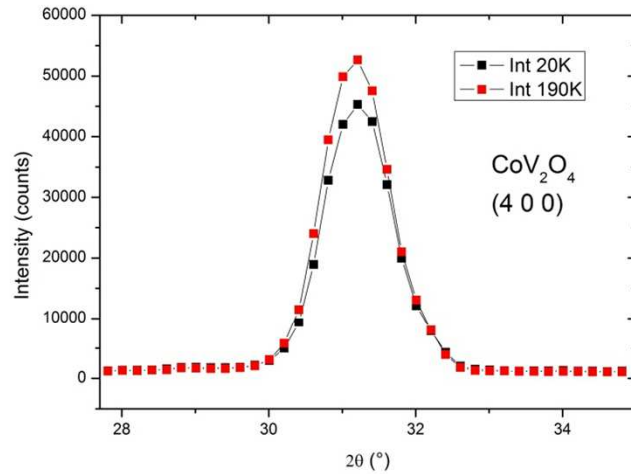
# Experiments

- E5, HZB
  - Single Crystal Diffractometer
  - $T = 4\text{K}$  with wavelength  $1\text{ \AA}$
- HB1A, HFIR
  - Thermal Triple axis spectrometer
  - $T = 4\text{K}$  with  $E_f = 65\text{meV}$

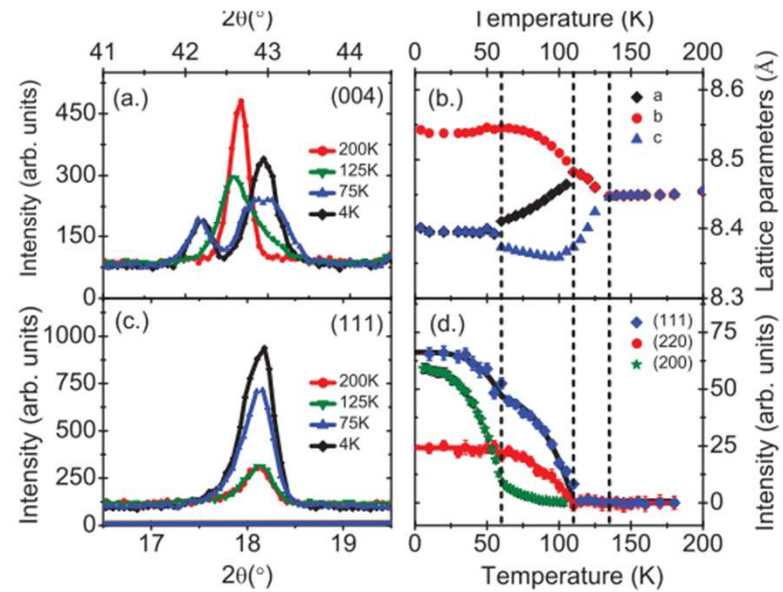
# Temperature dependence of nuclear + magnetic peaks



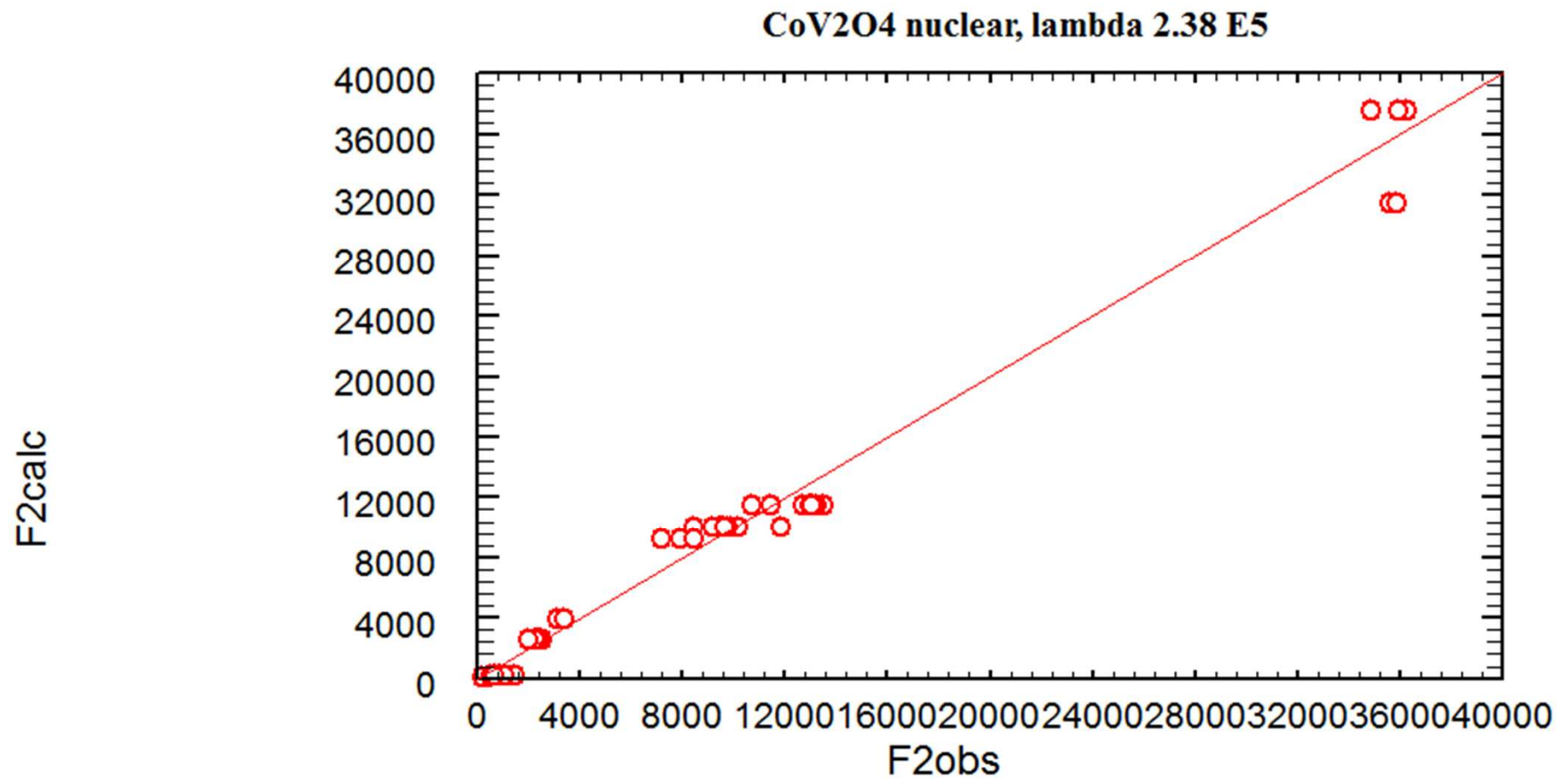
# No structural transition



FeV<sub>2</sub>O<sub>4</sub>

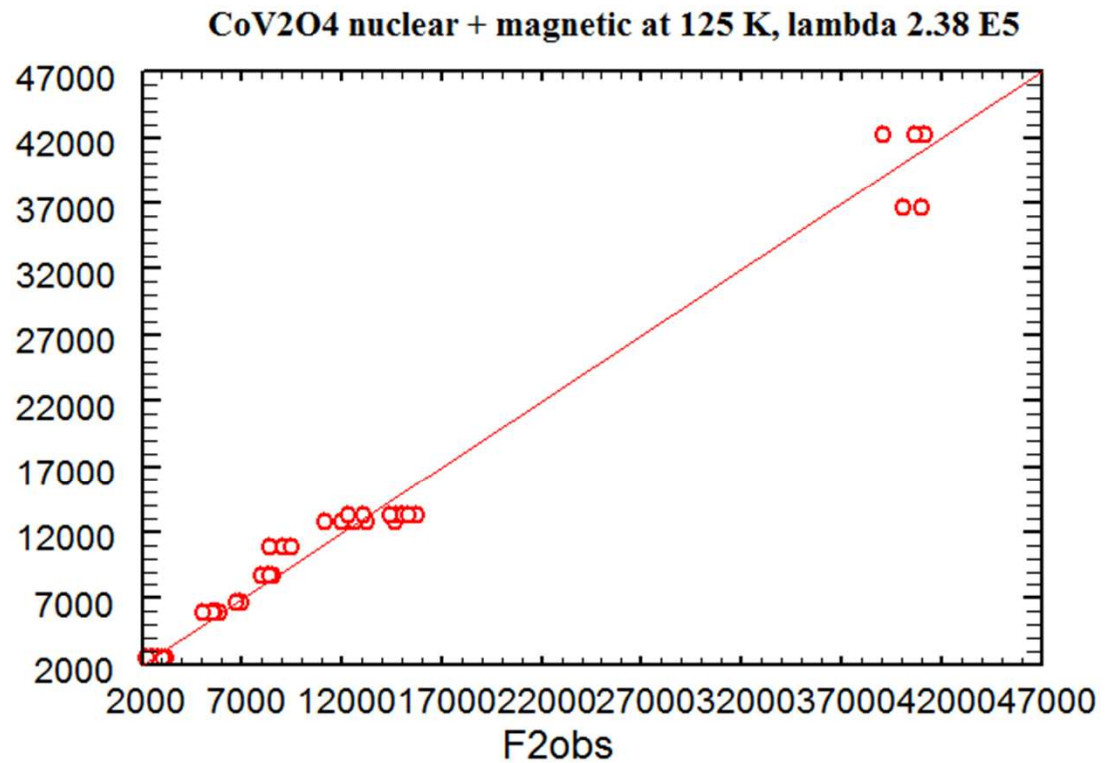
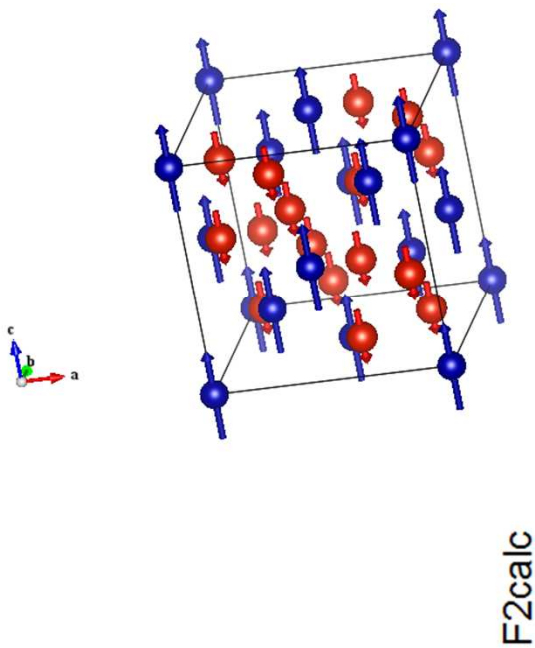


# Refinement – 200 K (nuclear)

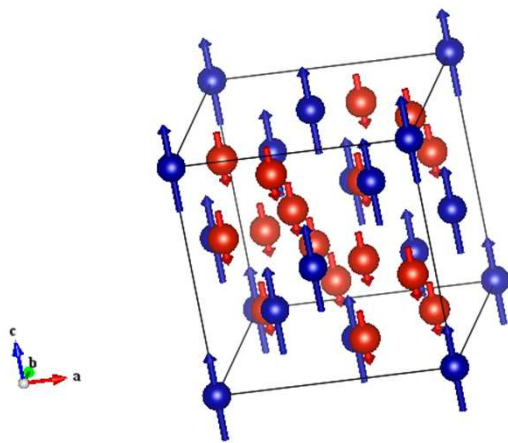


# Refinement – 115 K (magnetic)

Co	V
$2.58227 \mu_B$	$-0.68620 \mu_B$



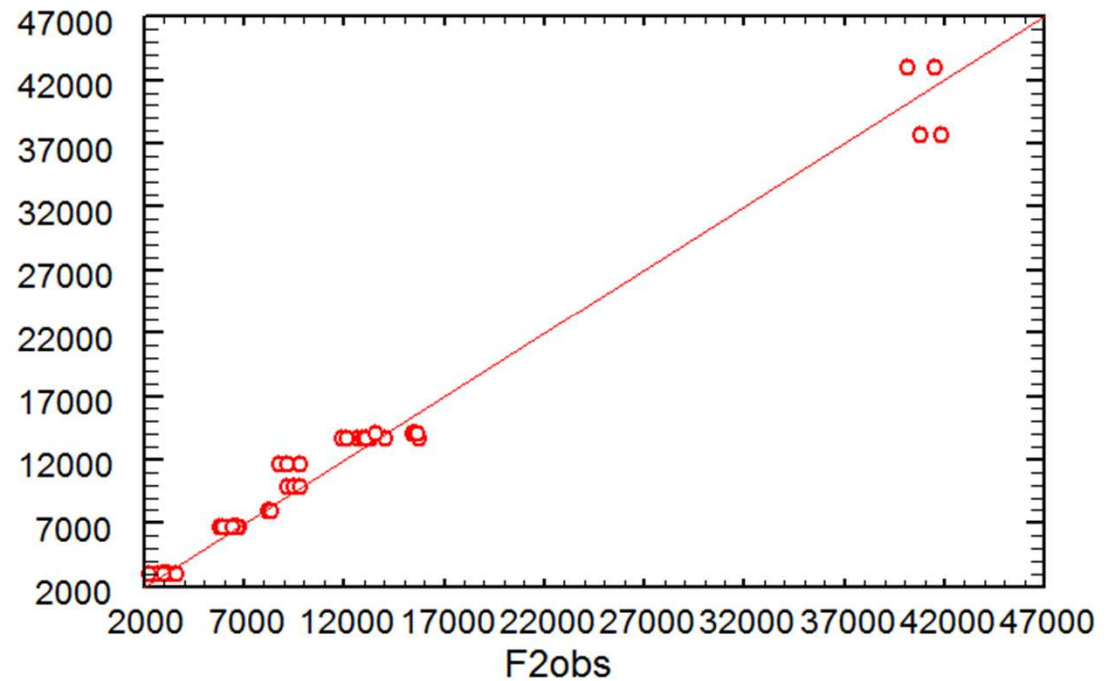
# Refinement – 75 K (magnetic)



F2calc

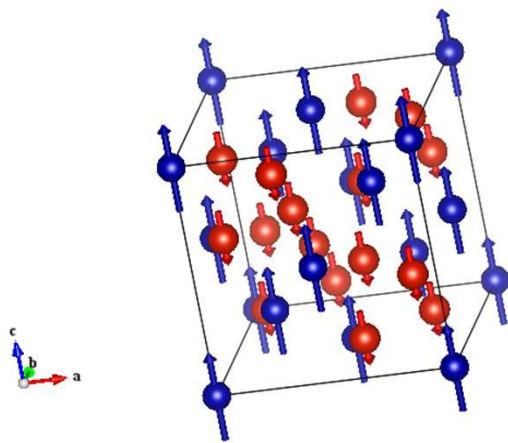
Co	V
$2.95270 \mu_B$	$-0.82596 \mu_B$

CoV2O4 nuclear + magnetic at 75 K, lambda 2.38 E5





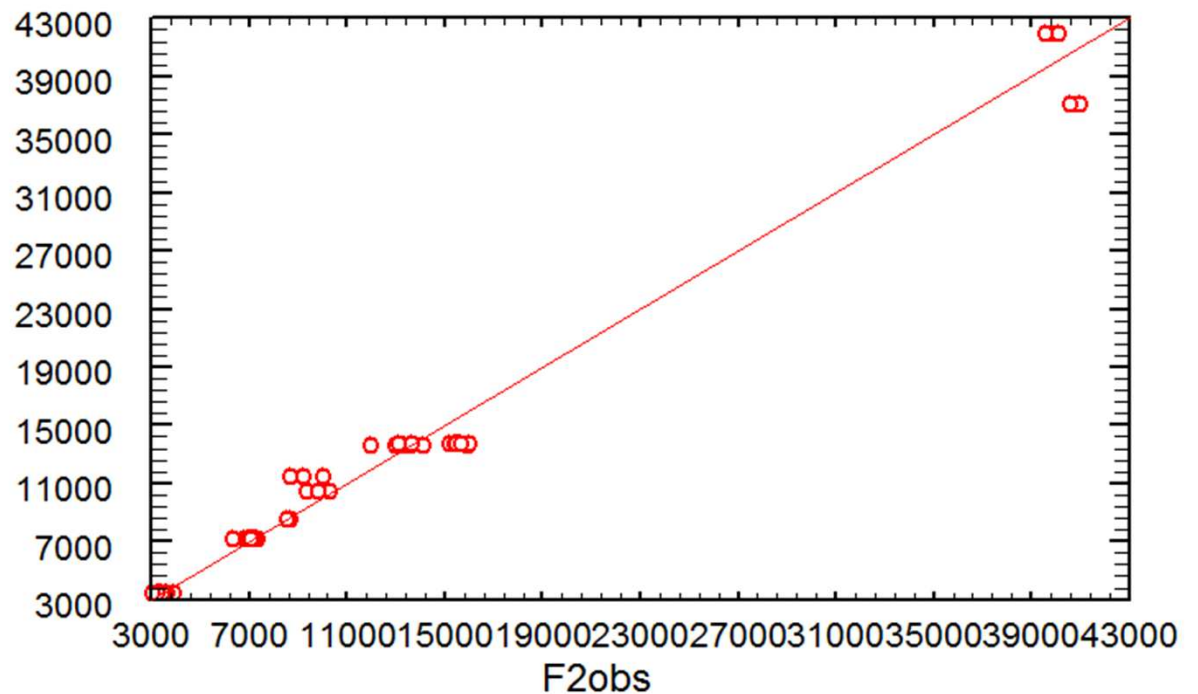
# Refinement – 10 K (magnetic)



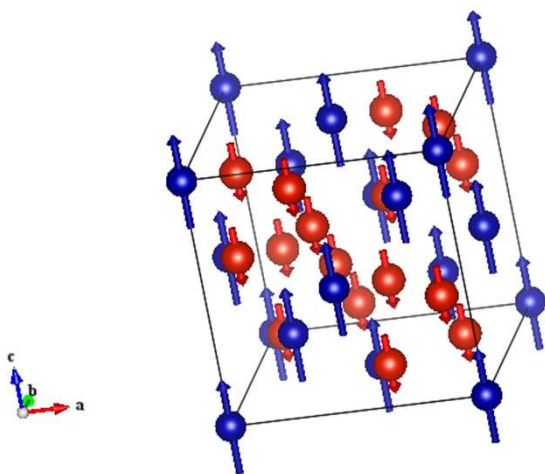
F2calc

Co	V
$3.27383 \mu_B$	$-0.87081 \mu_B$

CoV2O4 nuclear + magnetic at 10 K, lambda 2.38 E5

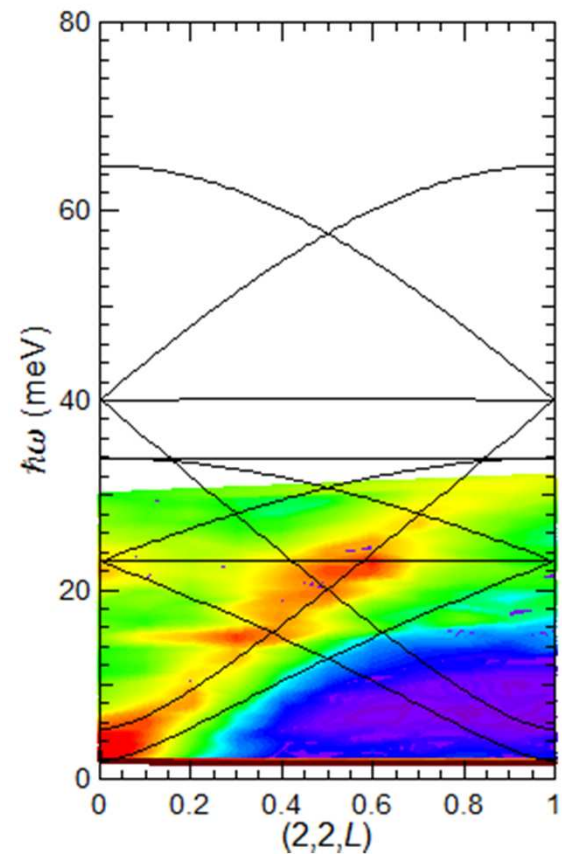
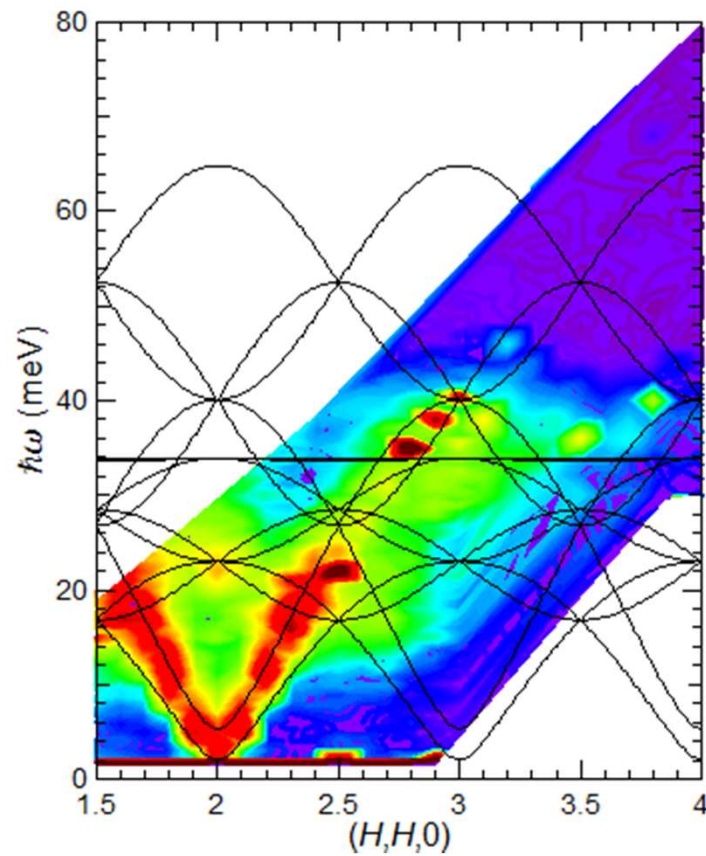


# Refinement – Magnetic Structure



Temp (K)	Co Moment ( $\mu_B$ )	V Moment ( $\mu_B$ )
115 K	2.582(1)	0.686(1)
75 K	2.952(2)	0.825(1)
10 K	3.273(2)	0.870(1)

# Magnetic excitations



# Linear Spinwave calculation

$$\begin{aligned}\mathcal{H} = & \sum_{i,j} J_{Co-V} \mathbf{S}_{i,Co} \cdot \mathbf{S}_{j,V} + \sum_{i,j} J_{Co-Co} \mathbf{S}_{i,Co} \cdot \mathbf{S}_{j,Co} + \sum_{i,j} J_{V-V} \mathbf{S}_{i,V} \cdot \mathbf{S}_{j,V} \\ & + \sum_i D_{Co} (S^z_{i,Co})^2 + \sum_i D_V (S^z_{i,V})^2\end{aligned}$$

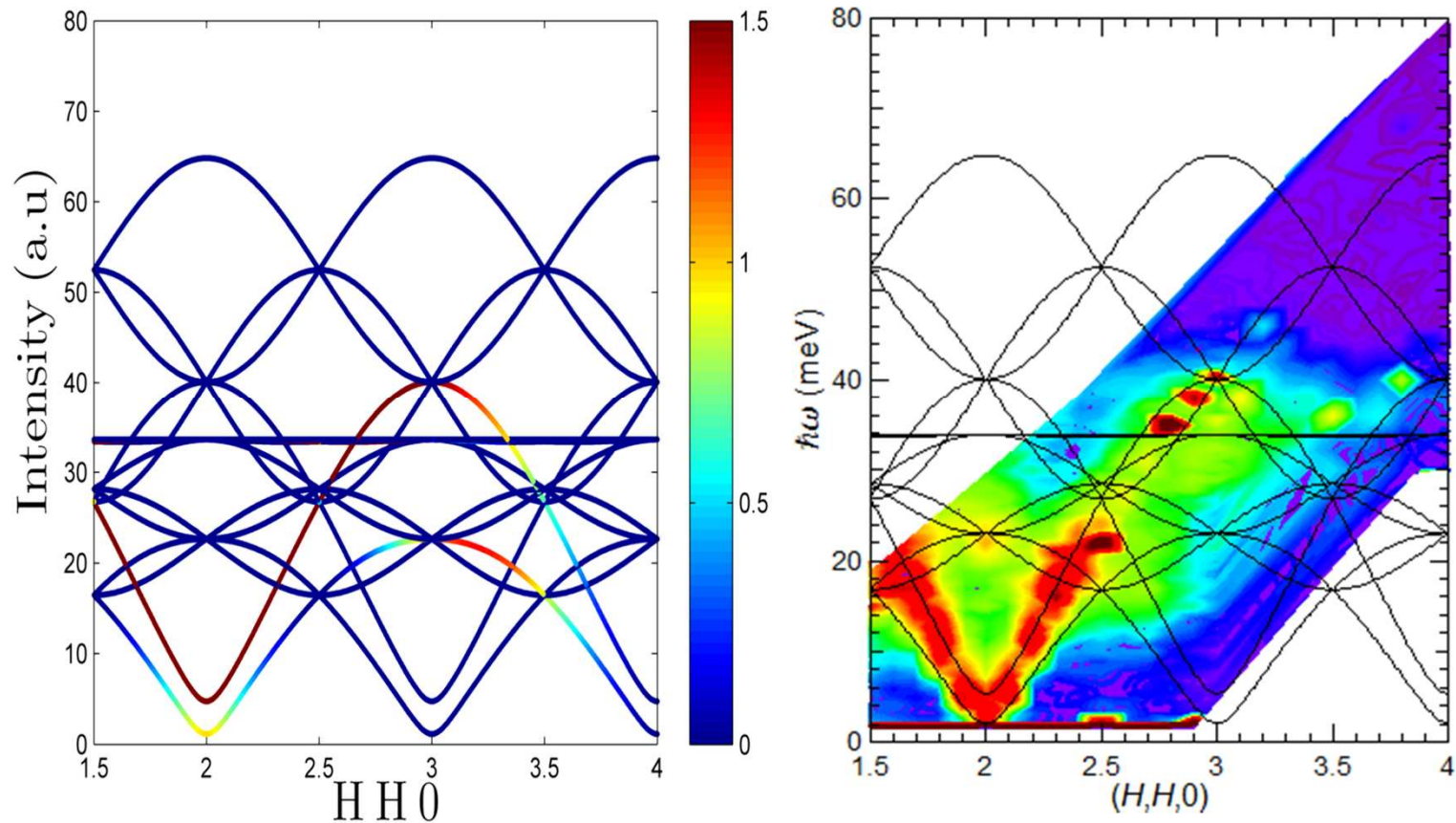
$J_{Co-Co}$ ,  $J_{V-V}$  and  $J_{Co-V}$  are nearest neighbor interactions.  $D_{Co}$  and  $D_V$  are the single ion anisotropy for Co and V along z direction.

A good fit with the observed magnetic excitations were obtained when  $J_{Co-V} = 1.3$  meV,  $J_{V-V} = -2.7$  meV,  $J_{Co-Co} = -4.1$  meV,  $D_V = -0.32$  meV and  $D_{Co} = -0.02$  meV.

# Linear Spinwave calculation

$J_{Co-V} = 1.3 \text{ meV}$  ,  $J_{V-V} = -2.7 \text{ meV}$  ,  $J_{Co-Co} = -4.1 \text{ meV}$

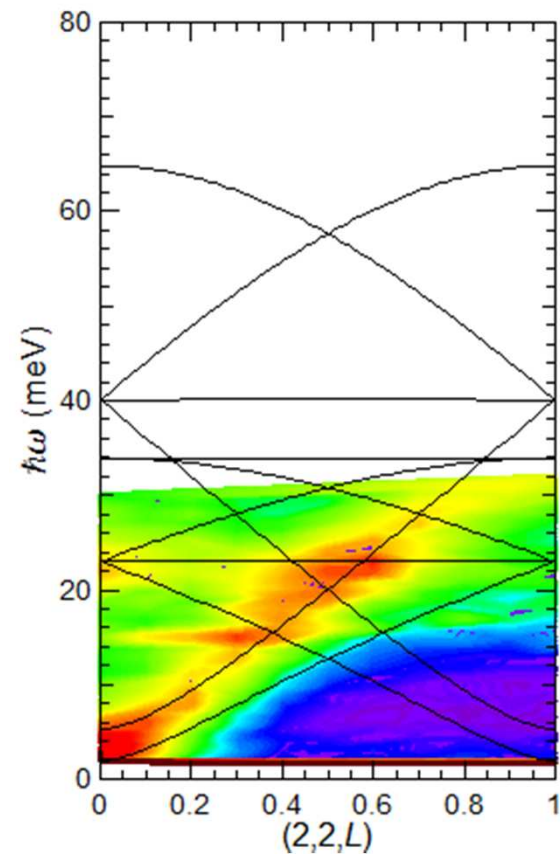
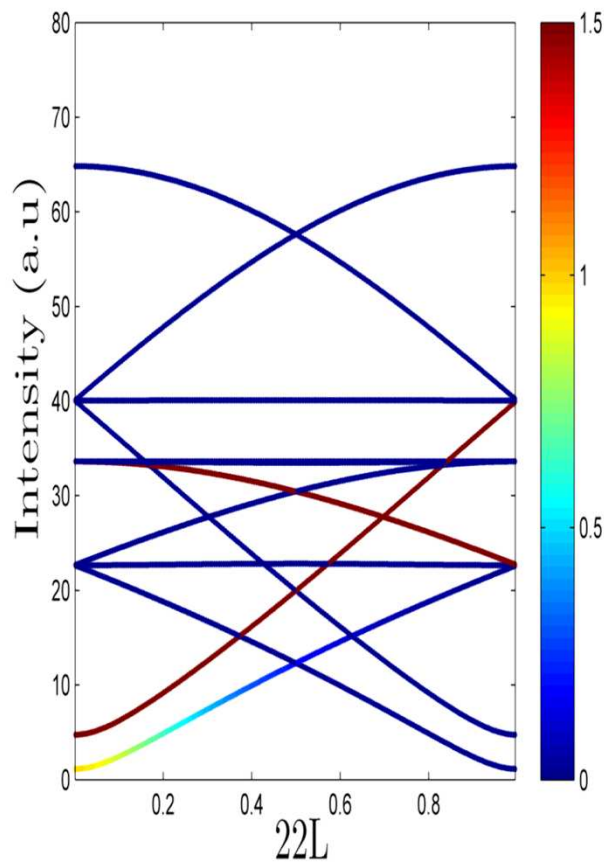
$D_V = -0.32 \text{ meV}$  ,  $D_{Co} = -0.02 \text{ meV}$ .





# Linear Spinwave calculation

$$J_{Co-V} = 1.3 \text{ meV}, J_{V-V} = -2.7 \text{ meV}, J_{Co-Co} = -4.1 \text{ meV}$$
$$D_V = -0.32 \text{ meV}, D_{Co} = -0.02 \text{ meV}.$$



# Summary and future work

- $\text{CoV}_2\text{O}_4$  does not show any structural transition as observed in  $\text{MnV}_2\text{O}_4$  and  $\text{FeV}_2\text{O}_4$ .
- Single crystal data can be successfully refined with a collinear ferrimagnetic structure.
- Collinear ferrimagnetic structure can be used to reproduce inelastic data using a Hamiltonian with nearest neighbor interactions and single ion anisotropy.
- Origin of the anomaly around 40 K is yet to be understood.



THANK YOU !

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